

Neutrino oscillations and beyond

J W F Valle
IFIC/CSIC – U Valencia



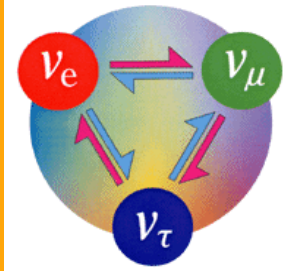
<http://astroparticles.es/>

Mini Workshop on Neutrinos, IPMU, Nov 8-11, 2010

LEPTON MIXING MATRIX

$$\mathbf{K} = \omega_{23} \cdot \omega_{13} \cdot \omega_{12}$$

Schechter & JV PRD22 (1980) 2227, PRD23 (1981) 1666
PDG



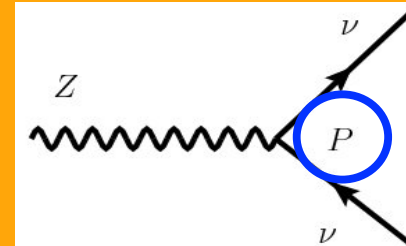
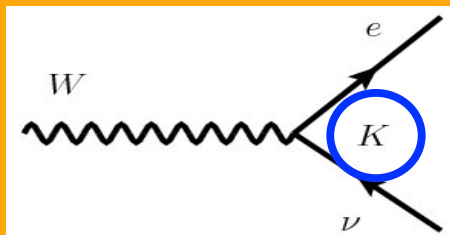
$$\begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & e^{i\phi_{23}} s_{23} \\ 0 & -e^{-i\phi_{23}} s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & e^{i\phi_{13}} s_{13} \\ 0 & 1 & 0 \\ -e^{-i\phi_{13}} s_{13} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & e^{i\phi_{12}} s_{12} & 0 \\ -e^{-i\phi_{12}} s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

23=atm+acc

13=reactor + ..

12=solar+KL

- Even in such simplest unitary form \mathbf{K} differs from quark mixing matrix, with two extra (Majorana) phases

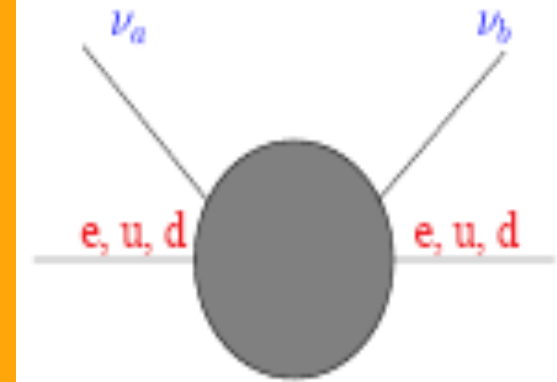


- Nontrivial structure of CC & NC in seesaw implies new phases & angles affecting neutrino propagation & inducing LFV

NONSTANDARD NEUTRINO INTERACTIONS

from non-trivial seesaw mixing matrix

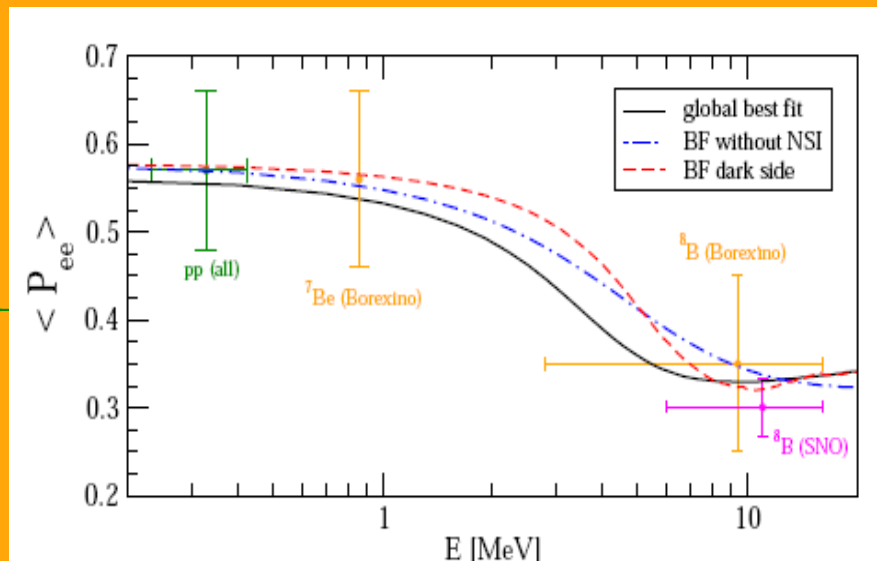
Schechter & JV PRD22 (1980) 2227



effective **non-unitarity** affects **propagation, source and detection**, e.g.

Resonant Oscillations Of Massless Neutrinos In Matter. Valle PLB199 (1987) 432

propagation



NSI-OSCILLATION INTERPLAY manifest @ Long-baseline studies

Fornengo et al, PRD65:013010,2002

Huber, Schwetz, JV

PRL88:101804,2002

PRD66 013006, 2002

Davidson et al JHEP (2003) 0303:011

Barranco, et al , D73 (2006) 113001, D77 (2008) 093014

Abada, Biggio Bonnet, Gavela, Hambye PRD78

Esteban Huber JV PLB668:197201,2008

Gavela, Hernandez, Ota, Winter, PRD79

Escrihuela et al PRD80:105009,2009; Err-D80:129908,2009

Malinsky et el, arXiv:0905.2889 [hep-ph]

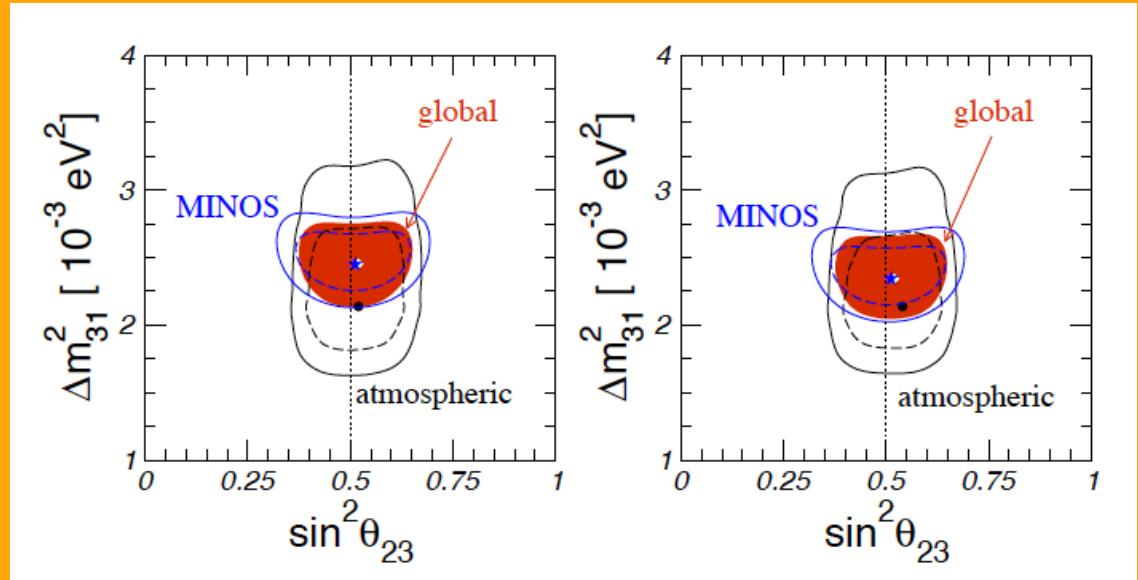
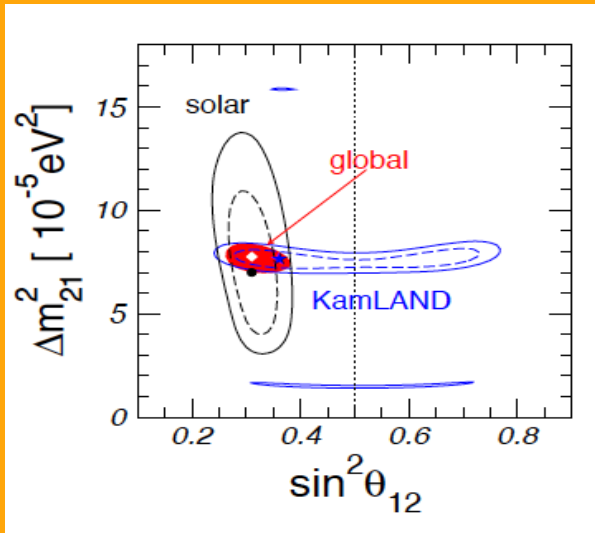
Bolaños et al PRD79 (2009) 113012

Kopp, Machado, Parke, arXiv:1009.0014

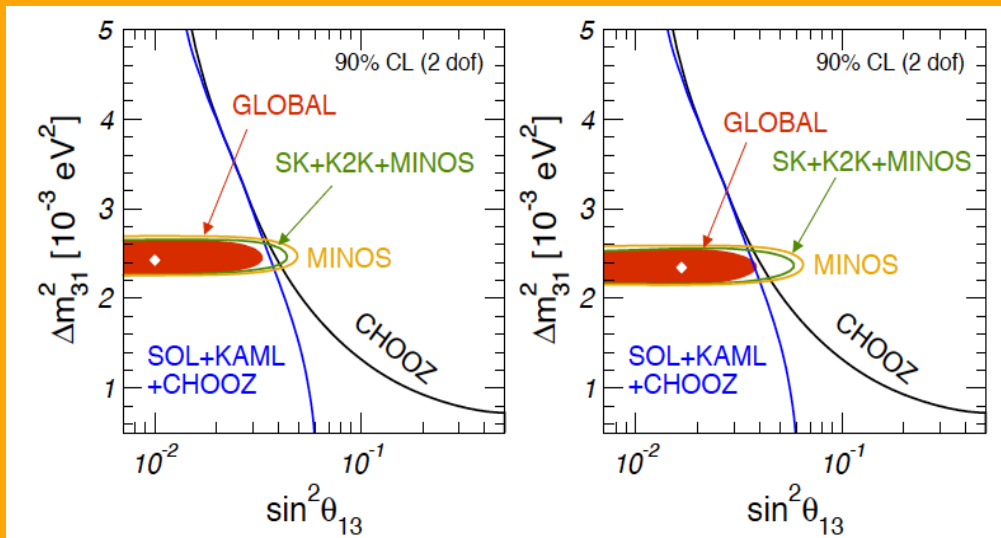
IN ANALYSIS OF CURRENT OSCILLATION DATA WE NEGLECT NSI

NEUTRINO OSCILLATIONS STATUS @ NUFACT2010

Update of Schwetz et al, NJP 10 (2008) 113011



Homestake, SAGE
 GALLEX/GNO,
 SK-III
 SNO-Ieta, SSM
 Borexino
 KamLAND (180 Km)



... SK-III
 thanks to SK ... E. Kearns

K2K (250 Km)
 MINOS latest
 (735 Km)

UPDATED THETA13

Update of NJP 10 (2008) 113011

We obtain the following bounds at 90% (3σ) CL:

$$\sin^2 \theta_{13} \leq \begin{cases} 0.054 (0.080) & (\text{solar+KamLAND}) \\ 0.040 (0.065) & (\text{MINOS NH}) \\ 0.053 (0.085) & (\text{MINOS IH}) \\ 0.025 (0.042) & (\text{CHOOZ+atm+K2K+MINOS, NH}) \\ 0.033 (0.052) & (\text{CHOOZ+atm+K2K+MINOS, IH}) \\ 0.029 (0.042) & (\text{global data, NH}) \\ 0.034 (0.049) & (\text{global data, IH}) \end{cases}$$

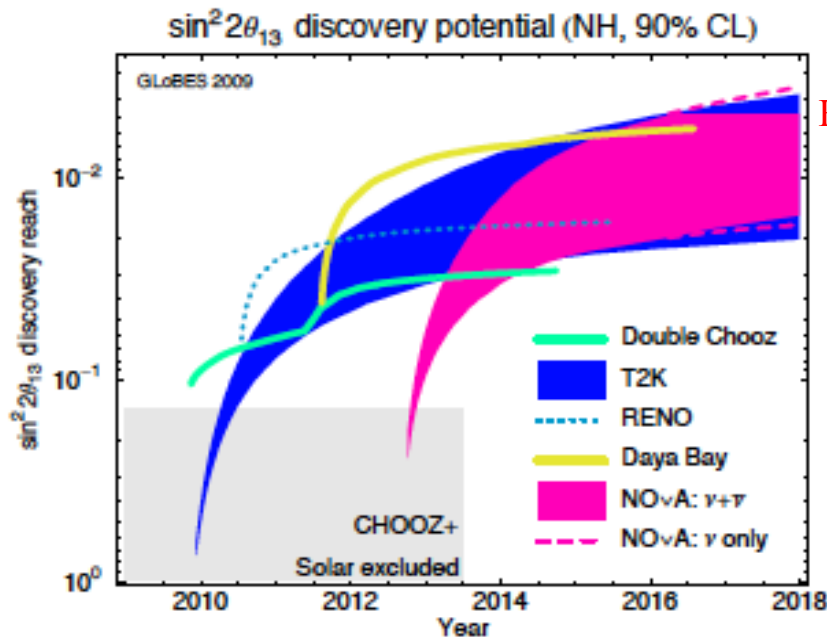
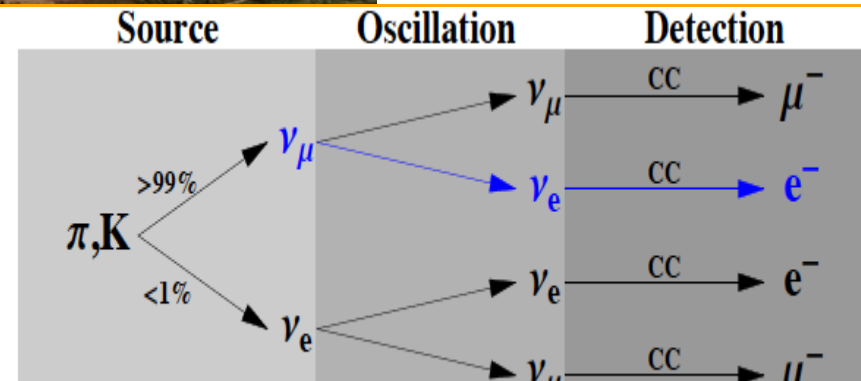
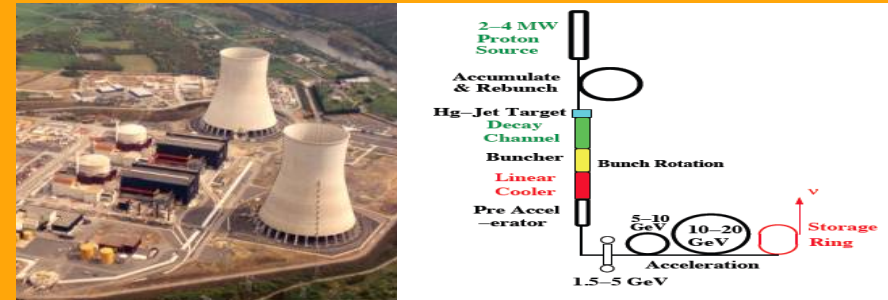
IH

$$\sin^2 \theta_{13} = 0.017 \pm 0.010$$

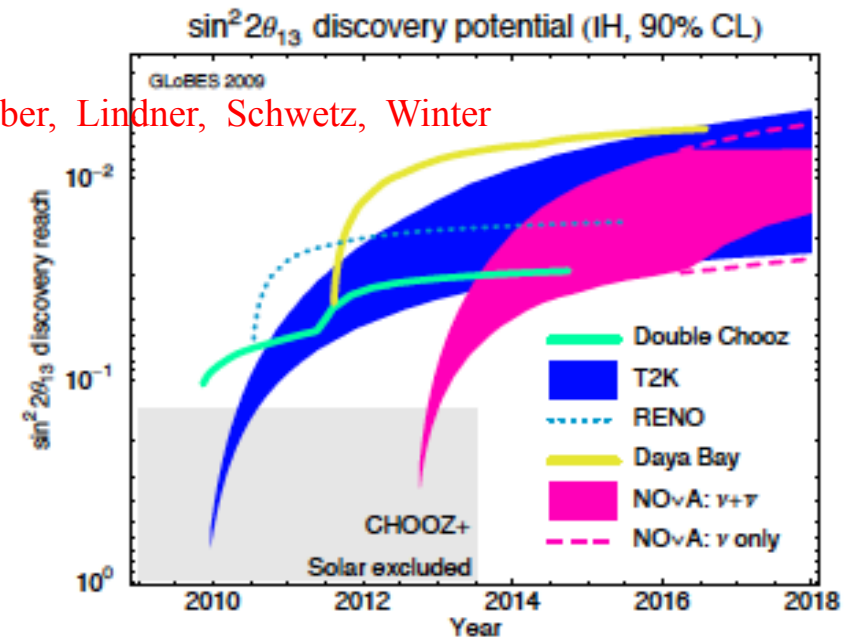
NH

$$\sin^2 \theta_{13} = 0.011^{+0.011}_{-0.007}$$

FUTURE

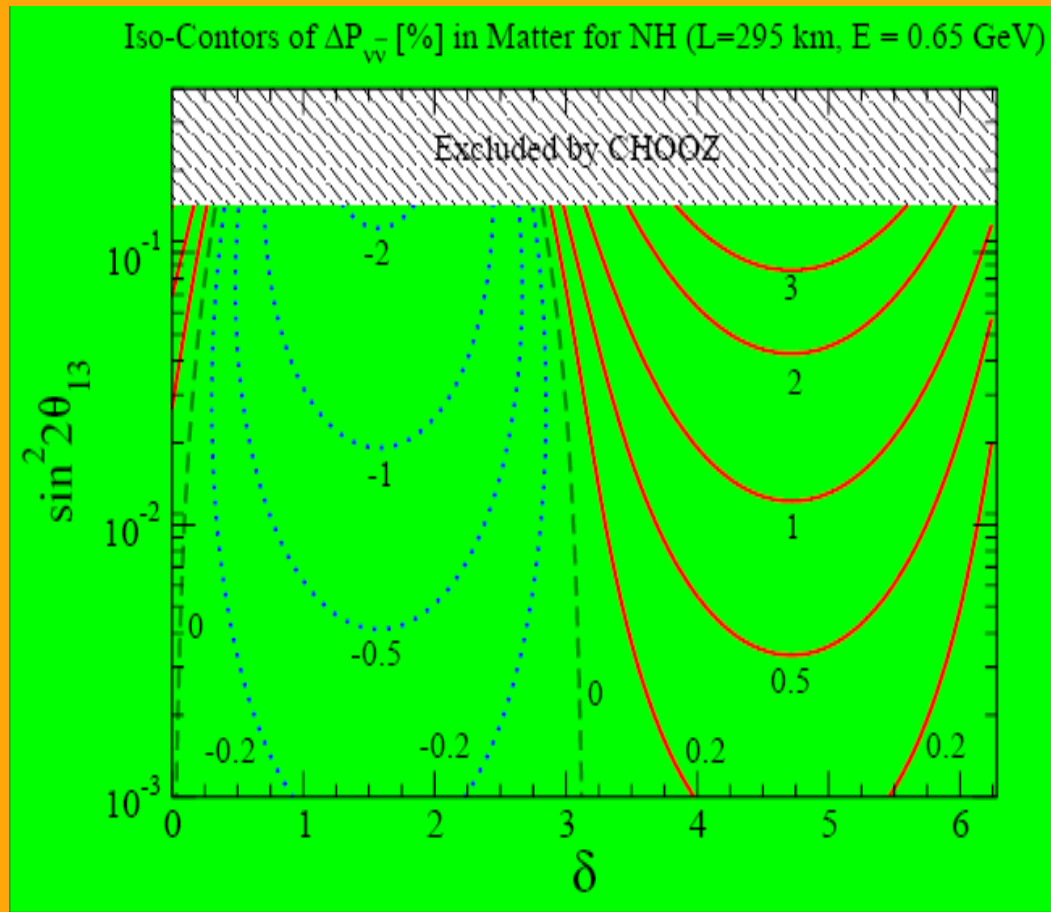


Huber, Lindner, Schwetz, Winter



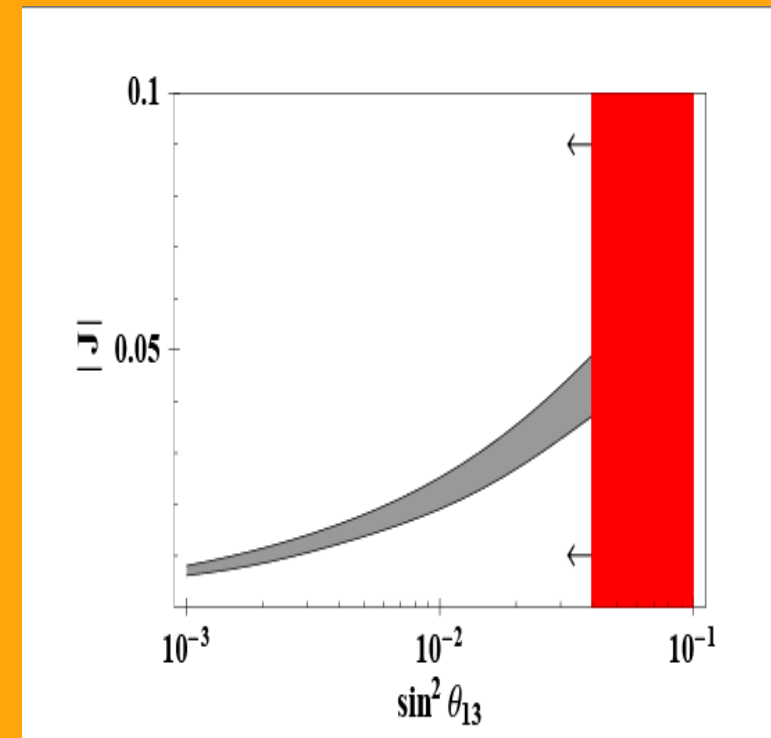
CPV & NU-OSCILLATIONS AT LONG BASELINES

asymmetries at % level



maximal CPV

Hirsch et al PRL99(2007)151802

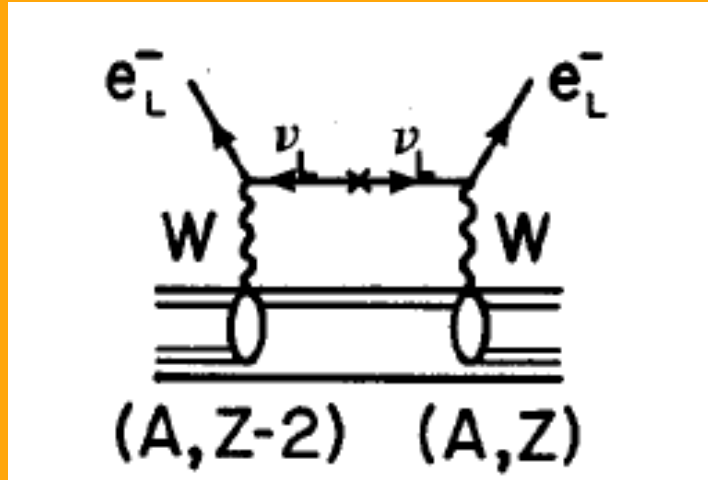


Nunokawa et al
Prog.Part.Nuc.Phys. 60 (2008) 338

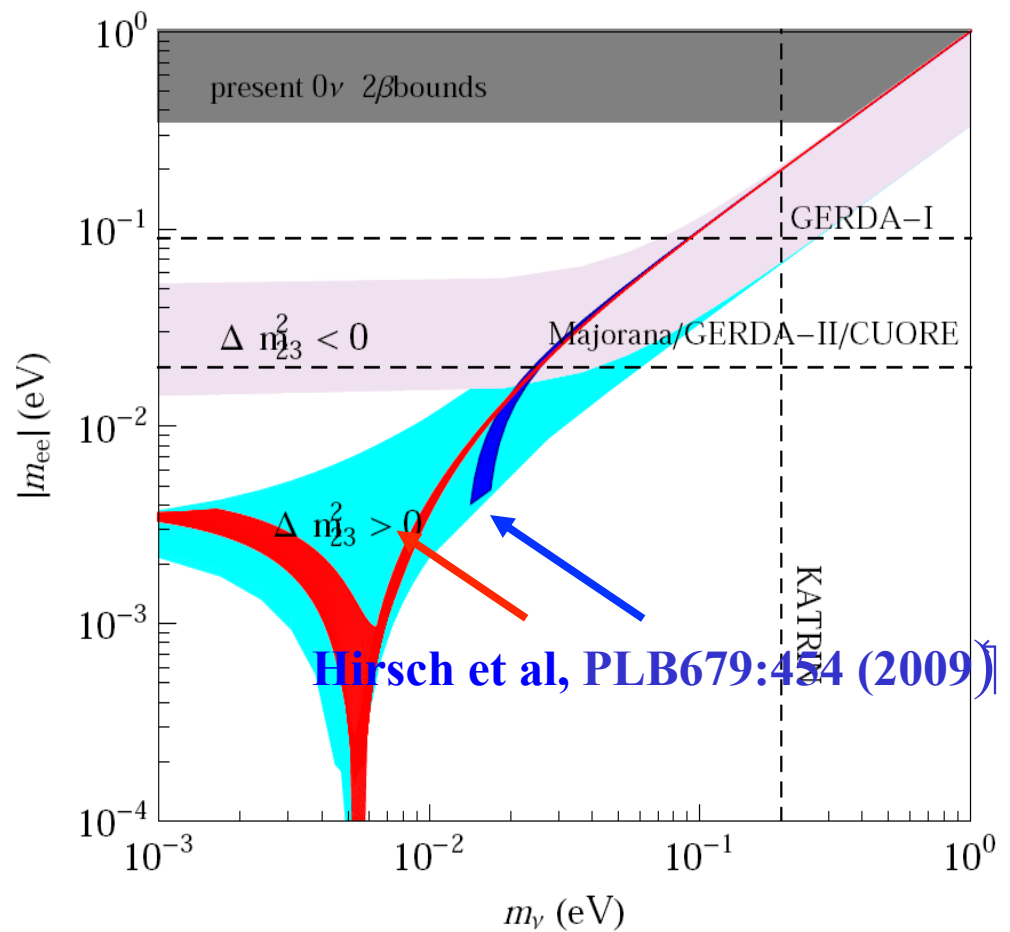
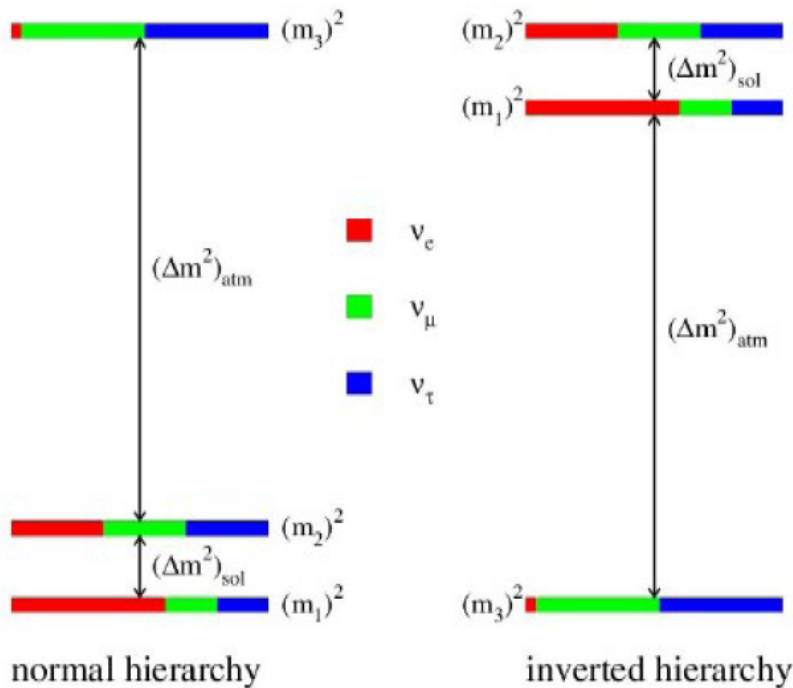
ISS report S. King et al Rep. Prog. Phys.

TESTING NEUTRINO SPECTRA WITH DBD

NH VERSUS IH



- **SBD & Cosmology ...**
- **Flavor sensitivity**

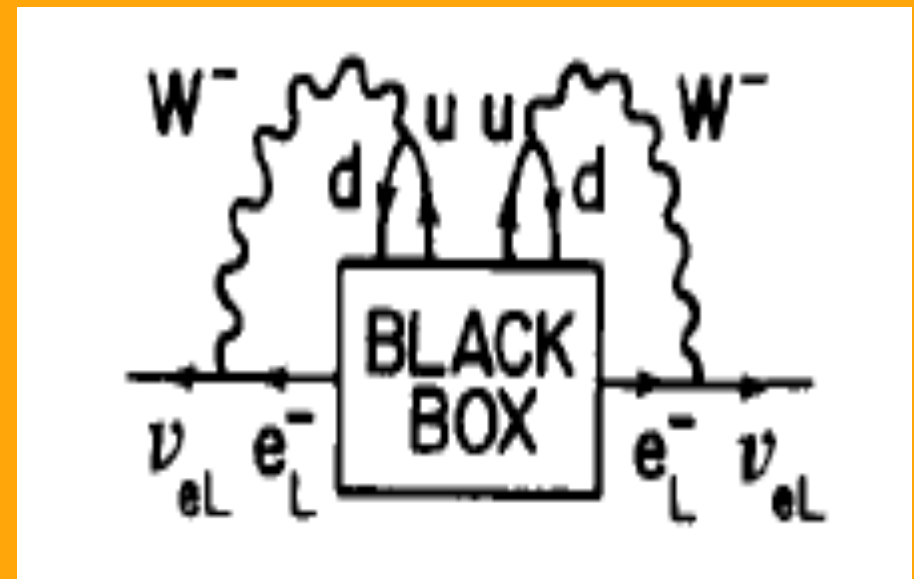


NEUTRINOLESS DOUBLE BETA DECAY

0-nu
DBD

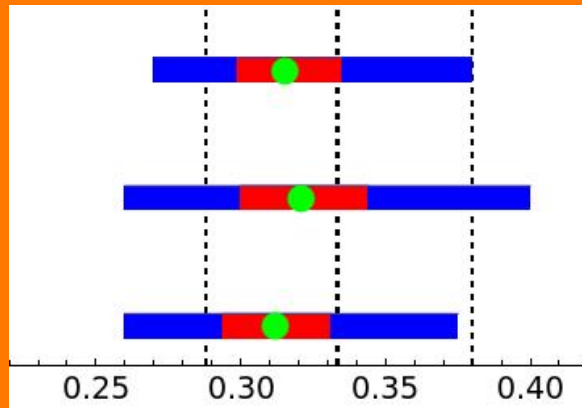


majorana

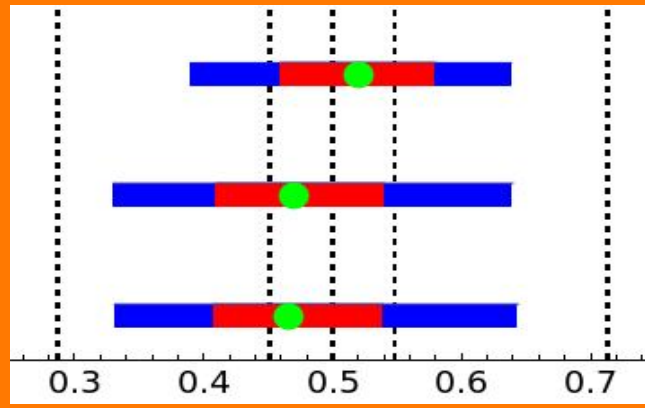


Schechter, JV PRD25 (1982) 2951

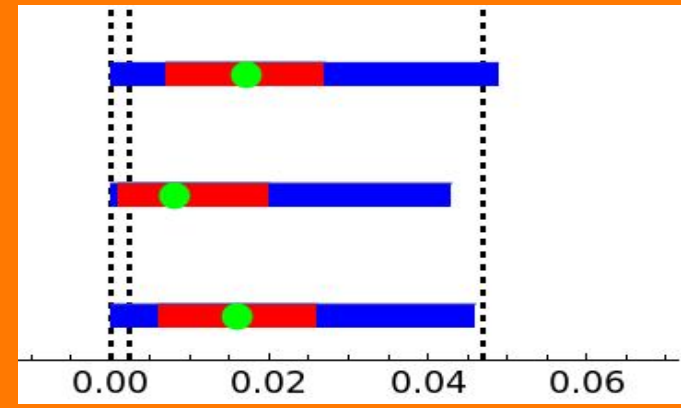
IN SHORT: OSCILLATION ANALYSES CONVERGE ...



$$\sin^2 \theta_{12} = 1/3$$



$$\sin^2 \theta_{23} = 0.5$$



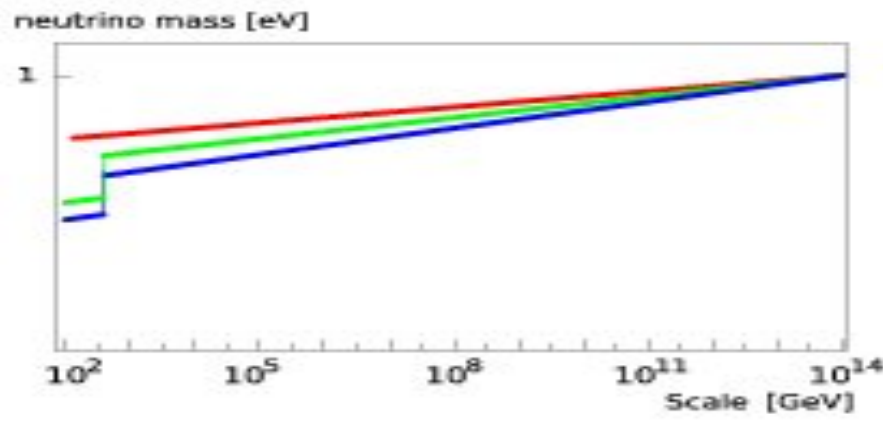
$$\sin^2 \theta_{13} = 0$$

TRI-BIMAXIMAL MIXING

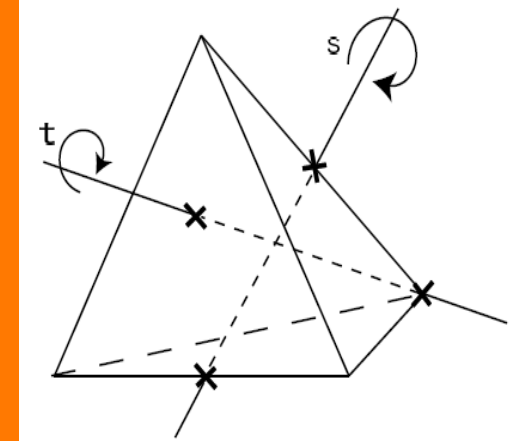
Harrison, Perkins & Scott

$$U_{\text{HPS}} = \begin{pmatrix} \sqrt{2/3} & 1/\sqrt{3} & 0 \\ -1/\sqrt{6} & 1/\sqrt{3} & -1/\sqrt{2} \\ -1/\sqrt{6} & 1/\sqrt{3} & 1/\sqrt{2} \end{pmatrix}$$

various **other** ansatze were proposed: mu-tau, bi-maximal, tri-maximal, tetra-maximal, symmetric mixing, hexagon mixing, golden mixing, QLC ...



Babu et al PLB552 (2003) 207
 Hirsch et al PRD69 (2004) 093006



ORIGIN OF TBM



Z_3

$\langle \phi \rangle \sim (1, 1, 1)$

Sectors are separated by an extra Abelian Z_n

Z_2

$\langle \phi' \rangle \sim (1, 0, 0)$

CHARGED LEPTONS

NEUTRINOS

$U_{lep}^\dagger U_\nu = V_{TBM}$

FLAVOR SYMMETRIES

Ishimori, Kobayashi, Ohki, Okada, Shimizu & Tanimoto arXiv:1003.3552

Frampton and Kephart, PRD64 (01)

order	groups
6	$S_3 \equiv D_3$
8	$D_4, Q = Q_4$
10	D_5
12	$D_6, Q_6, T \equiv A_4$
14	D_7
16	$D_8, Q_8, Z_2 \times D_4, Z_2 \times Q$
18	$D_9, Z_3 \times D_3$
20	D_{10}, Q_{10}
22	D_{11}
24	$D_{12}, Q_{12}, Z_2 \times D_6, Z_2 \times Q_6, Z_2 \times T, Z_3 \times D_4, Z_3 \times Q, Z_4 \times D_3, S_4$
26	D_{13}
28	D_{14}, Q_{14}
30	$D_{15}, D_5 \times Z_3, D_3 \times Z_5$

A4

Babu, Ma, Valle PLB552 (2003)
 Altarelli, Feruglio NPB72 (2005)
 Hirsch, Morisi, Valle PRD78 & D79
 (2008) & PLB679 (2009) 454
 Hagedorn, Molinaro, Petcov (2009)
 Ibanez, Morisi, Valle, PRD80 (2009)

S3

Grimus, Lavoura, JHEP0904
 Mohapatra, Nasri, Yu, PLB627
 Mondragón, Mondragón, Peinado

S4

Lam PRL101
 Bazzocchi, Morisi, PRD80

T'

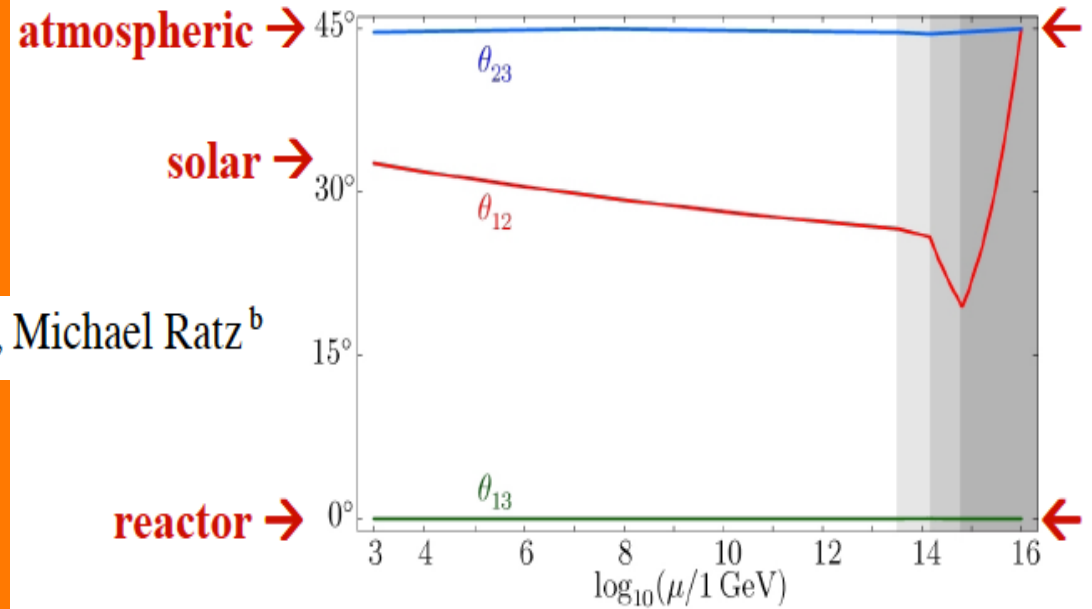
Feruglio, Hagedorn, Lin, Merlo (2007)
 Carr, Frampton (2007)
 Aranda, Carone, Lebed PLB474

D(27)

Medeiros, King, Ross PLB648

TBM DEVIATIONS

Stefan Antusch ^a, Jörn Kersten ^a, Manfred Lindner ^a, Michael Ratz ^b

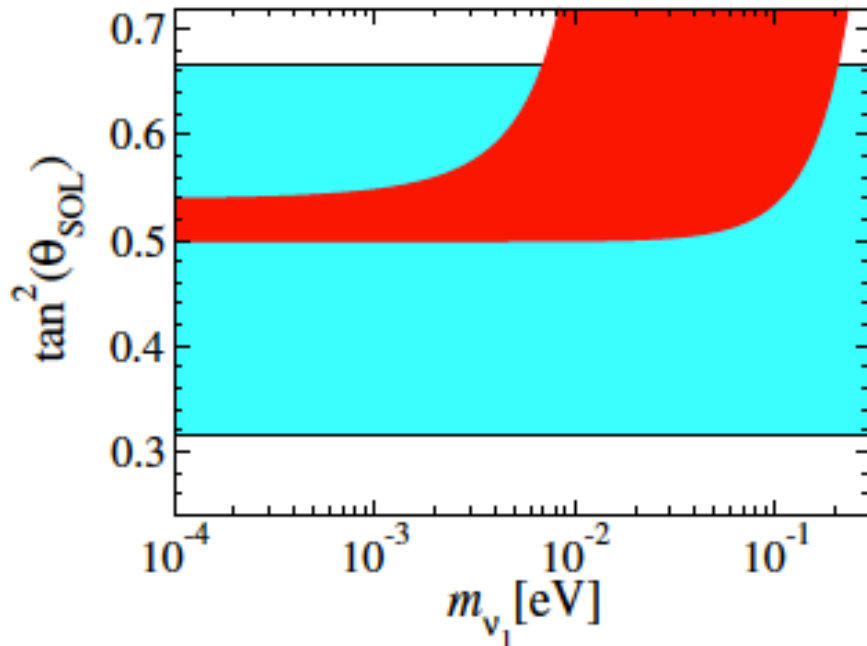


Hirsch et al

PHYSICAL REVIEW D 75, 053006 (2007)

Luo, Xing, Plentinger, Rodejohann, ...

Minimal supergravity radiative effects on the tribimaximal neutrino mixing pattern



$$\epsilon_{23} = \frac{\delta_{23}^0 (m_2 + m_3)}{(-1 - \delta_{22}^0)m_2 + (1 + \delta_{33}^0)m_3},$$

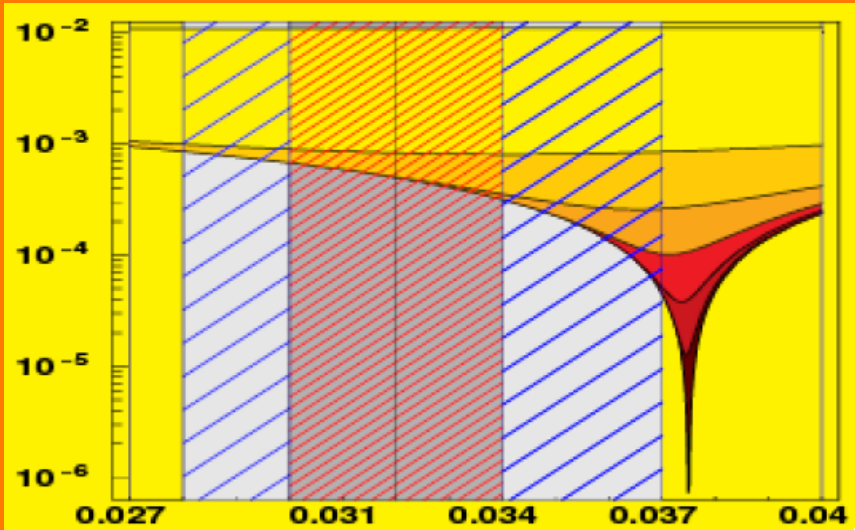
$$\epsilon_{13} = \frac{\delta_{13}^0 (m_1 + m_3)}{(-1 - \delta_{11}^0)m_1 + (1 + \delta_{33}^0)m_3},$$

$$\epsilon_{12} = \frac{\delta_{12}^0 (m_1 + m_2)}{(-1 - \delta_{11}^0)m_1 + (1 + \delta_{22}^0)m_2}.$$

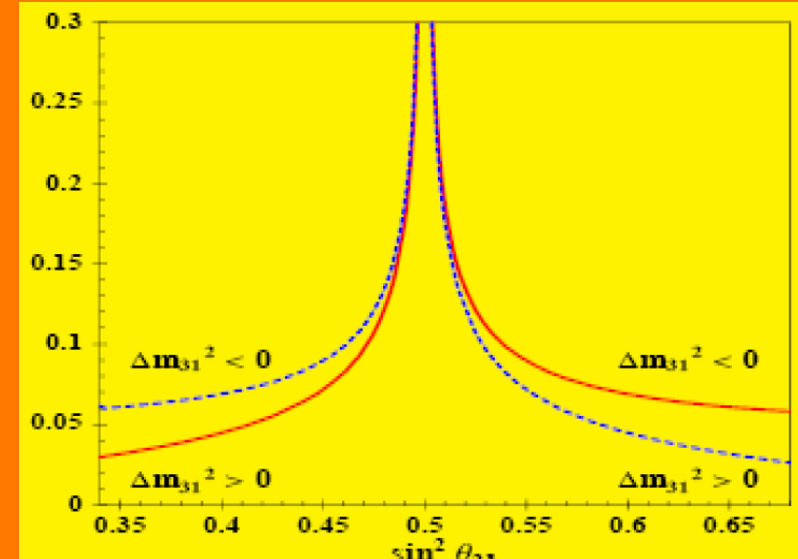
$\theta = \mu$ DBD & FLAVOR

PRL 99 (2007) 151802, PRD82 (2010) 073008

PRD78:093007 (2008)



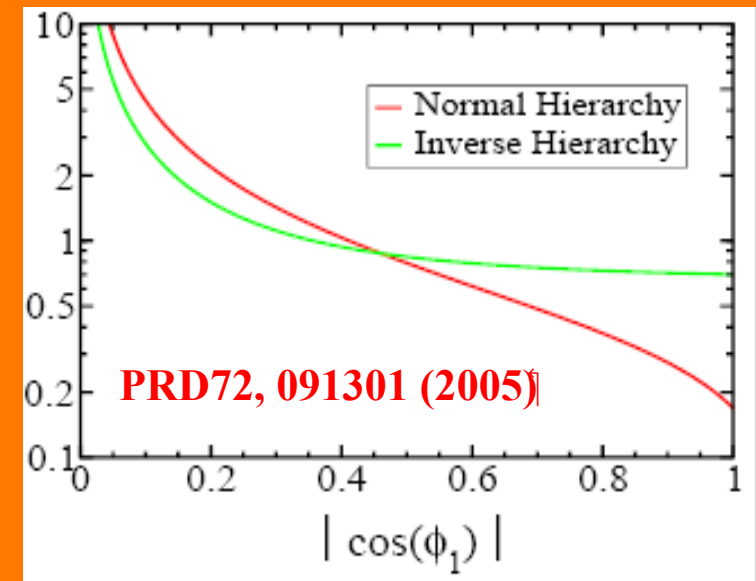
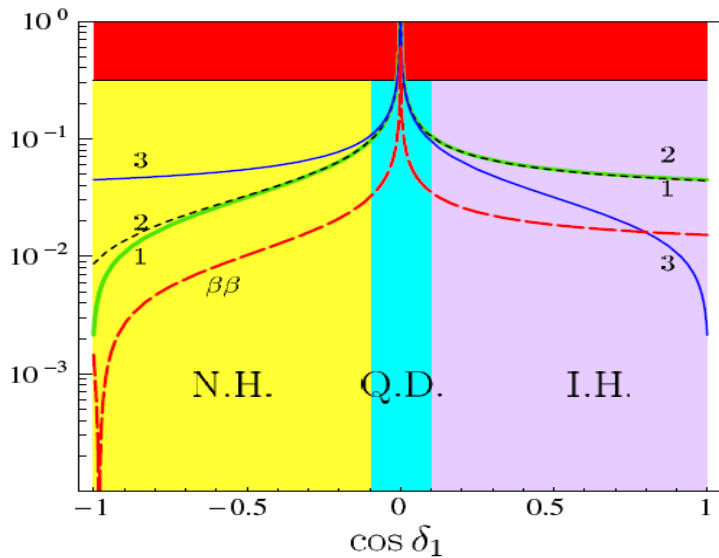
correlates with $\alpha = \frac{\Delta m_{\text{SOL}}^2}{\Delta m_{\text{ATM}}^2}$



correlates with ATM angle

A4

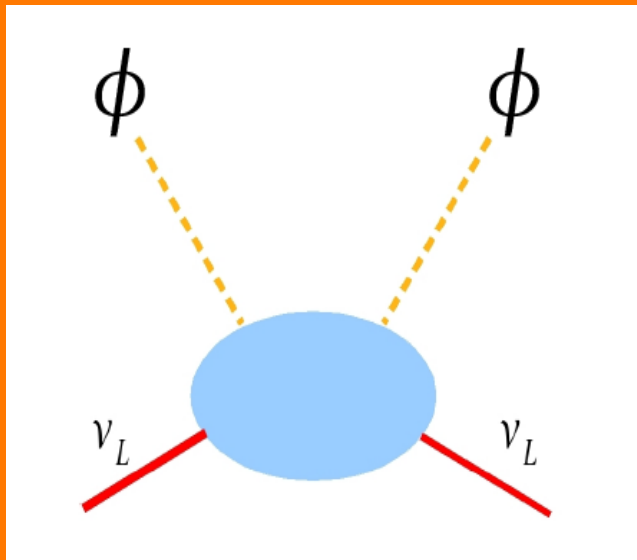
PHYSICAL REVIEW D 79, 016001 (2009)



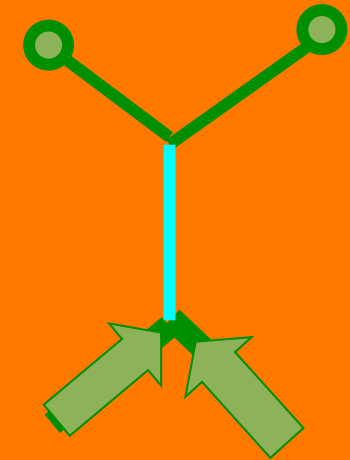
PRD72, 091301 (2005)

correlates with Majorana phase

ORIGIN OF NU-MASSES & MIXINGS



Fermion exchange
Type I & III



Scalar exchange
type II

Schechter-Valle 80/82

scale

mechanism

flavor structure

Minkowski 77

Gellman Ramond Slansky 80

Glashow, Yanagida 79

Mohapatra Senjanovic 80

Schechter-Valle, 80 & 82

Lazarides Shafi Weterrich 81

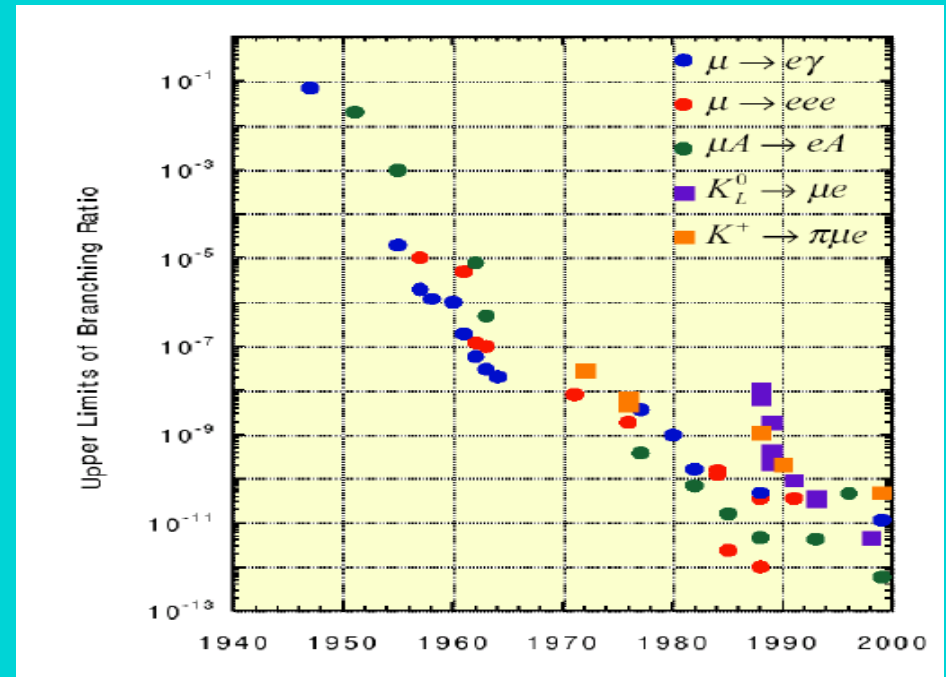
Foot et al 89

LOW-SCALE SEESAW Mohapatra-Valle 86

LFV & neutrino oscillations

Flavor is violated in neutrino
Propagation !!

Okada@NuFact2010



- **HEAVY** vs **LIGHT MESSENGERS**
- **VIRTUAL** vs **DIRECT EFFECTS**

F. del Aguila et al. Eur.Phys.J.C57:183-308,2008
Hall, Kostelecky, Raby 86, Borzumati, Masiero, 86, ..

LOW-SCALE SEESAW

INVERSE SEESAW

Mohapatra-Valle, 86
 Ibanez Morisi JV, PRD80 (2009) 053015
 Bazzocchi, et al, PRD81 (2010) 051701

LINEAR SEESAW

Malinsky et al PRL95(2005)161801
 Hirsch, et al PLB679:454,2009

- **LFV & CPV survive in massless neutrino limit**
- **hence unsuppressed by m-nu**

Bernabeu et al 87, Branco et al 89, Rius JV 90,
 Gonzalez-Garcia, JV, Mod.Phys.Lett.A7:477,1992
 Ilakovac Kniehl Pilaftsis 95,
 Deppish et al PRD72:036001,2005 & NPB752 (2006) 80
 Deppish, Kosmas & JV 2006,
 Malinsky, Ohlsson, Zhang, PRD79, 073009
 Gavela, Hambye, Hernández, Hernández...
 Okada, et al

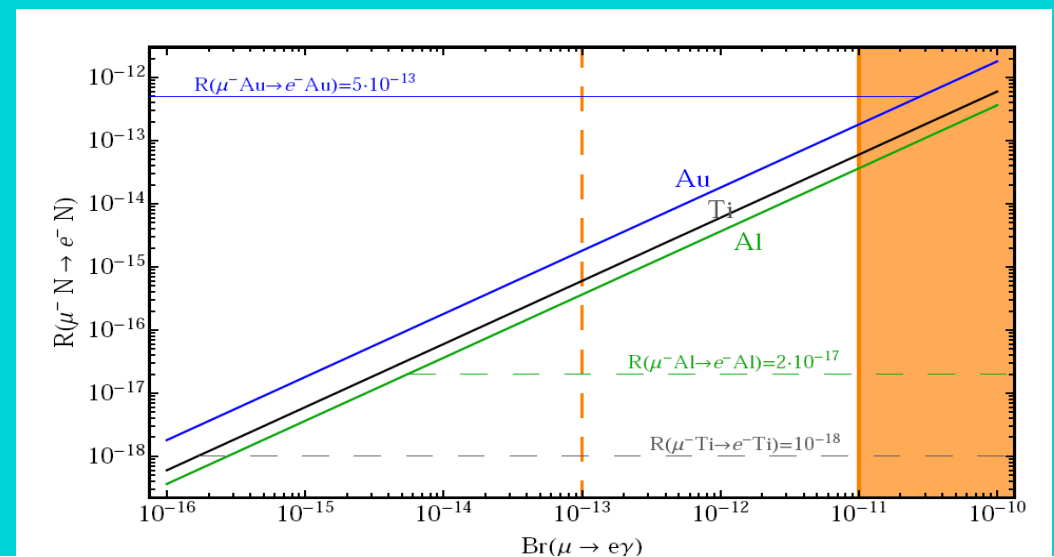
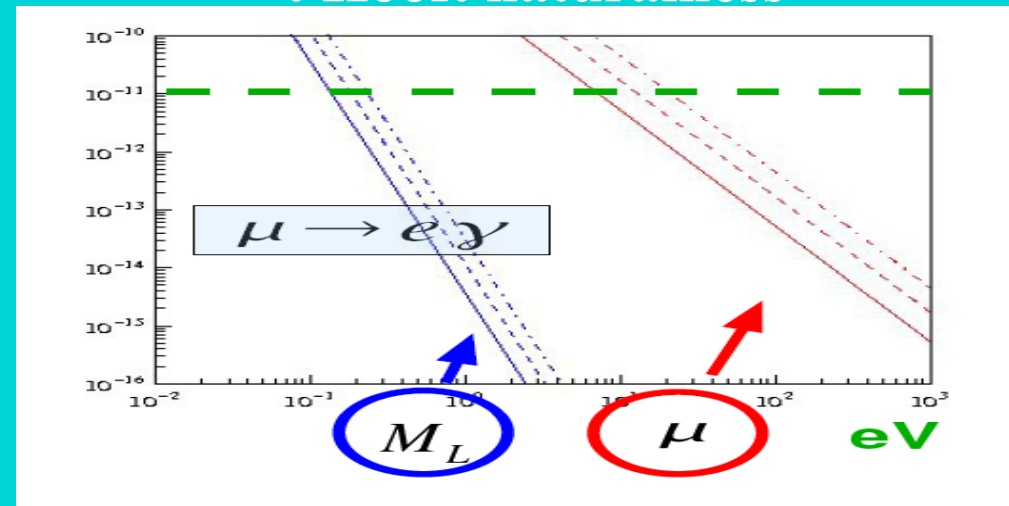
LFV from NHIL exchange

M=0.2-1 TeV

$$\begin{pmatrix} 0 & M_D & M_L \\ M_D^T & 0 & M \\ M_L^T & M^T & 0 \end{pmatrix}$$

$$\begin{pmatrix} 0 & M_D & 0 \\ M_D^T & 0 & M \\ 0 & M^T & \mu \end{pmatrix}$$

t'Hooft naturalness



PROBING LFV IN SUSY DECAYS AT LHC

Hirsch et al PRD 78 (2008) 013006
Esteves et al JHEP05 (2009) 3

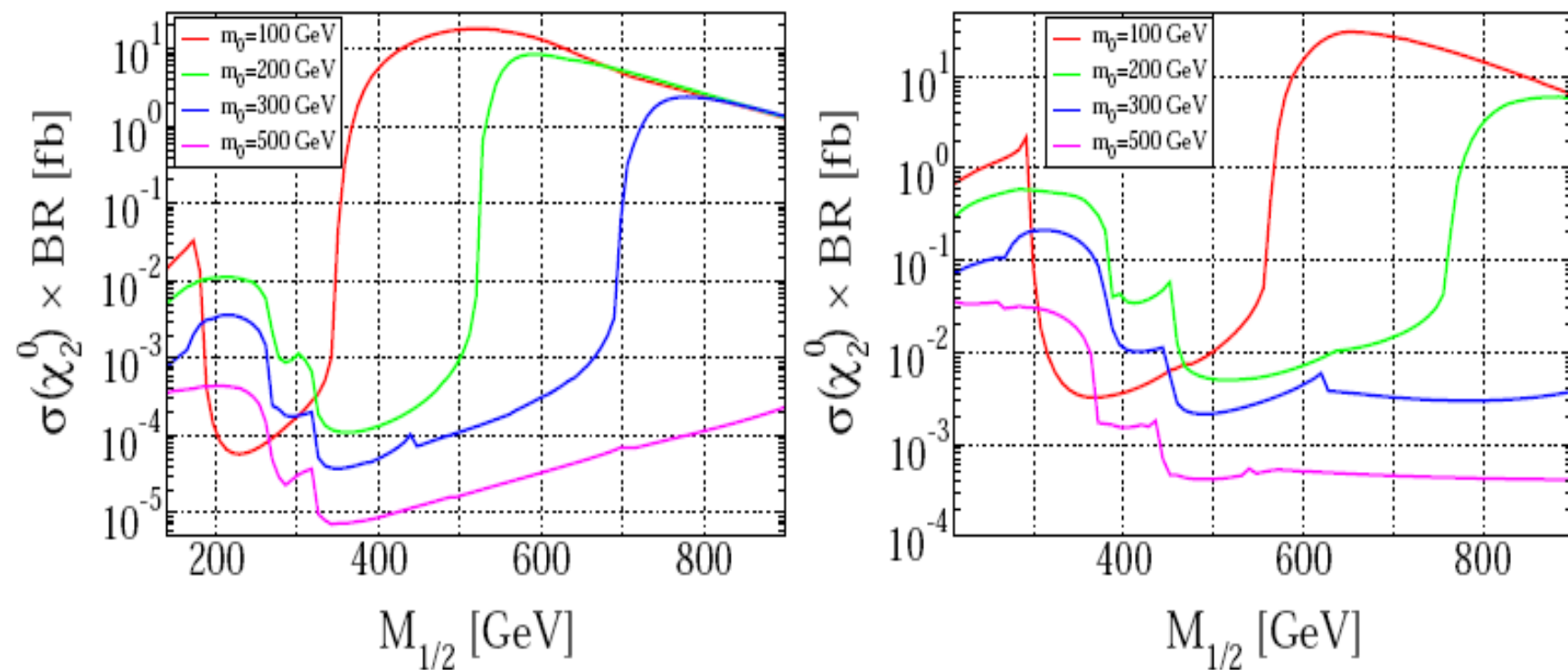
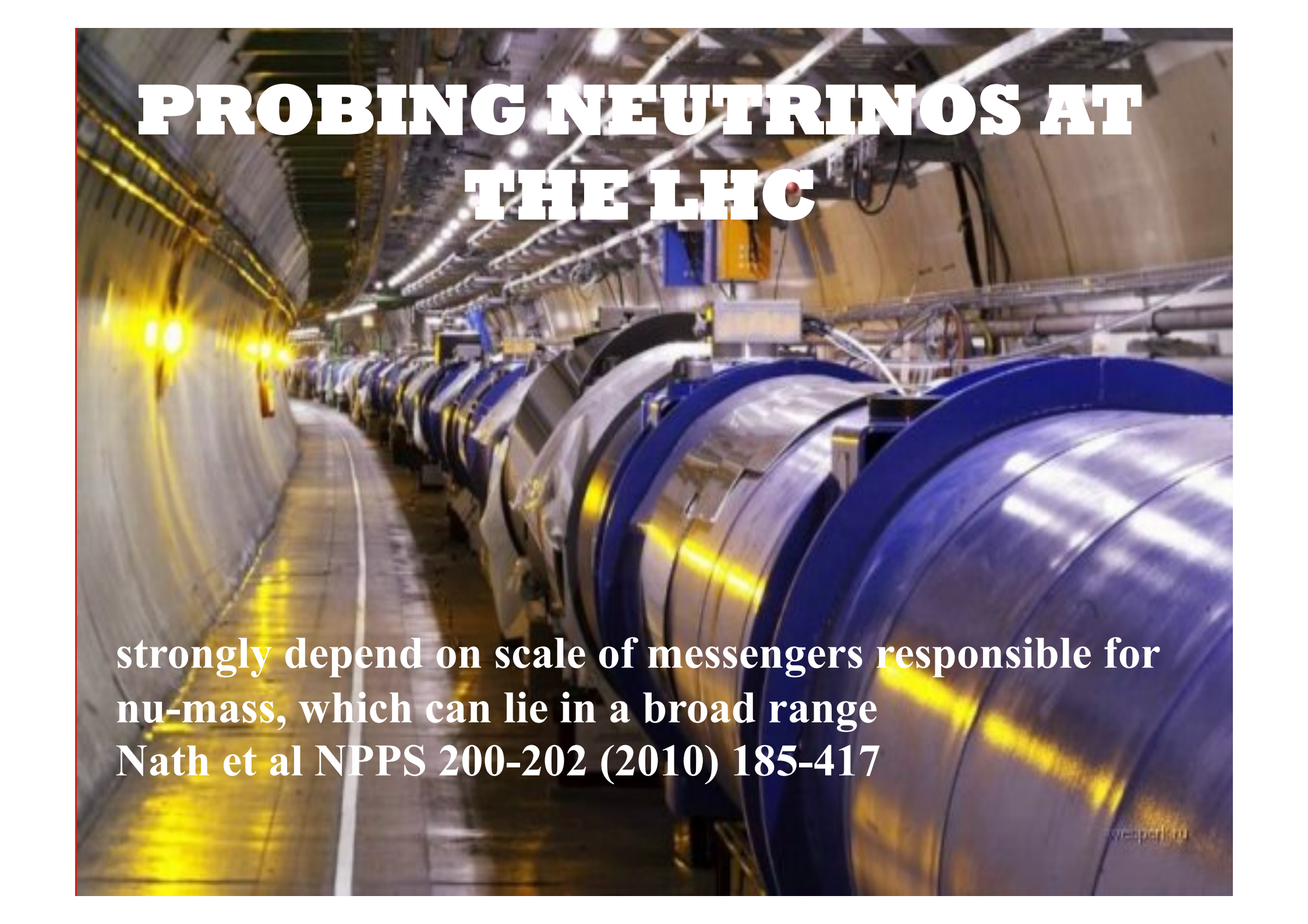


FIG. 12: Production cross section (at leading order) of χ_2^0 times BR of χ_2^0 going to $\mu\text{-}\tau$ lepton pair versus $M_{1/2}$ for $m_0 = 100$ GeV (red), 200 GeV (green), 300 GeV (blue) and 500 GeV (magenta), and for our standard choice of parameters: $\mu > 0$, $\tan\beta = 10$ and $A_0 = 0$ GeV, for type-I (left panel) and for type-II seesaw (right panel) with $\lambda_1 = 0.02$ and $\lambda_2 = 0.5$, imposing $\text{Br}(\mu \rightarrow e + \gamma) \leq 1.2 \cdot 10^{-11}$.



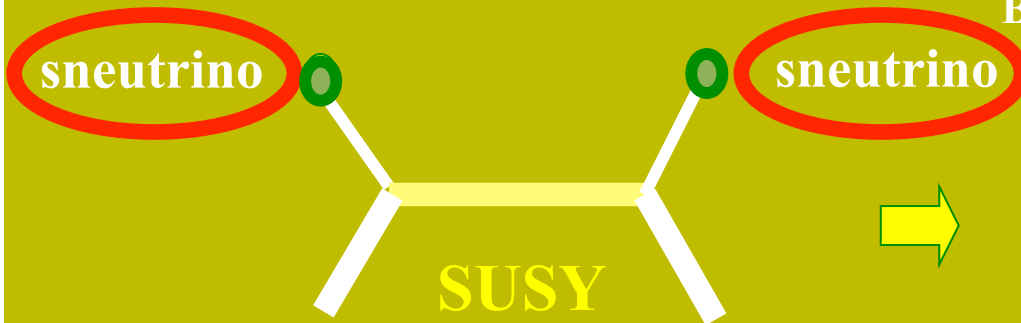
PROBING NEUTRINOS AT THE LHC

**strongly depend on scale of messengers responsible for
nu-mass, which can lie in a broad range
Nath et al NPPS 200-202 (2010) 185-417**

SPONTANEOUS/**BILINEAR** R-PARITY VIOLATION

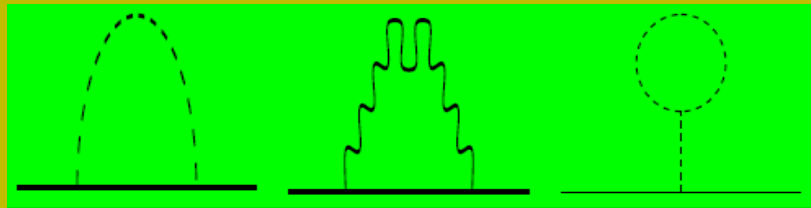
Masiero & Valle, PLB251 (1990) 273

Bhattacharyya & Pal, PRD82 (2010) 055013



ATM SCALE
SUSY-SEESAW

Hall & Suzuki, Ross & JV 85,
Ellis et al, 85, Santamaria JV, ...



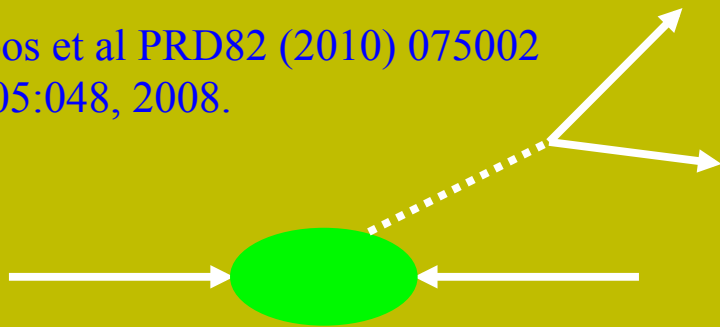
SOLAR SCALE
RADIATIVE

Diaz et al PRD68 (2003) 013009, PRD62 (2000) 113008
PRD65 (2002) 119901; PRD61 (2000) 071703

LSP decay

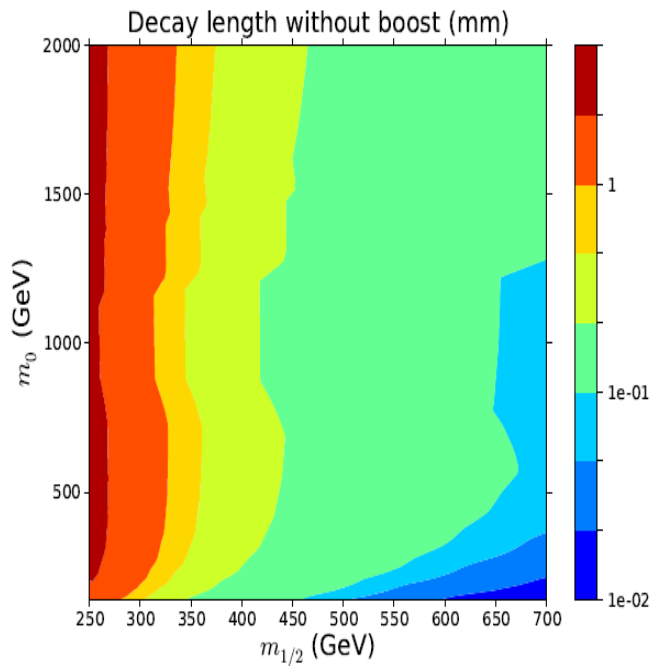
DISPLACED VERTICES

De Campos et al PRD82 (2010) 075002
 JHEP 0805:048, 2008.

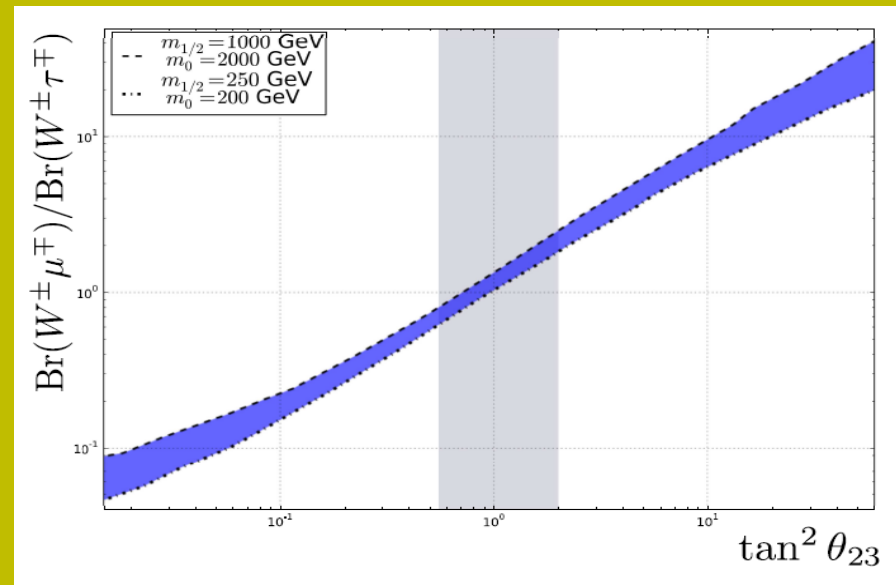


The LSP can live long enough to leave a **displaced vertex** in the detector

χ_{10}^0	$\rightarrow \nu l^+ l^-$
χ_{11}^0	$\rightarrow \nu \tau^+ \tau^-$
χ_{11}^0	$\rightarrow \nu \tau l$
χ_{11}^0	$\rightarrow \nu q \bar{q}$
χ_{10}^0	$\rightarrow \tau q \bar{q}$
χ_{10}^0	$\rightarrow l q \bar{q}$
χ_{11}^0	$\rightarrow \nu b \bar{b}$
χ_{11}^0	$\rightarrow \nu \nu \nu$



$\tilde{\chi}_1^0$ decay length in the plane $m_0, m_{1/2}$ for $A_0 = -100$ GeV, $\tan \beta = 10$ and $\mu > 0$.



$\frac{BR(\chi \rightarrow \mu W)}{BR(\chi \rightarrow \tau W)}$ vs \tan_{atm}^2

PROBING OSCILLATION PARAMETERS @ LHC

- LSP decays deplete pT-miss, increasing multiplicities
- leaving displaced vertices
- decay pattern correlates with oscillation angles

Simulation reveals that 23-mixing angle can be extracted competitive with Super-K

work needed ...

12 & 13-angles, magnitude of nu-mass ...

ROBUST:

similar features in any SUSY breaking scenario & hence with any LSP profile

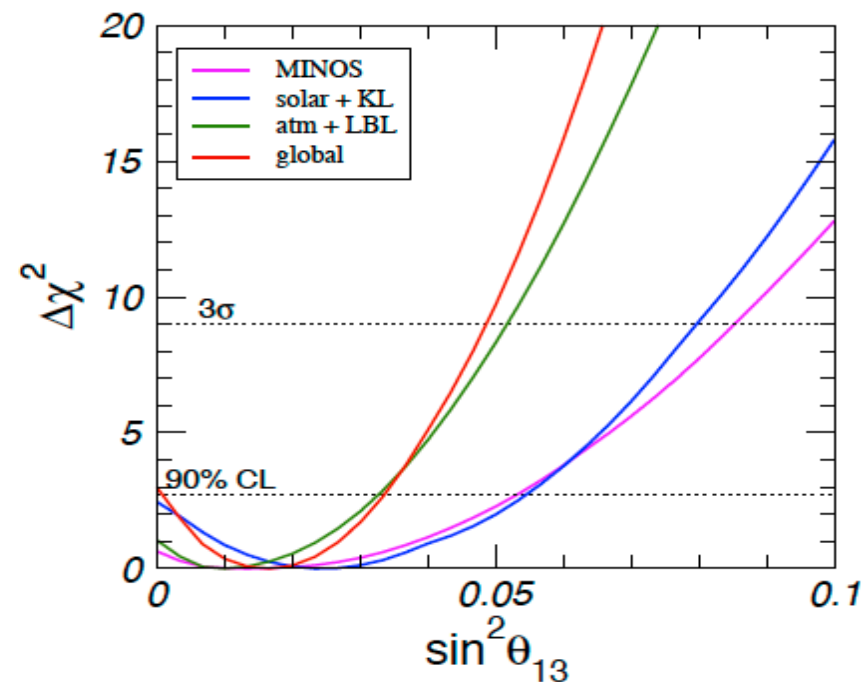
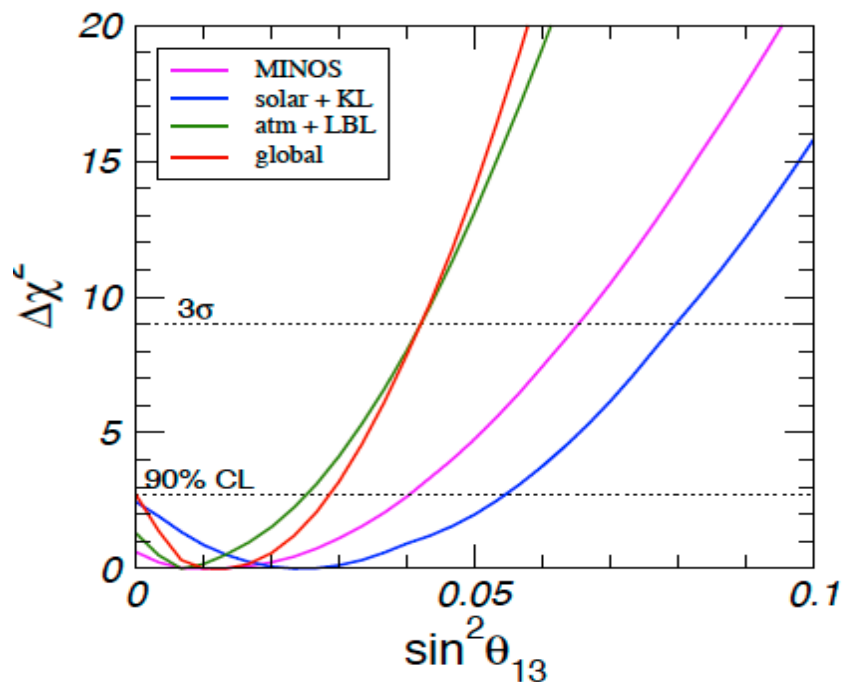
stop Restrepo et al, PRD64 (2001) 055011
stau Hirsch et al, PRD66 (2002) 095006
others D68 (2003) 115007

ROBUSTNESS OF OSCILLATIONS

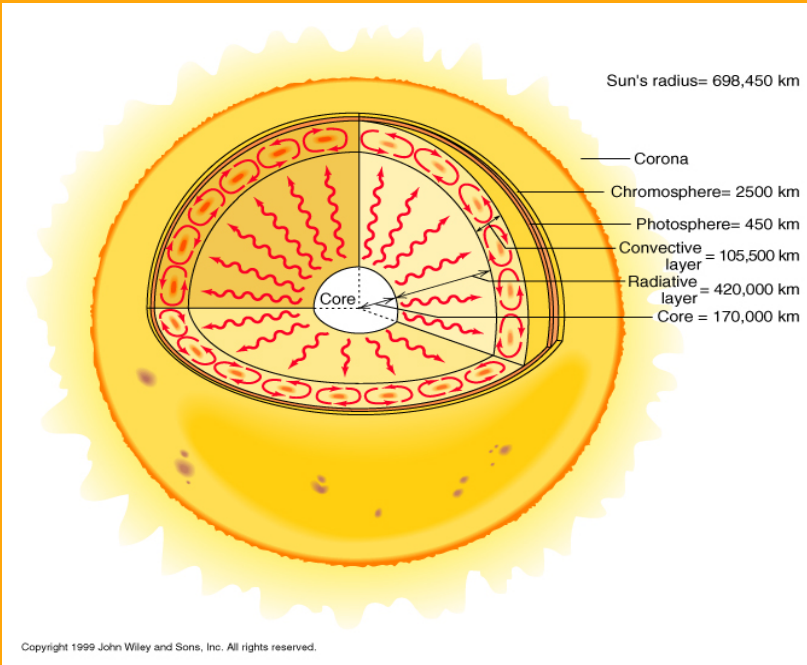
parameter	best fit	2σ	3σ
$\Delta m_{21}^2 [10^{-5} \text{eV}^2]$	$7.65^{+0.22}_{-0.20}$	7.23–8.08	7.05–8.31
$ \Delta m_{31}^2 [10^{-3} \text{eV}^2]$	$2.35^{+0.10}_{-0.09}$	2.17–2.54	2.08–2.64
$\sin^2 \theta_{12}$	$0.315^{+0.020}_{-0.016}$	0.29–0.36	0.27–0.38
$\sin^2 \theta_{23}$	0.52 ± 0.06	0.42–0.61	0.39–0.64
$\sin^2 \theta_{13}$	0.017 ± 0.010	≤ 0.037	≤ 0.049

NEUTRINO OSCILLATIONS STATUS @ Nufact2010

upd Schwetz et al,
NJP 10 (2008) 113011
Maltoni et al, NJP 6 (2004) 122



HOW ROBUST ARE SOLAR NU-OSCILLATIONS



RZ magn fields

Burgess et al JCAP0401 (2004) 007

CZ magn fields

Miranda et al

Both strongly disfavored by KamLAND

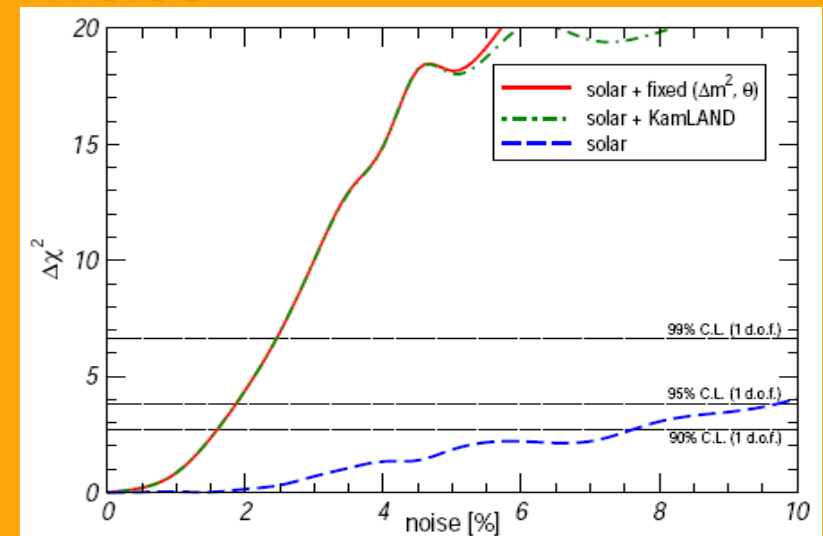
FROM NEUTRINO PROPERTIES TO SOLAR PHYSICS

Test SSM fluxes

Probe RZ mag-fields & RZ-density fluctuations

helioseismology & magneto-gravity waves ...

Burgess et al MNRAS.348 (2004) 609



TRANSITION MAGNETIC MOMENTS

AFFECT NU-PROPAGATION

Miranda et al PRL93 (2004) 051304
PRD70 (2004) 113002

KamLAND anti-nu-e flux limit

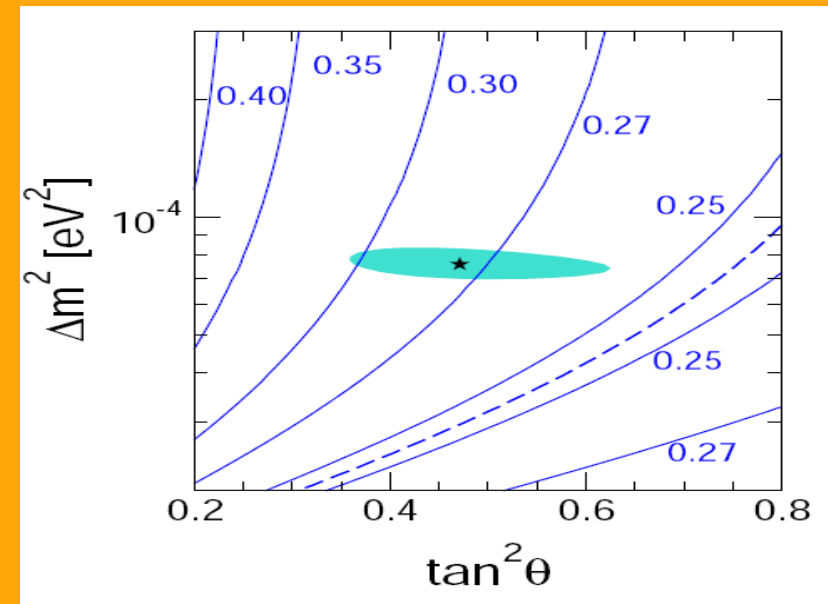
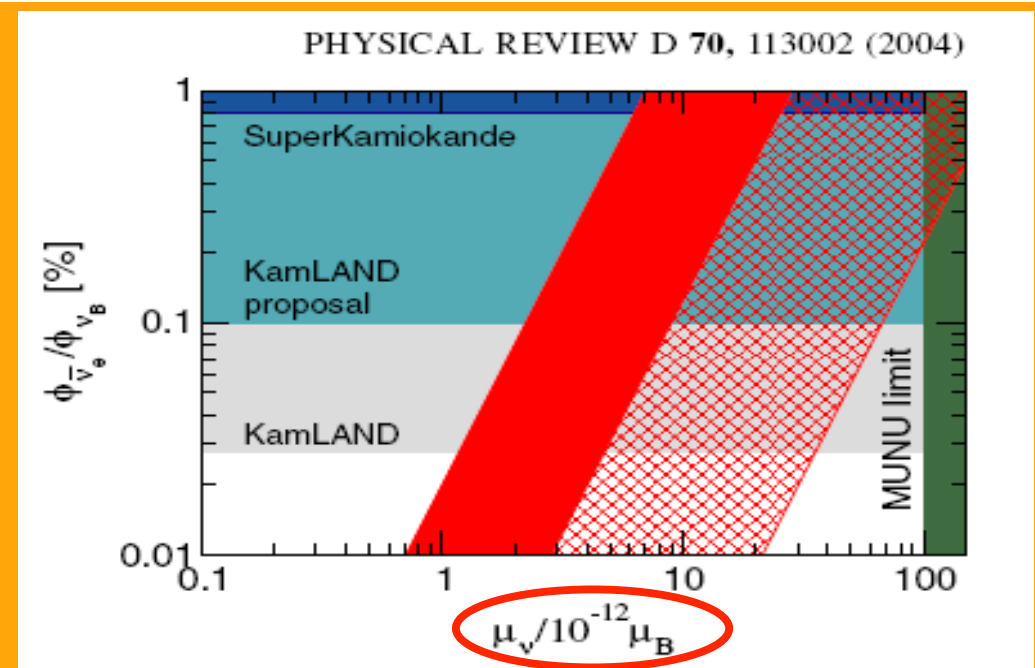
SFP must be sub-leading w.r.t. oscillations

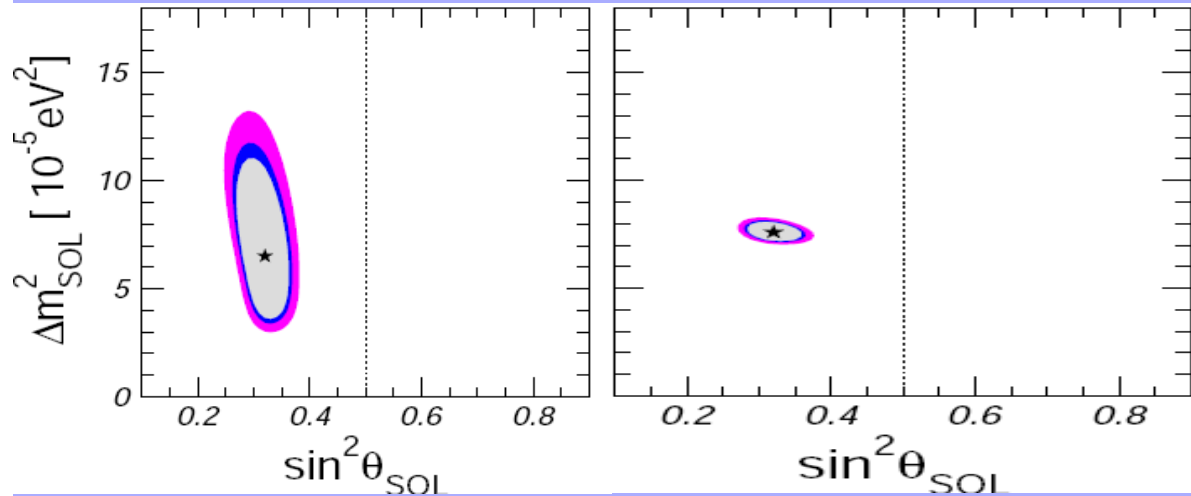
bounds on transition mag-moment in a turbulent magnetic field

AFFECT DETECTION NEUTRINO-ELECTRON SCATTERING X-SECTION

Grimus et al, NPB648, 376 (2003)

cf. Borexino sensitivities





SOL

SOL+KAML

← pure oscillation

OSC+NSI

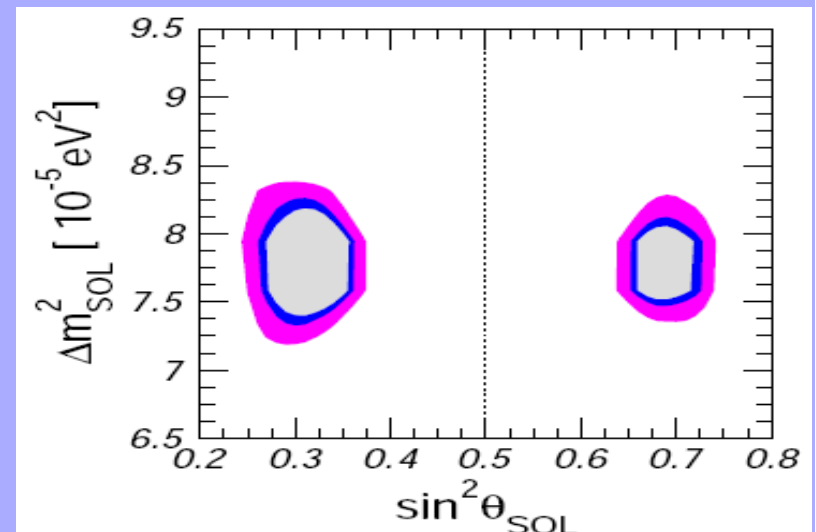
↓ SOL+KAML

The Hamiltonian describing solar neutrino evolution in the presence of NSI contains, in addition to the standard oscillation term,

$$\begin{pmatrix} -\frac{\Delta m^2}{4E} \cos 2\theta + \sqrt{2} G_F N_e & \frac{\Delta m^2}{4E} \sin 2\theta \\ \frac{\Delta m^2}{4E} \sin 2\theta & \frac{\Delta m^2}{4E} \cos 2\theta \end{pmatrix}, \quad (2)$$

a term H_{NSI} accounting for an effective potential induced by the NSI with matter, which may be written as

$$H_{\text{NSI}} = \sqrt{2} G_F N_d \begin{pmatrix} 0 & \varepsilon \\ \varepsilon & \varepsilon' \end{pmatrix}. \quad (3)$$



Combining with accelerator data: tau-neutrino NSI parameters

PHYSICAL REVIEW D **80**, 105009 (2009)

TABLE I. Sensitivity of neutrino experiments to flavor-conserving NSI parameters.

Data	ϵ_{ee}^{dV}	$\epsilon_{\tau\tau}^{dV}$	ϵ_{ee}^{dA}	$\epsilon_{\tau\tau}^{dA}$
Solar propagation	↓	↓		
Solar NC detection			↓	↓
KamLAND propagation	↓	↓		
CHARM detection	↓		↓	

$$\epsilon = -\sin\theta_{23}\epsilon_{e\tau}^{dV}, \quad \epsilon' = \sin^2\theta_{23}\epsilon_{\tau\tau}^{dV} - \epsilon_{ee}^{dV}$$

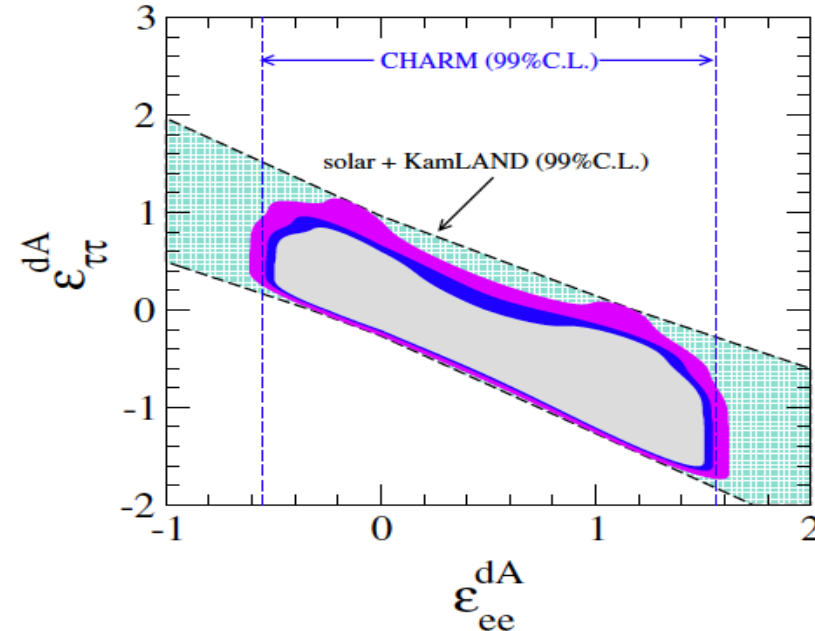
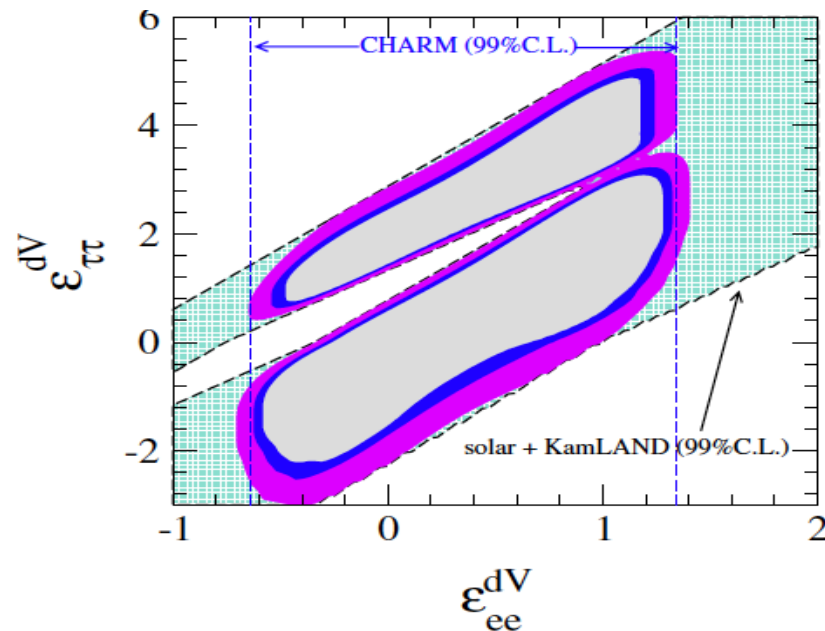


FIG. 6 (color online). Constraints on the vector (left panel) and axial-vector (right panel) NSI couplings from our global analysis at 90%, 95%, and 99% C.L., and from the separate solar + KamLAND and CHARM data sets (dashed lines).

Interplay of nonstandard neutrino interactions with oscillations

Palazzo & JV PRD80:091301,2009

similar confusion as in
LBL neutrino oscillations

Huber, Schwetz, JV PRL88:101804,2002 & PRD66:013006,2002

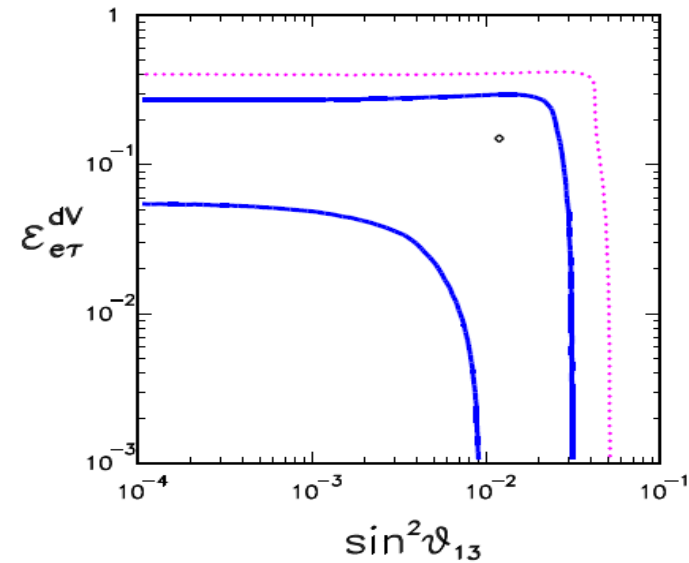


FIG. 2: Region allowed by the combination of solar and KamLAND data at two C.L.'s [$\Delta\chi^2 = 1$ (solid line) and $\Delta\chi^2 = 4$ (dashed line)] after marginalization of δm^2 and θ_{12} .

NSI in MINOS Parke et al, Akhmedov et al

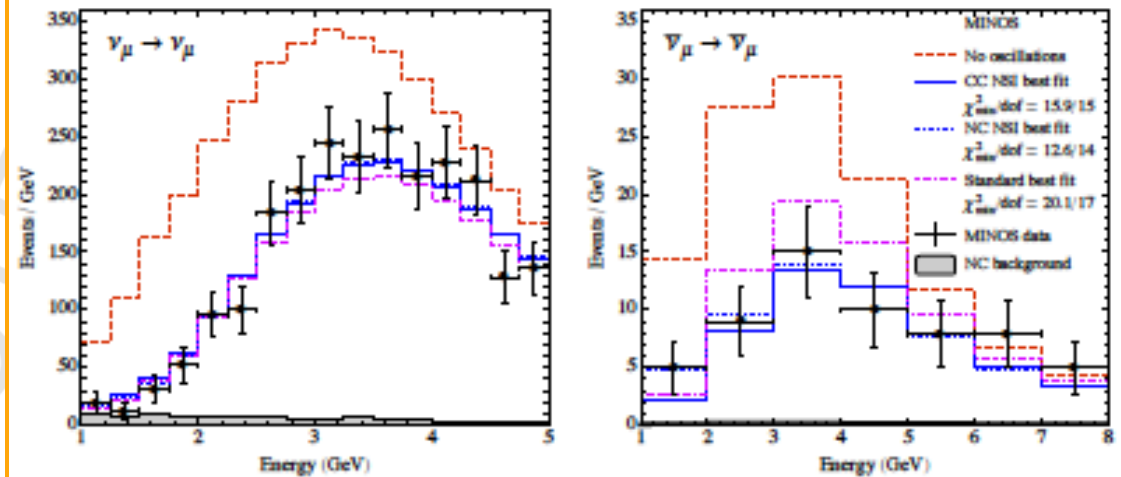
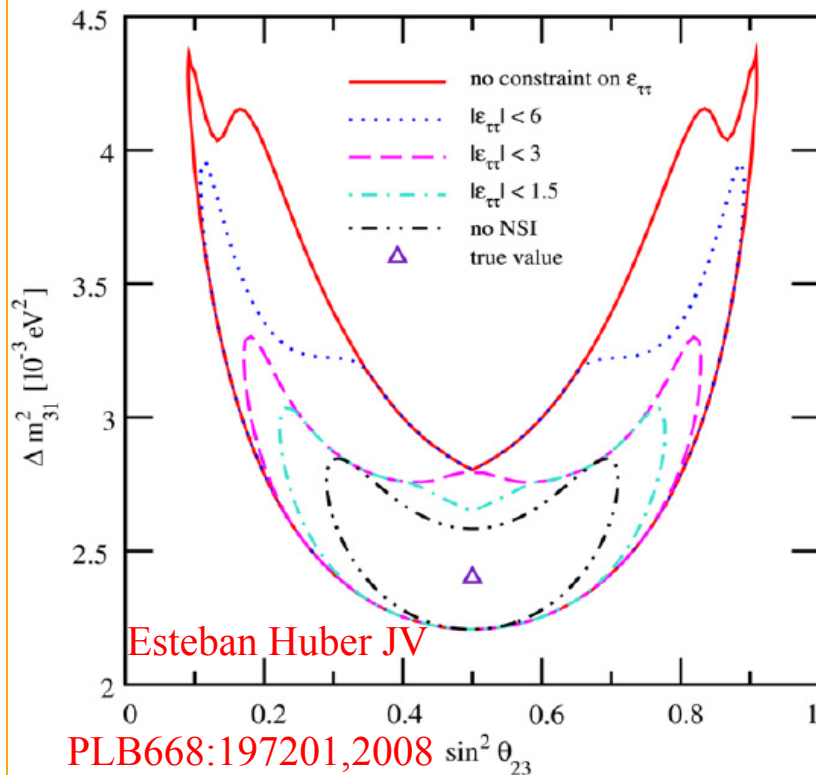
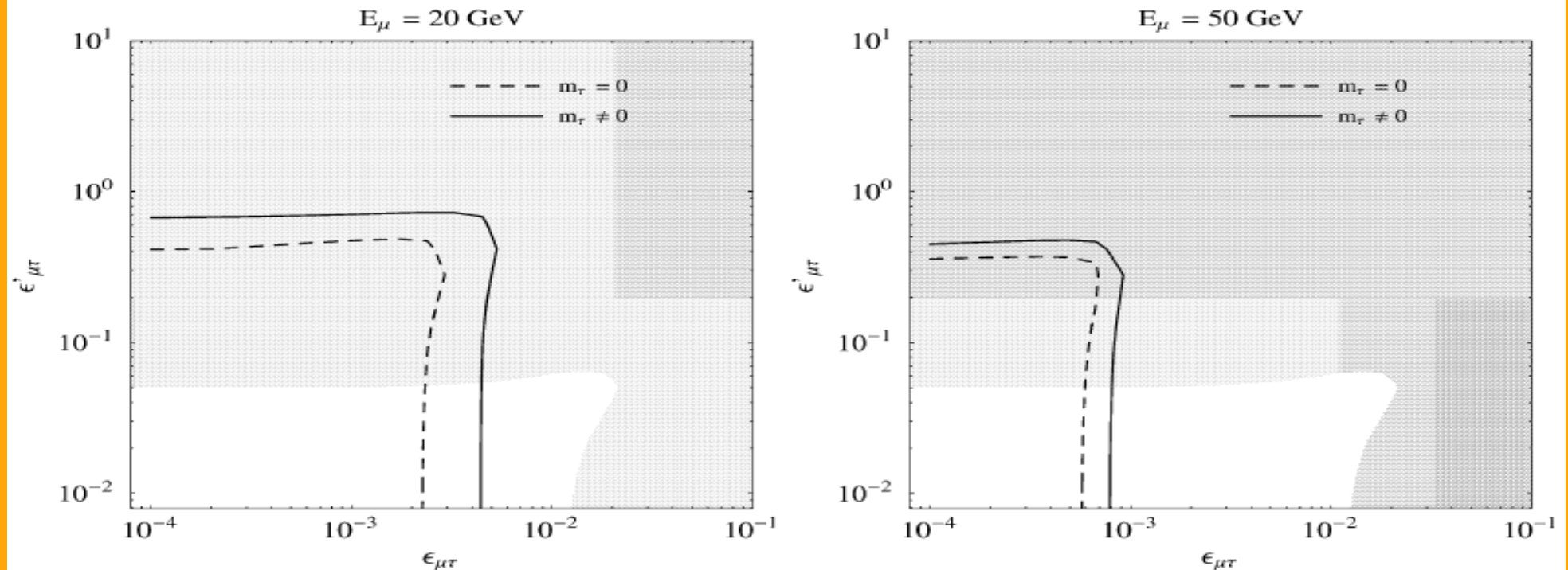


Figure 3: Comparison of MINOS data (black dots and error bars) to theoretical predictions including neutral-current NSI parameterized by the best fit point eq. (20) (blue dotted histograms) and charged current NSI parameterized by the best fit point eq. (21) (blue solid histograms). For comparison, the red dashed histograms show the theoretical prediction in the absence of neutrino oscillations, and the pink dash-dotted histograms represent the results of a two-flavor standard oscillation fit to the combined ν_μ and $\bar{\nu}_\mu$ data.

improving FC NSI sensitivity @ future NuFact

P. Huber, J.W.F. Valle / Physics Letters B 523 (2001) 151–160



PROBING NSI WITH ELECTRONS

current bounds (90% CL):

$$-0.14 < \epsilon_{ee}^{eL} < 0.09$$

$$-0.03 < \epsilon_{ee}^{eR} < 0.18$$



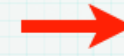
$$-0.036 < \epsilon_{ee}^{eL} < 0.063$$

$$-0.27 < \epsilon_{ee}^{eR} < 0.59$$

current bounds (90% CL):

$$-0.6 < \epsilon_{\tau\tau}^{eL} < 0.4$$

$$-0.4 < \epsilon_{\tau\tau}^{eR} < 0.6$$



$$-0.16 < \epsilon_{\tau\tau}^{eL} < 0.11$$

$$-1.05 < \epsilon_{\tau\tau}^{eR} < 0.31$$

Bolaños et al PRD79 (2009) 113012 improves LL.

Texono 1006.1947 improves ee-RR

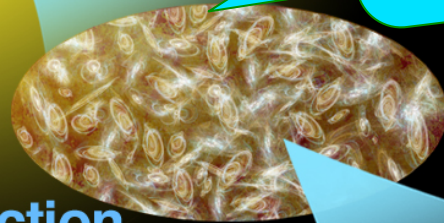
Wrt. Barranco et al PRD77:093014,2008 & PRD73:113001,2006

COSMO-CONNECTION

DAWN
OF
TIME
?

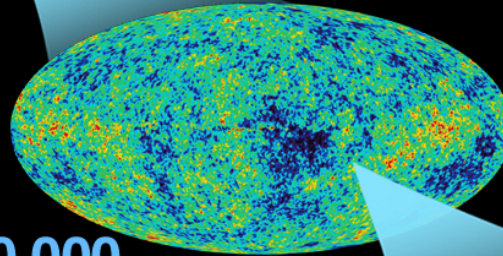
neutrinos may be relevant
here LG & DM

tiny fraction
of a second



inflation

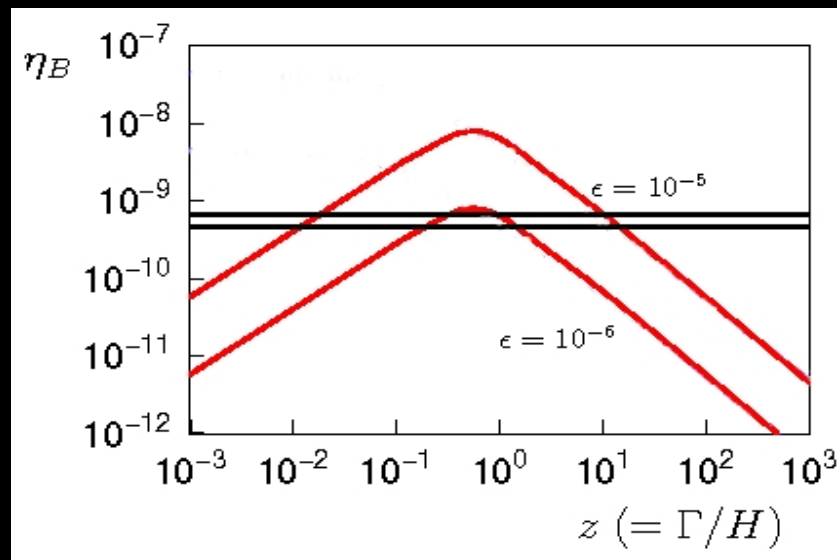
380,000
years



13.7
billion
years



THERMAL SEESAW LEPTOGENESIS



From PRD77 (2008) 055002

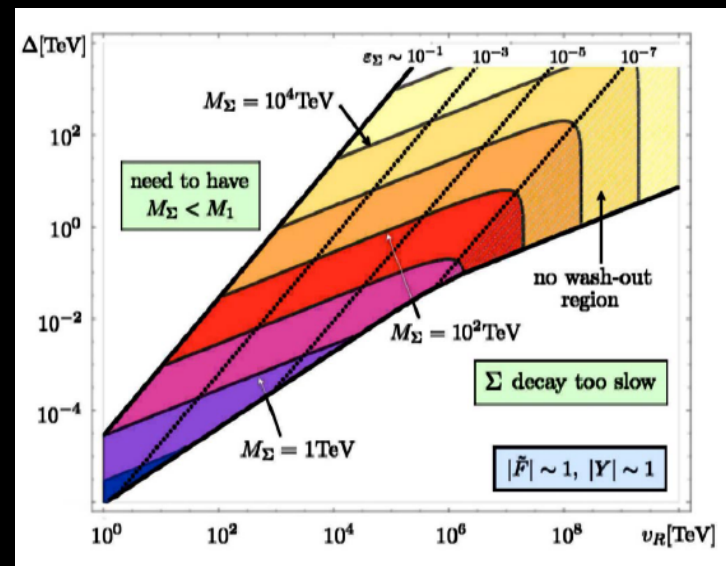
BUT inconsistency with BBN

Kawasaki, Kohri & Moroi, PRD71 (2005) 083502

Gives lower bound on M_1

Sakharov, KRS, Fukugita, Yanagida

Low-scale LG in non-minimal seesaw



Dirac phase suffices

PRL 96, 011601 (2006)

PHYSICAL REVIEW LETTERS

week ending
13 JANUARY 2006

R Parity Violation Assisted Thermal Leptogenesis in the Seesaw Mechanism

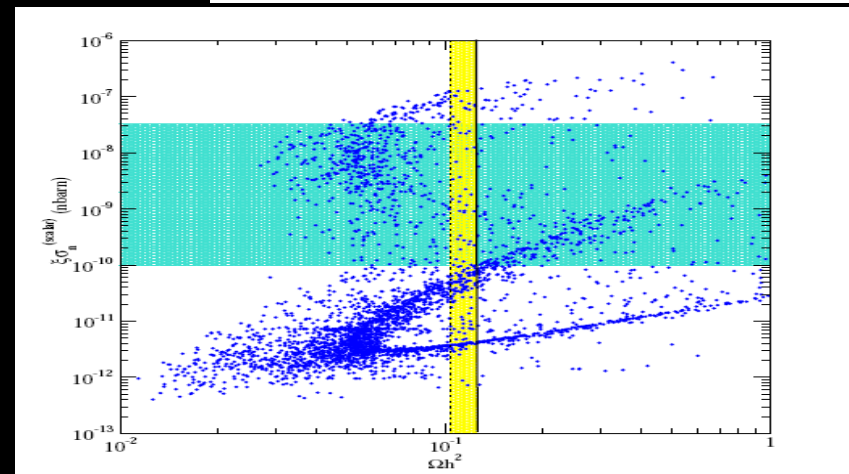
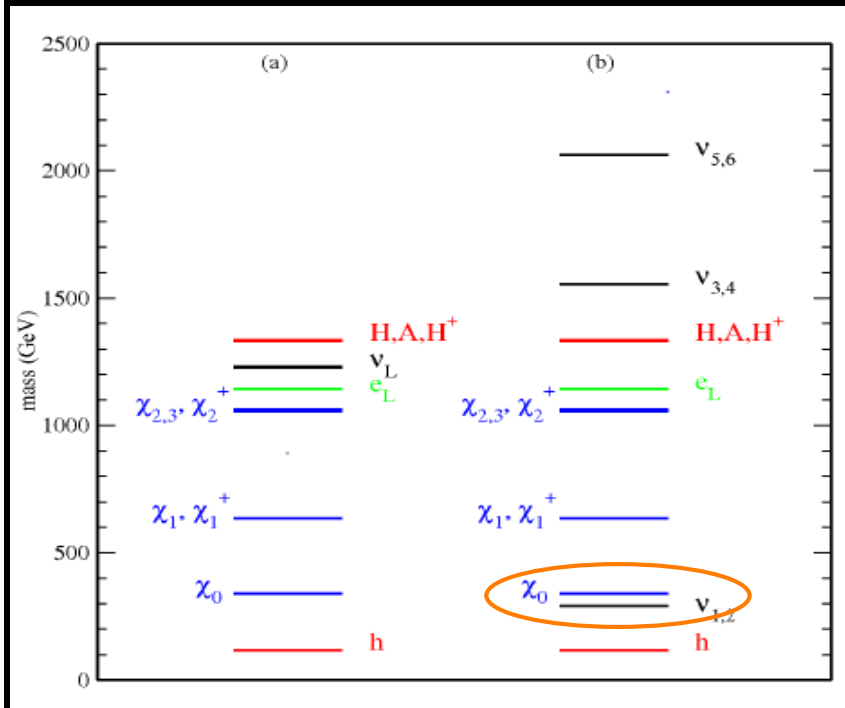
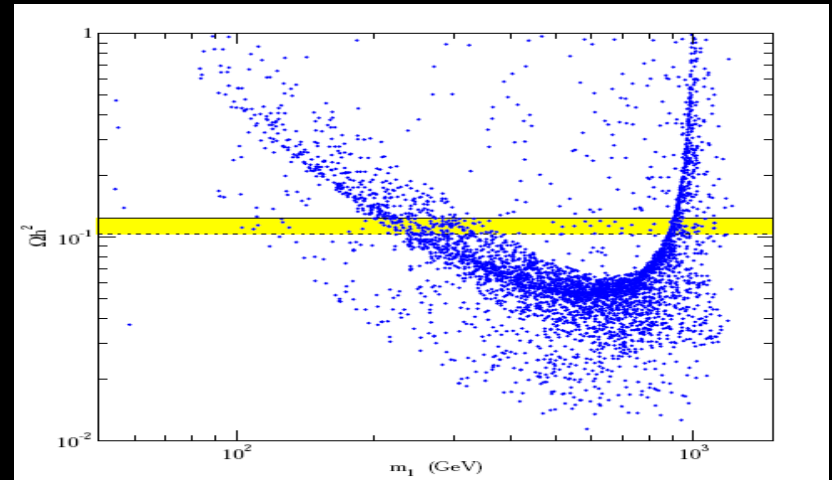
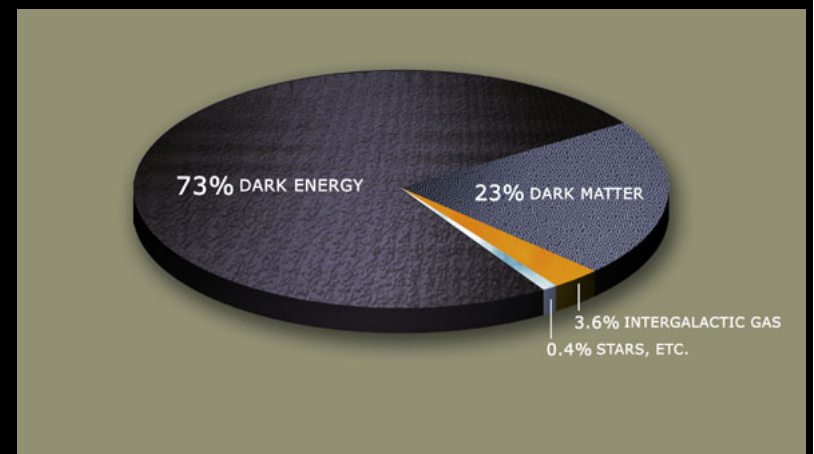
Even if not the source **neutrinos**
may give the clue to DM

Neutrinos masses may change the SUSY spectrum : e.g. in inverse seesaw

one may have SNEUTRINO-like DM

Arina & al PRL101 (2008) 161802

Bazzocchi, Cerdeno, Munoz, Valle, PRD81:051701,2010



Gravity → No DM strictly stable

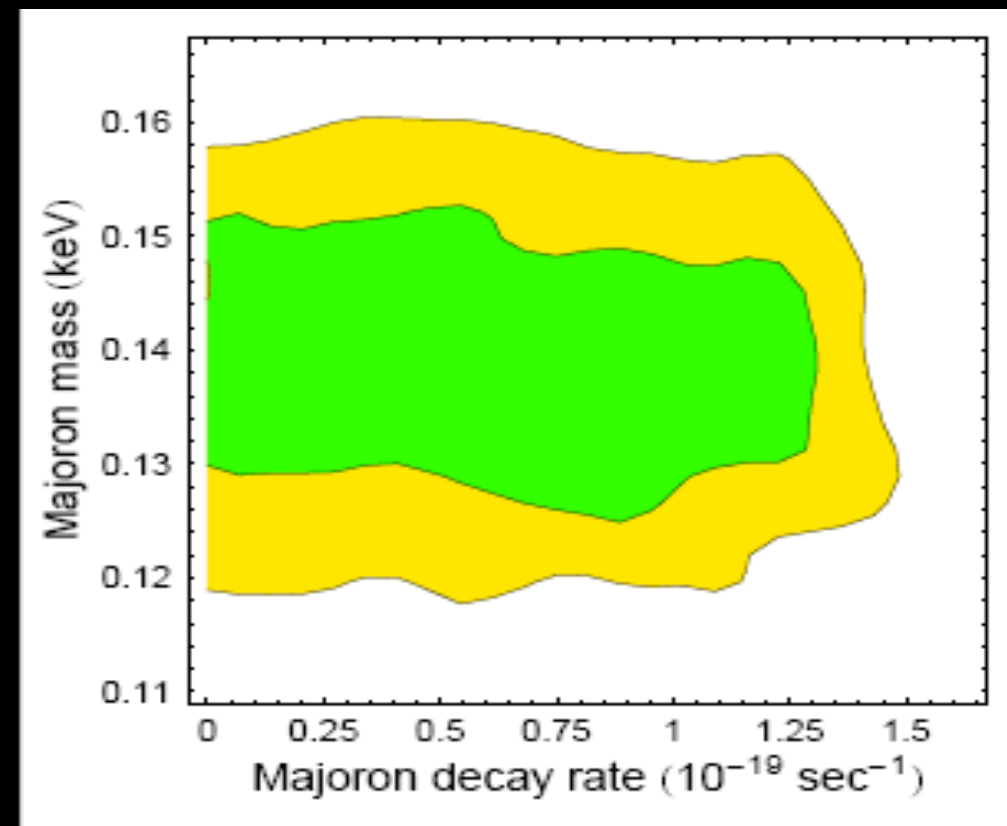
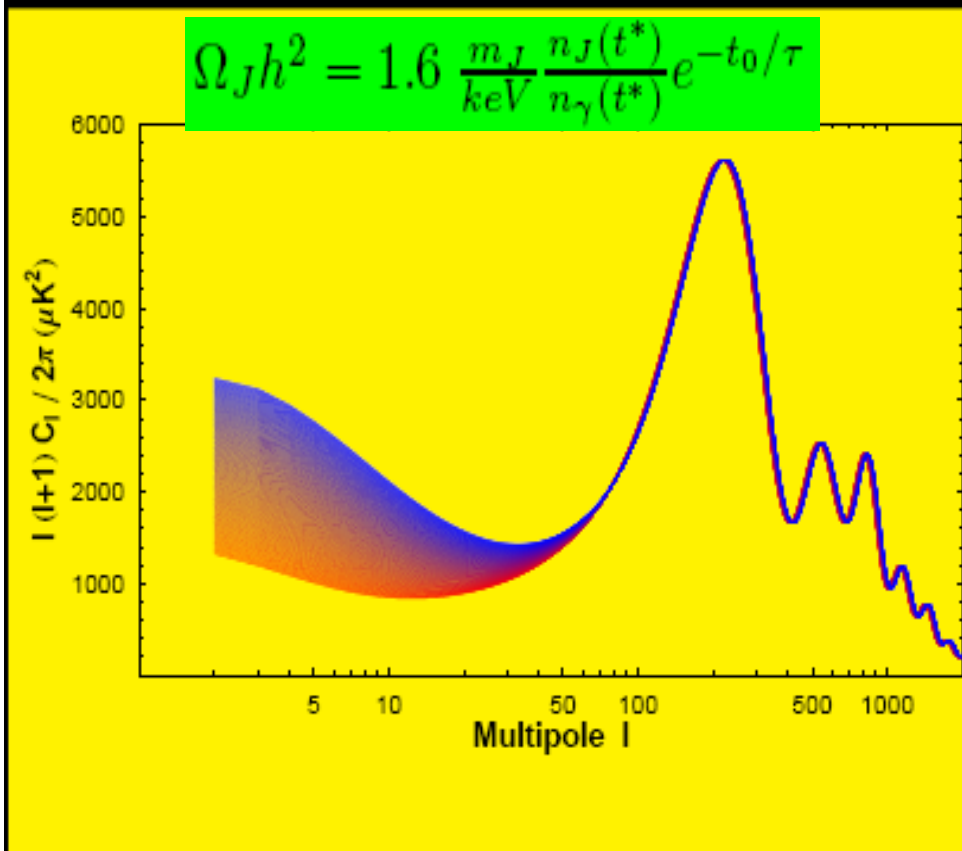
Coleman 88, Kallosh, Linde, Susskind, Nelson, Seiberg, ...

Majoron decaying dark matter

Berezinsky et al PLB318 (1993) 360, PRD57(1998) 147

Consistency with CMB

Lattanzi & Valle, PRL99 (2007) 121301

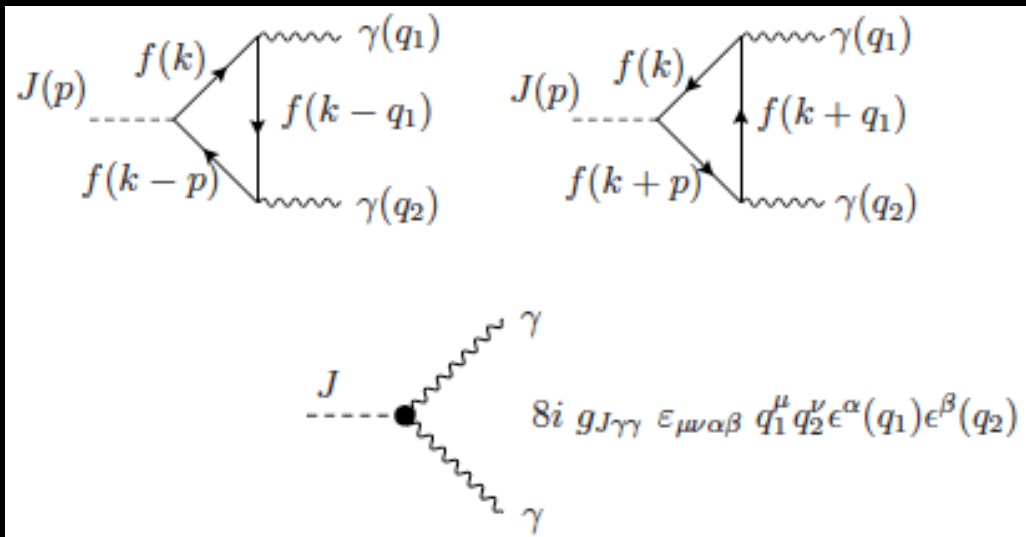


Bazzocchi & al JCAP 0808 (2008) 013

XENIA arXiv:0906.1788

TYPE-II SEESAW MAJORON DECAYING-DM

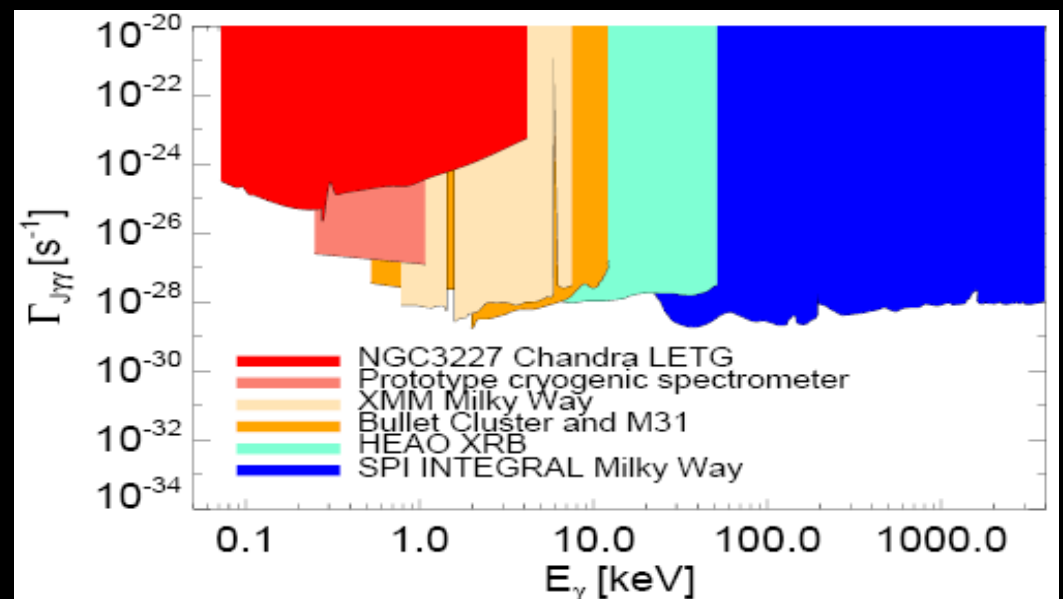
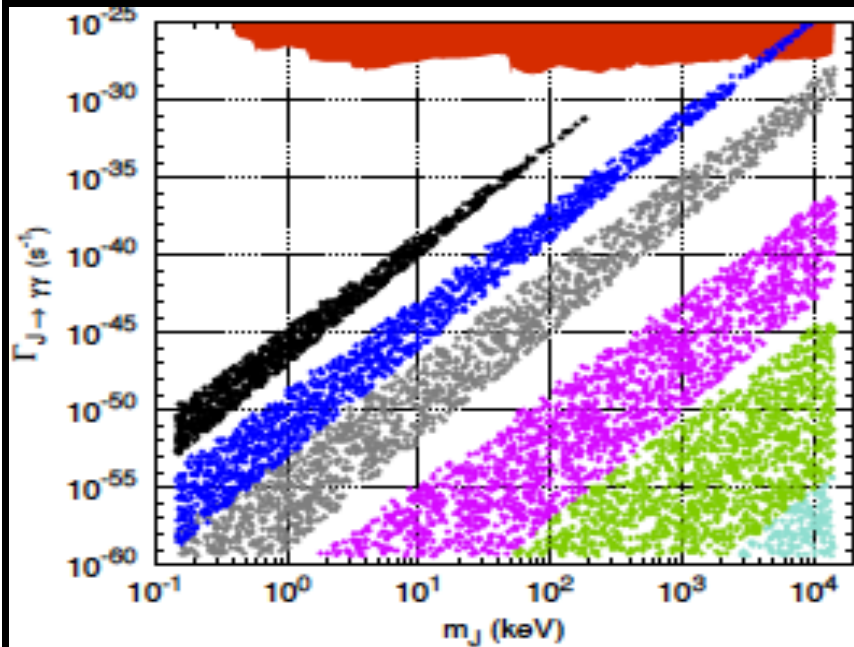
$$g_{J\gamma\gamma} J \epsilon^{\nu\mu\rho\sigma} F_{\nu\mu} F_{\rho\sigma}$$



Esteves et al, PRD 82, 073008 (2010)

TABLE I: Lepton multiplet structure ($Q = T_3 + Y/2$)

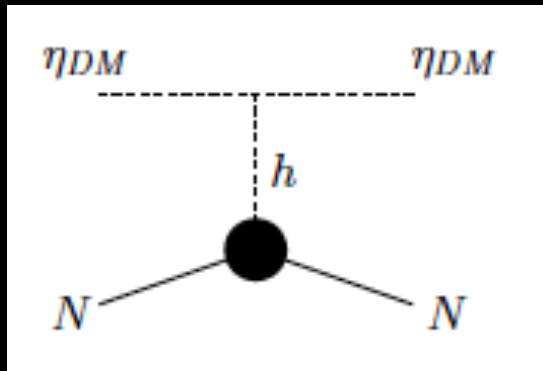
	L_1	L_2	L_3	l_{RH}	ν_{iR}	Φ_i	Δ	σ	S_i
$SU(2)$	2	2	2	1	1	2	3	1	1
$U(1)_Y$	-1	-1	-1	-2	0	-1	2	0	0
A_4	$1'$	1	$1''$	3	3	3	$1''$	$1''$	3
L	1	1	1	1	1	0	-2	-2	1



Discrete dark matter

Non abelian discrete symmetries are motivated by neutrino oscillation data [5, 6]. Here we propose that the same symmetry explaining neutrino mixing angles is also responsible for the dark matter stability. In our simplest type-I seesaw [7] realization the flavor symmetry A_4 spontaneously breaks to Z_2 providing a stable DM

	L_e	L_μ	L_τ	l_e^c	l_μ^c	l_τ^c	N_T	N_4	H	η
$SU(2)$	2	2	2	1	1	1	1	1	2	2
A_4	1	1'	1''	1	1''	1'	3	1	1	3



Inverse hierarchy

$$\theta_{13} = 0$$

Just an example !!

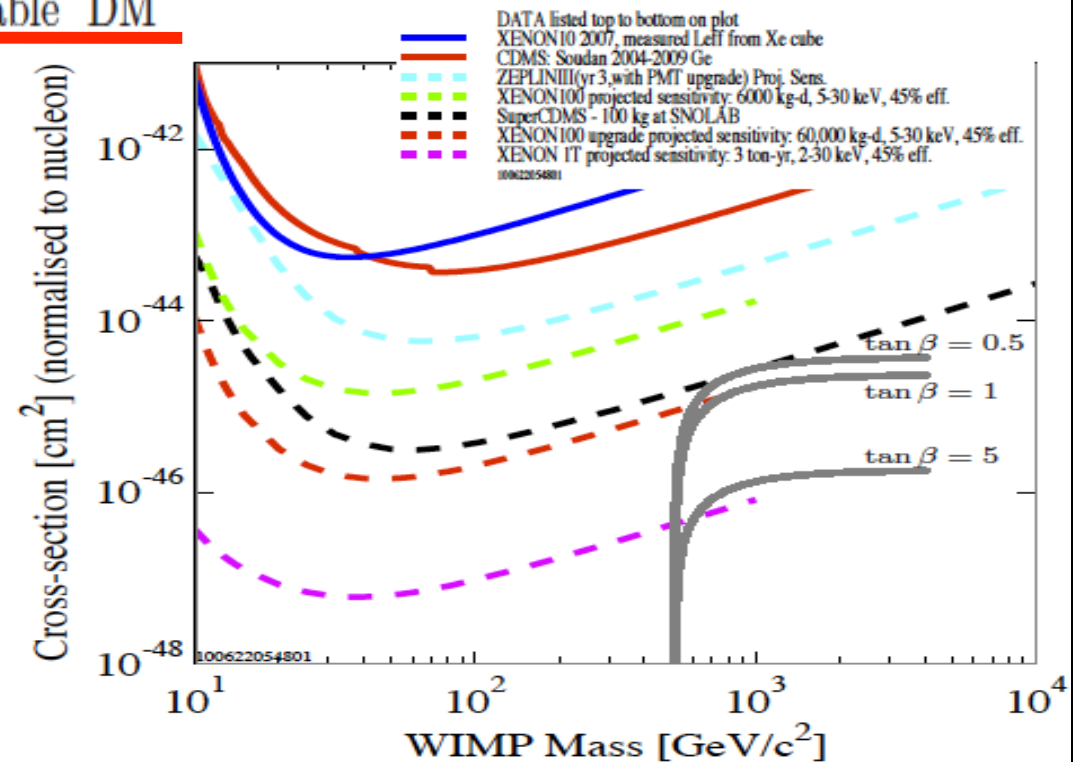
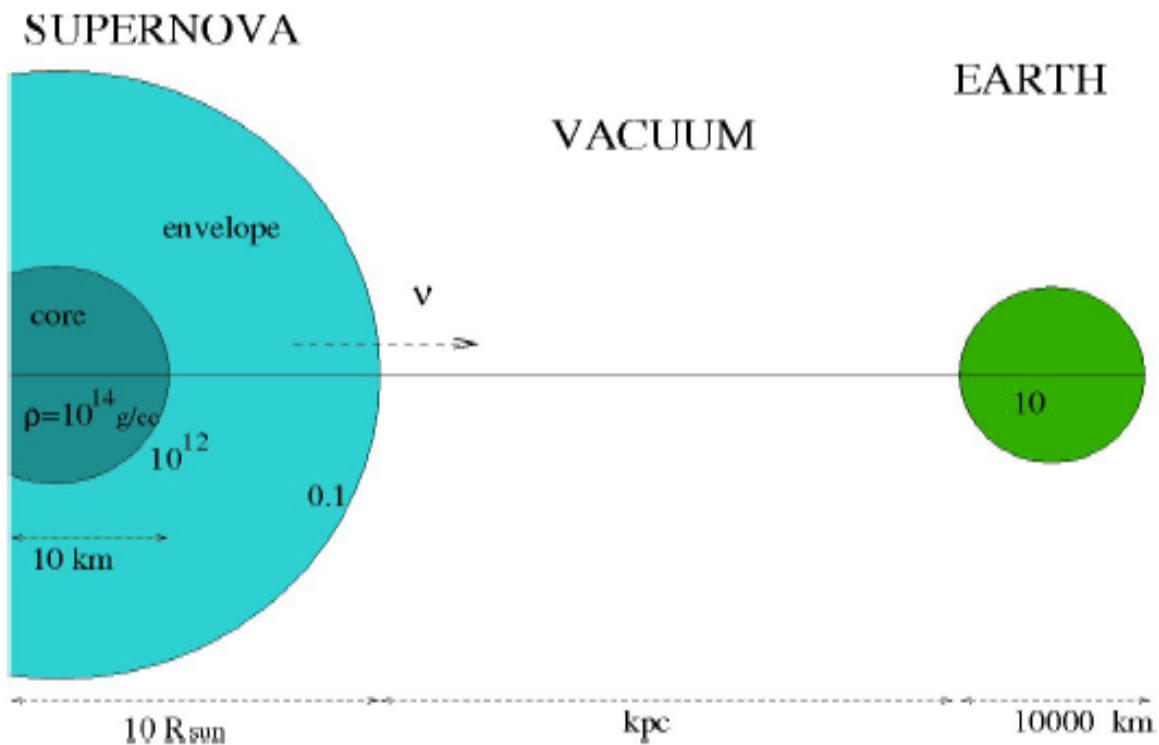
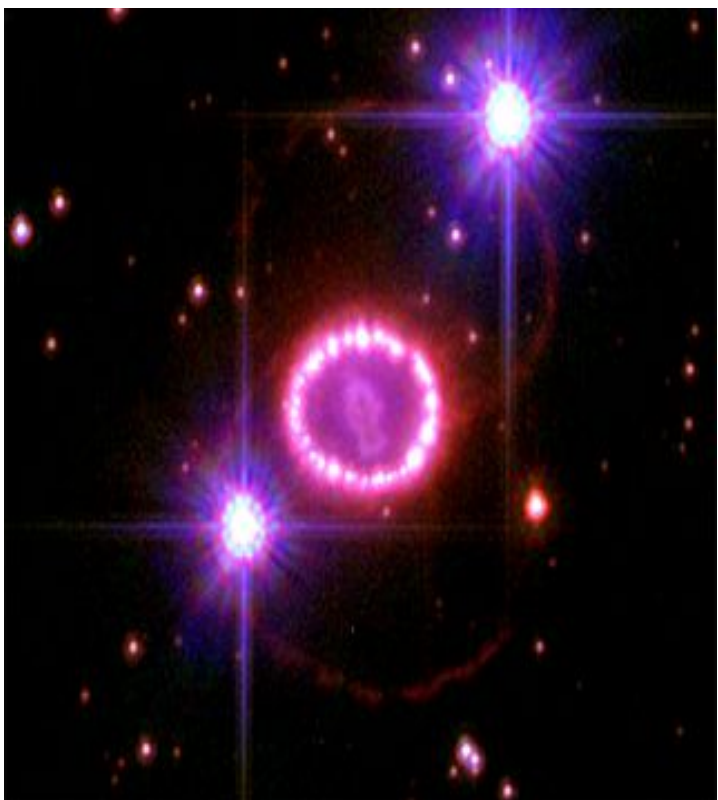
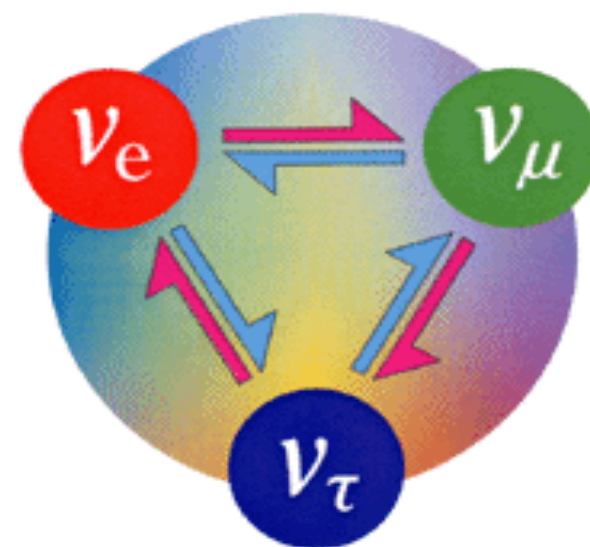
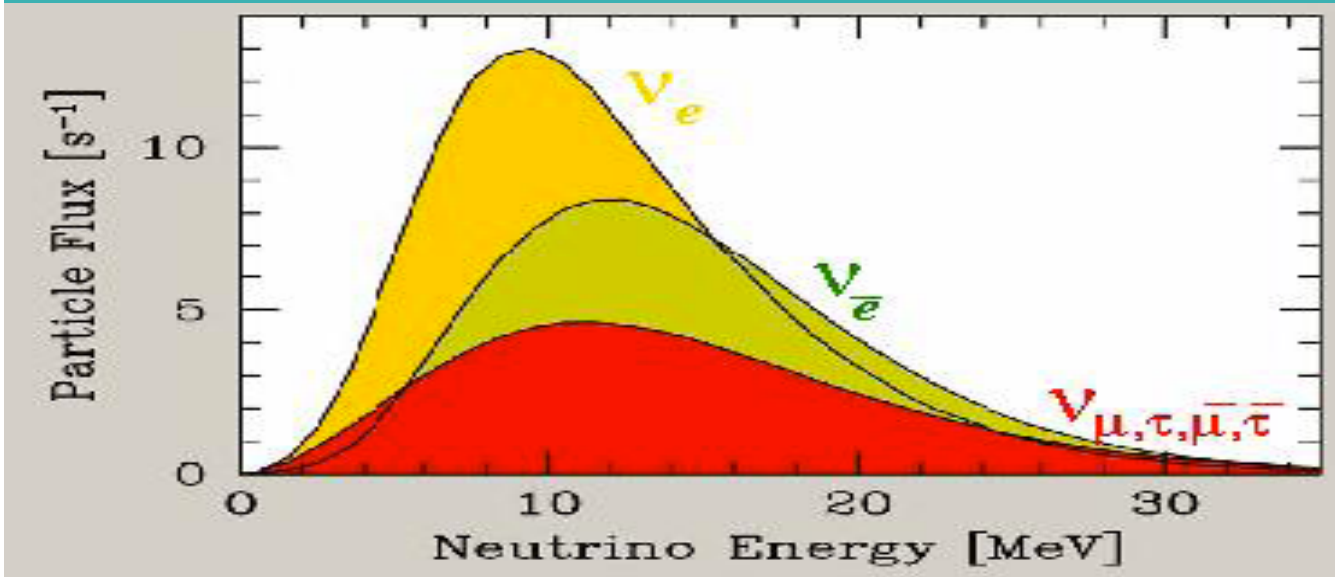


FIG. 2: Elastic DM scattering cross section with a nucleon versus DM mass. We compare present [16, 17] and future [18, 19] sensitivities with our model expectations, for $m_H = 120$ GeV and $\tan \beta = 0.5, 1, 5$ (grey solid lines).



NEUTRINO PROPERTIES FROM SUPERNOVAE



NSI induced flavor conversion near SN-core

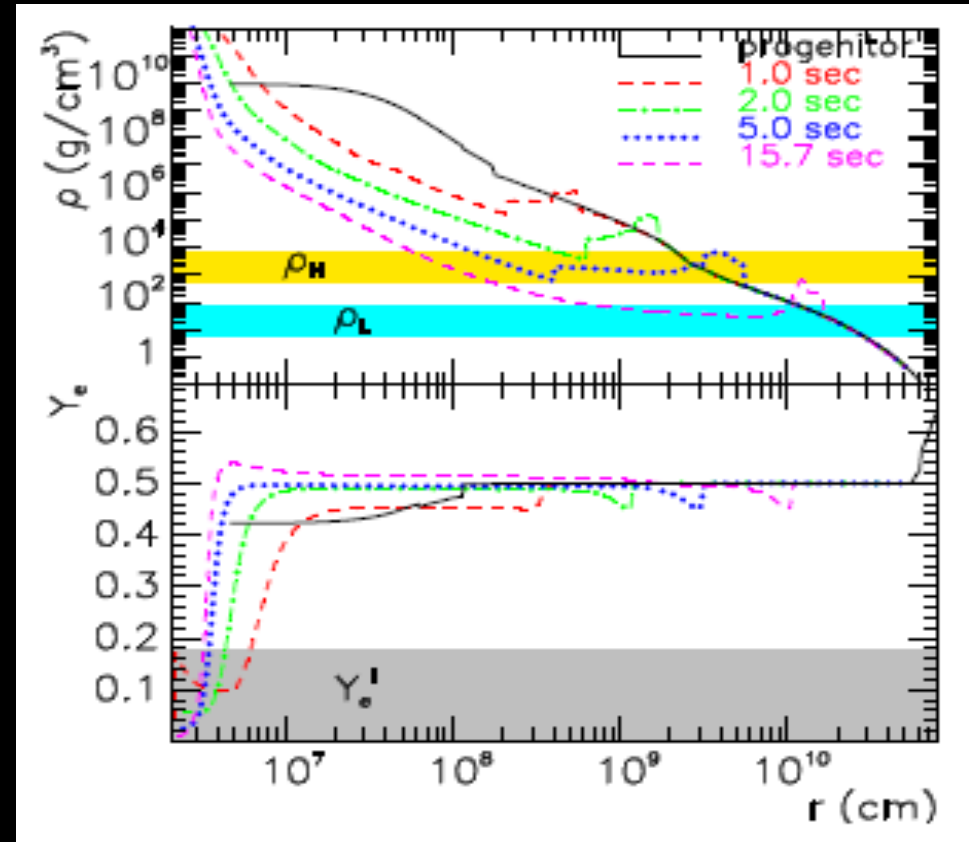
atm

sol

Interplay between collective effects and non-standard interactions in supernova

Esteban-Pretel et al PRD76 (2007) 053001 & PRD81 (2010) 063003

NSI



OSCILLATIONS MAINLY ROBUST but **NEED TO GO BEYOND**

THE ORIGIN OF NEUTRINO MASS REMAINS A MYSTERY

NSI and LFV searches help distinguish the **heavy** from **light** messenger models
in low-scale seesaw models, e.g. inverse & linear, **messengers may be produced**

SUSY ORIGIN OF NEUTRINO MASS TESTABLE AT LHC

DISPLACED VERTEX searches probe neutrino mass scale
LSP DECAY PATTERN probes neutrino **mixing**

STILL DO NOT UNDERSTAND FLAVOR

flavor models **correlate LFV phenomena & oscillations**
if flavor is linked to unification, **hard to reconcile lepton & quark mixings**

NEUTRINOS -- COSMO CONNECTION (not covered)

- **thermal LEPTOGENESIS** from low-scale seesaw
- sneutrinos as **DARK MATTER** in inverse seesaw
- **DARK MATTER** stabilized by neutrino flavor symmetry
- **majoron** as Decaying **DARK MATTER**

THANK YOU

next are BACKUP slides

