Vainshtein mechanism in quasi-dilaton massive gravity

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In progress

dRGT massive gravity

de Rham, Gabadadze, Tolley (2011)

 \checkmark The action

$$S_{MG} = \frac{M_{\rm Pl}^2}{2} \int d^4x \sqrt{-g} \left[R - \frac{m^2}{4} \left(\mathcal{U}_2 + \alpha_3 \mathcal{U}_3 + \alpha_4 \mathcal{U}_4 \right) \right] + S_m[g_{\mu\nu}, \psi]$$

where the potential is given by

$$\mathcal{U}_{2} = 2\varepsilon_{\mu\alpha\rho\sigma}\varepsilon^{\nu\beta\rho\sigma}\mathcal{K}^{\mu}_{\ \nu}\mathcal{K}^{\alpha}_{\ \beta} = 4\left([\mathcal{K}^{2}] - [\mathcal{K}]^{2}\right)$$
$$\mathcal{U}_{3} = \varepsilon_{\mu\alpha\gamma\rho}\varepsilon^{\nu\beta\delta\rho}\mathcal{K}^{\mu}_{\ \nu}\mathcal{K}^{\alpha}_{\ \beta}\mathcal{K}^{\gamma}_{\ \delta} = -[\mathcal{K}]^{3} + 3[\mathcal{K}][\mathcal{K}^{2}] - 2[\mathcal{K}^{3}]$$
$$\mathcal{U}_{4} = \varepsilon_{\mu\alpha\gamma\rho}\varepsilon^{\nu\beta\delta\sigma}\mathcal{K}^{\mu}_{\ \nu}\mathcal{K}^{\alpha}_{\ \beta}\mathcal{K}^{\gamma}_{\ \delta}\mathcal{K}^{\rho}_{\ \sigma} = -[\mathcal{K}]^{4} + 6[\mathcal{K}]^{2}[\mathcal{K}^{2}] - 3[\mathcal{K}^{2}]^{2} - 8[\mathcal{K}][\mathcal{K}^{3}] + 6[\mathcal{K}^{4}]$$

and

$$\mathcal{K}^{\mu}_{\ \nu} = \delta^{\mu}_{\ \nu} - \sqrt{\eta_{ab} g^{\mu\alpha} \partial_{\alpha} \phi^a \partial_{\nu} \phi^b}$$

(BD) Ghost free massive gravity

Hassan and Rosen (2011)

Quasi-dilaton theory

D'Amico, Gabadadze, Hui, Pirtskhalava (2012)

✓ Symmetry

$$\sigma \to \sigma - \alpha M_{\rm Pl}, \qquad \phi^a \to e^\alpha \phi^a$$

 \checkmark Define the new tensor

$$\mathcal{K}^{\mu}_{\ \nu} = \delta^{\mu}_{\ \nu} - e^{\sigma/M_{\rm Pl}} \sqrt{\eta_{ab} g^{\mu\alpha} \partial_{\alpha} \phi^a \partial_{\nu} \phi^b}$$

K satisfies the symmetry

 \checkmark The action

$$S = \frac{M_{\rm Pl}^2}{2} \int d^4x \sqrt{-g} \left[R - \frac{\omega}{M_{\rm Pl}^2} g^{\mu\nu} \partial_\mu \sigma \partial_\nu \sigma - \frac{m^2}{4} \left(\mathcal{U}_2 + \alpha_3 \mathcal{U}_3 + \alpha_4 \mathcal{U}_4 \right) \right] + S_m[g_{\mu\nu}, \psi]$$

massive graviton + scalar (DOF=6)

Decoupling limit in QMG

D'Amico, Gabadadze, Hui, Pirtskhalava (2012)

✓ Decoupling limit

 $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu} \qquad \phi^a = \delta^a_\mu x^\mu - \eta^{a\mu} \partial_\mu \pi / M_{\rm Pl} m^2$ $M_{\rm Pl} \to \infty, \qquad m \to 0, \qquad \Lambda = (M_{\rm Pl} m^2)^{1/3} = \text{fixed}, \qquad \frac{T_{\mu\nu}}{M_{\rm Pl}} = \text{fixed}$

\checkmark The action

$$\mathcal{L}_{\rm DL} = -\frac{1}{4} h^{\mu\nu} \mathcal{E}^{\alpha\beta}_{\mu\nu} h_{\alpha\beta} - \frac{\omega}{2} \partial^{\mu} \sigma \partial_{\mu} \sigma - h^{\mu\nu} \left[\frac{1}{4} \varepsilon_{\mu} \varepsilon_{\nu} \Pi - \frac{\alpha}{4\Lambda^{3}} \varepsilon_{\mu} \varepsilon_{\nu} \Pi \Pi - \frac{\beta}{2\Lambda^{6}} \varepsilon_{\mu} \varepsilon_{\nu} \Pi \Pi \Pi \right] + \sigma \left[\gamma_{0} \varepsilon \varepsilon \Pi + \frac{\gamma_{1}}{\Lambda^{3}} \varepsilon \varepsilon \Pi \Pi + \frac{\gamma_{2}}{\Lambda^{6}} \varepsilon \varepsilon \Pi \Pi \Pi + \frac{\gamma_{3}}{\Lambda^{9}} \varepsilon \varepsilon \Pi \Pi \Pi \Pi \right] + \frac{1}{M_{\rm Pl}} h^{\mu\nu} T_{\mu\nu}$$

 $\varepsilon \varepsilon \Pi \equiv \varepsilon^{\mu \alpha \beta \gamma} \varepsilon^{\nu}_{\ \alpha \beta \gamma} \partial_{\mu} \partial_{\nu} \pi \qquad \alpha, \beta, \text{ and } \gamma_{i} = \text{functions of } \alpha_{3}, \text{ and } \alpha_{4}$ $\varepsilon_{\mu} \varepsilon_{\nu} \Pi \equiv \varepsilon^{\alpha \gamma \delta}_{\mu} \varepsilon^{\beta}_{\nu \gamma \delta} \partial_{\alpha} \partial_{\beta} \pi \qquad \Pi_{\mu \nu} = \partial_{\mu} \partial_{\nu} \pi$

Decoupling limit in QMG

✓ Transformation

$$h_{\mu\nu} \to h_{\mu\nu} + \pi \eta_{\mu\nu} + \frac{\alpha}{\Lambda^3} \partial_\mu \pi \partial_\nu \pi$$

 \checkmark The action with $\beta=0$

Galileon terms

$$\begin{split} \mathcal{L}_{\mathrm{DL}} &= -\frac{1}{4} h^{\mu\nu} \mathcal{E}_{\mu\nu}^{\alpha\beta} h_{\alpha\beta} - \frac{\omega}{2} \partial^{\mu} \sigma \partial_{\mu} \sigma \left(-\frac{1}{8} \pi \left[\varepsilon \varepsilon \Pi - \frac{2\alpha}{\Lambda^{3}} \varepsilon \varepsilon \Pi \Pi + \frac{\alpha^{2}}{\Lambda^{6}} \varepsilon \varepsilon \Pi \Pi \Pi \right] \right) \\ &+ \left(\sigma \left[\gamma_{0} \varepsilon \varepsilon \Pi + \frac{\gamma_{1}}{\Lambda^{3}} \varepsilon \varepsilon \Pi \Pi + \frac{\gamma_{2}}{\Lambda^{6}} \varepsilon \varepsilon \Pi \Pi \Pi + \frac{\gamma_{3}}{\Lambda^{9}} \varepsilon \varepsilon \Pi \Pi \Pi \Pi \right] \right) \\ &+ \frac{1}{M_{\mathrm{Pl}}} h^{\mu\nu} T_{\mu\nu} + \frac{1}{M_{\mathrm{Pl}}} \pi T + \frac{\alpha}{M_{\mathrm{Pl}} \Lambda^{3}} \partial_{\mu} \pi \partial_{\nu} \pi T^{\mu\nu} \quad \text{Bi-galileon interaction term} \end{split}$$

where

$$\begin{aligned} \alpha &= -\frac{3}{4}\alpha_3 - 1, \qquad \beta = -\frac{1}{8}\alpha_3 - \frac{1}{2}\alpha_4, \qquad \gamma_0 = \frac{1}{2} - \frac{2}{3}\alpha_5, \\ \gamma_1 &= \frac{3}{8}\alpha_3 - \frac{1}{2} - \alpha_5, \qquad \gamma_2 = \frac{1}{2}\alpha_4 - \frac{3}{8}\alpha_3 - \frac{2}{3}\alpha_5, \qquad \gamma_3 = -\frac{1}{2}\alpha_4 - \frac{1}{6}\alpha_5. \end{aligned}$$

Spherically symmetric case

 \checkmark Master equation for λ

$$\frac{3}{2}\left(1-\frac{6}{\omega}\right)\lambda - \left(3\alpha + \frac{36\gamma_1}{\omega}\right)\lambda^2 + \left(\alpha^2 - \frac{32\gamma_1^2 + 24\gamma_2}{\omega}\right)\lambda^3 - 40\frac{\gamma_1\gamma_2}{\omega}\lambda^4 - \frac{12\gamma_2^2}{\omega}\lambda^5 = \left(\frac{r_*}{r}\right)^3$$
$$\lambda_{\sigma} = \frac{1}{\omega}(3\lambda + 4\gamma_1\lambda^2 + 2\gamma_2\lambda^3)$$

where

$$\lambda \equiv \frac{\pi'}{\Lambda^3 r}, \qquad \lambda_{\sigma} \equiv \frac{\sigma'}{\Lambda^3 r}, \qquad r_* \equiv \left(\frac{M}{M_{\rm Pl}^2 m^2}\right)^{1/3}$$

✓ Solution inside Vainshtein radius

$$\lambda \simeq -\left(rac{3\omega}{16(1+lpha)^2}
ight)^{1/5} \left(rac{r_*}{r}
ight)^{3/5} rac{1/5}{2}$$
 Negative λ (repulsive force)

 π is suppressed by usual gravity \rightarrow Newton gravity

Spherically symmetric case

✓ Solution outside Vainshtein radius

$$\frac{3}{2}\left(1-\frac{6}{\omega}\right)\lambda - \left(3\alpha + \frac{36\gamma_1}{\omega}\right)\lambda^2 + \left(\alpha^2 - \frac{32\gamma_1^2 + 24\gamma_2}{\omega}\right)\lambda^3 - 40\frac{\gamma_1\gamma_2}{\omega}\lambda^4 - \frac{12\gamma_2^2}{\omega}\lambda^5 = \left(\frac{r_*}{r}\right)^3$$

 $\lambda = 0$ and 4 constants

0

 \checkmark λ =0 solution (Asymptotically flat solution)

$$\lambda \simeq \frac{2\omega}{3(\omega - 6)} \left(\frac{r_*}{r}\right)^3 \qquad \qquad \pi' \propto \frac{1}{r^2}$$
vDVZ solution

 \checkmark λ =constant solutions (Asymptotic cosmological solutions)

$$\lambda\simeq\lambda_1,\,\,\lambda_2,\,\,\lambda_3,\,\,\lambda_4$$

 $h_{00} \propto h_{ij} \propto r^2$

 $\pi' \propto r$

Ghost ??

 $\pi \to \Phi(r) + \phi(t, x)$ $\sigma \to \Psi(r) + \psi(t, x)$

✓ Quadratic Lagrangian

$$\mathcal{L}_{\mathrm{DL}}^{(2)} = \mathcal{A}_1(\partial_t \phi)^2 - \mathcal{A}_2(\partial_r \phi)^2 - \mathcal{A}_3(\partial_\Omega \phi)^2 + \mathcal{B}_1(\partial_t \psi)^2 - \mathcal{B}_2(\partial_r \psi)^2 - \mathcal{B}_3(\partial_\Omega \psi)^2 + \mathcal{C}_1\partial_t \phi \,\partial_t \psi - \mathcal{C}_2\partial_r \phi \,\partial_r \psi - \mathcal{C}_3\partial_\Omega \phi \,\partial_\Omega \psi$$

where

$$\begin{split} \mathcal{A}_{1} &= \frac{3}{4} - \frac{1}{\Lambda^{3}} \left[\frac{3}{2} \alpha \left(\Phi'' + 2\frac{\Phi'}{r} \right) + 2\gamma_{1} \left(\Psi'' + 2\frac{\Psi'}{r} \right) \right] \\ &+ \frac{1}{\Lambda^{6}} \left[\frac{3}{2} \alpha^{2} \left(\frac{\Phi'^{2}}{r^{2}} + 2\frac{\Phi'\Phi''}{r} \right) - 6\gamma_{2} \left(\frac{\Psi'\Phi'}{r^{2}} + \frac{\Psi''\Phi'}{r} + \frac{\Psi'\Phi''}{r} \right) \right] \\ &- \frac{12\gamma_{3}}{\Lambda^{9}} \left(\frac{\Psi''\Phi'^{2}}{r^{2}} + 2\frac{\Psi'\Phi'\Phi''}{r^{2}} \right) \\ \mathcal{A}_{2} &= \frac{3}{4} + \frac{3}{2} \alpha a + 2\gamma_{1} b - \frac{1}{\Lambda^{3}} \left[3(\alpha + \alpha^{2} a - 2\gamma_{2} b) \frac{\Phi'}{r} + 2(2\gamma_{1} - 3\gamma_{2} a) \frac{\Psi'}{r} \right] \\ &+ \frac{1}{\Lambda^{6}} \left[\frac{3}{2} \left(\alpha^{2} + 8\gamma_{3} b \right) \frac{\Phi'^{2}}{r^{2}} - 6(\gamma_{2} - 4\gamma_{3} a) \frac{\Psi'\Phi'}{r^{2}} \right] \\ \mathcal{A}_{3} &= \frac{3}{4} + \frac{3}{2} \alpha a + 2\gamma_{1} b - \frac{1}{\Lambda^{3}} \left[\frac{3}{2} (\alpha + \alpha^{2} a - 2\gamma_{2} b) \left(\Phi'' + \frac{\Phi'}{r} \right) + (2\gamma_{1} - 3\gamma_{2} a) \left(\Psi'' + \frac{\Psi'}{r} \right) \right] \\ &+ \frac{1}{\Lambda^{6}} \left[\frac{3}{2} \left(\alpha^{2} + 8\gamma_{3} b \right) \frac{\Phi'\Phi''}{r} - 3(\gamma_{2} - 4\gamma_{3} a) \frac{\Phi'\Psi'' + \Psi'\Phi''}{r} \right] \\ \mathcal{B}_{1} &= \frac{\omega}{2} \\ \mathcal{B}_{2} &= \frac{\omega}{2} \\ \mathcal{B}_{3} &= \frac{\omega}{2} \\ \mathcal{C}_{1} &= -6\gamma_{0} - \frac{4\gamma_{1}}{\Lambda^{3}} \left(\Phi'' + 2\frac{\Phi'}{r} \right) - 6\frac{\gamma_{2}}{\Lambda^{6}} \left(\frac{\Phi'^{2}}{r^{2}} + 2\frac{\Phi'\Phi''}{r} \right) - \frac{24\gamma_{3}}{\Lambda^{9}} \frac{\Phi'^{2}\Phi''}{r^{2}} \\ \mathcal{C}_{2} &= -6\gamma_{0} + 4\gamma_{1} b - \frac{8\gamma_{1} - 6\gamma_{2}(a + b)}{\Lambda^{3}} \frac{\Phi'}{r} - \frac{6\gamma_{2} - 8\gamma_{3}(2a + b)}{\Lambda^{6}} \frac{\Phi'^{2}}{r^{2}} \\ \mathcal{C}_{3} &= -6\gamma_{0} + 4\gamma_{1} b - \frac{4\gamma_{1} - 3\gamma_{2}(a + b)}{\Lambda^{3}} \left(\Phi'' + \frac{\Phi'}{r} \right) - \frac{6\gamma_{2} - 8\gamma_{3}(2a + b)}{\Lambda^{6}} \frac{\Phi'\Phi''}{r} \end{split}$$

Ghost ?? $\pi \to \Phi(r) + \phi(t, x)$ $\sigma \to \Psi(r) + \psi(t, x)$

✓ The static clump of dust of constant density $T_{\mu\nu} = \rho \delta^0_{\mu} \delta^0_{\nu} \theta(R-r)$

✓ Quadratic Lagrangian (Inside the source + Vainshtein radius)

$$\begin{split} \mathcal{L}_{DL}^{(2)} \supset \left[\alpha \frac{\rho}{M_{\text{Pl}} \Lambda^3} - \left[\frac{3^{4/5}}{10 \times 2^{1/5}} \frac{(1+\alpha)(10+7\alpha)}{\omega} \left(\frac{\omega}{(1+\alpha)^2} \right)^{4/5} \left(\frac{r_*}{R} \right)^{12/5} \right] (\partial_t \phi)^2 \\ \frac{\alpha}{M_{\text{Pl}} \Lambda^3} \partial_\mu \pi \partial_\nu \pi T^{\mu\nu} = \alpha \left(\frac{r_*}{R} \right)^3 & \text{from galileon and bi-galileon} \\ \\ \text{For} \quad R \ll \xi^{-5/3} r_* \qquad \alpha > 0 \qquad \qquad \xi \equiv \frac{(1+\alpha)(10+7\alpha)}{\alpha \omega} \left(\frac{\omega}{(1+\alpha)^2} \right)^{4/5} \end{split}$$

✓ Quadratic Lagrangian (Outside the source, but inside Vainshtein radius)

$$\mathcal{L}_{DL}^{(2)} \supset -\frac{3^{4/5}}{10 \times 2^{1/5}} \frac{(1+\alpha)(10+7\alpha)}{\omega} \left(\frac{\omega}{(1+\alpha)^2}\right)^{4/5} \left(\frac{r_*}{r}\right)^{12/5} (\partial_t \phi)^2$$

Negative sign of the kinetic term

Ghost ??

 $\pi \to \Phi(r) + \phi(t, x)$ $\sigma \to \Psi(r) + \psi(t, x)$

 \checkmark Quadratic action for π and σ

$$\mathcal{L}_{\mathrm{DL}}^{(2)} \supset \mathcal{A}_1(\partial_t \phi)^2 + \mathcal{B}_1(\partial_t \psi)^2 + \mathcal{C}_1(\partial_t \phi)(\partial_t \psi)$$
$$= \mathcal{A}_1 \left(\partial_t \phi + \frac{\mathcal{C}_1}{2\mathcal{A}_1} \partial_t \psi \right)^2 + \left(\mathcal{B}_1 - \frac{\mathcal{C}_1^2}{4\mathcal{A}_1} \right) (\partial_t \psi)^2$$

✓ Inside Vainshtein radius,

$$\mathcal{B}_1 = \omega/2$$
$$\frac{\mathcal{C}_1^2}{4\mathcal{A}_1} \propto \left(\frac{r_*}{r}\right)^{12/5}$$

$$\mathcal{L}_{DL}^{(2)} \simeq \mathcal{A}_1 (\partial_t \chi)^2 - \frac{\mathcal{C}_1^2}{4\mathcal{A}_1} (\partial_t \psi)^2$$

One of the scalar field has always ghost !!!

Summary



Quasi-dilaton massive gravity with $\beta=0$ (Restricted bi-galileon)

• does not have "healthy" asymptotically flat solution (vDVZ

solution) as well as cosmological solution

It seems that there is some "healthy" time-dependent solutions (in progress)