

Workshop Summary

as a basis for discussion
immediately after this

Caution:

This is based on talks that I listened to during the workshop and thus does not necessarily reflect my own view on the subject. Moreover, this may differ from what the speakers originally intended (e.g. because of my misunderstanding).

Shinji Mukohyama (Kavli IPMU)

Monday (14:00-18:00)

- Cedric Deffayet (APC Paris)
Ghost free massive gravity with vierbeins
- Fawad Hassan (Stockholm)
Higher Derivative and Conformal Gravity
Actions from Bimetric and Partially Massless
Bimetric Theories
- Keisuke Izumi (LeCosPA)
Is there acausality in Massive Gravity?
- Robert Caldwell (Dartmouth)
Nonlinearity of Superposition in Massive
Gravity

Deffayet “Ghost free massive gravity with vierbeins”

- Fierz-Pauli
10 (metric components) – 4 (Bianchi) – 1 (extra constraint) = 5 d.o.f. of massive spin-2
- Nonlinear FP
 $m^2 V(M)$ with $M^\mu{}_\nu = g^{\mu\rho} f_{\rho\nu}$
BD ghost @ nonlinear level
- Metric dRGT
 $m^2 \beta_{1,2,3} F_{1,2,3}(\gamma)$
 $\gamma^\mu{}_\sigma \gamma^\sigma{}_\nu = g^{\mu\sigma} f_{\sigma\nu}$
no BD ghost
 $F_1(X) = [X]$
 $F_2(X) = ([X]^2 - [X^2])/2$
 $F_3(X) = ([X]^3 - 3[X][X^2] + 2[X^3])/6$

- 2 metrics \rightarrow 2 vielbeins

$$\eta_{AB} E^A{}_{\mu} E^B{}_{\nu} = g_{\mu\nu} \quad \eta_{AB} L^A{}_{\mu} L^B{}_{\nu} = f_{\mu\nu}$$

$$E_A{}^{\mu} L_{B\mu} = E_B{}^{\mu} L_{A\mu} \rightarrow 6 \text{ constraints}$$

$$16 \text{ (vielbein components)} - 6 \text{ (constraints)} = 10 \text{ (metric components)}$$

- Vielbein dRGT

$$S = M_P^{D-2} \int \Omega^{AB} \wedge E_{AB}^* - M_P^{D-2} m^2 \sum_{n=0}^{D-1} \beta_n \int L^{A_1} \wedge \dots \wedge L^{A_n} \wedge E_{A_1 \dots A_n}^*$$

$$G_A = t_A$$

- $10 - 4 \text{ (Dt}_A=0) - 1 \text{ (} m^A \wedge G_A = m^A \wedge t_A \text{ ?)} = 5?$

works when only $\beta_1 \neq 0$ or when only $\beta_2 \neq 0$

not clear for general cases

Hassan “Higher Derivative and Conformal Gravity Actions from Bimetric and Partially Massless Bimetric Theories”

- FP, nonlinear FP, BD ghost
- dRGT: decoupling limit analysis of $g^{-1}f$ theory
→ necessary condition
HR “it is indeed sufficient”
- Give dynamics to $f \rightarrow$ bi-metric theory
solutions with $f = c^2 g$ ($c \leftarrow \Lambda_f = \Lambda_g/c^2$)
FP mass around $f = c^2 g \rightarrow 2 + 5 = 7$ d.o.f.
- New gauge symmetry @ Higuchi bound \rightarrow
partially massless FP theory with $5 - 1 = 4$ d.o.f.

- Nonlinear extension \rightarrow PM bimetric
 new gauge symmetry would detune the c of
 $f = c^2 g \rightarrow c$ undetermined by $\Lambda_f = \Lambda_g / c^2 \rightarrow$
 one-parameter (β_2) subset of bimetric theories
- HD (or conformal) gravity from bigravity (or
 PM bigravity)
 solve g -eom w.r.t. f & substitute to the action
 & derivative expansion \rightarrow series of higher
 curvature terms
 truncation at finite order \rightarrow ghost
 infinite sum \rightarrow no-ghost completion of HD (or
 conformal) gravity?

Izumi “Is there acausality in Massive Gravity?”

- Deser-Waldron’ claim on acausality of dRGT is due to misuse of characteristics analysis
- Characteristics analysis: $\det(\text{time kinetic matrix}) = 0$ AFTER solving constraints \rightarrow instantaneous mode \rightarrow acausality
- dRGT eliminates BD ghost \rightarrow would-be BD ghost has vanishing time kinetic term $\rightarrow \det(\text{time kinetic matrix}) = 0$ before solving constraints \rightarrow this does NOT imply acausality

Caldwell “Nonlinearity of Superposition in Massive Gravity”

- Nonlinearity \rightarrow Vainshtein screening
- What happens if two Vainshtein regions overlap?
- Nonlinear solution $<$ naive superposition
- Modified gravity effects tend to be hidden
- Relaxation time scale \gg dynamical time scale?
If yes, then we may still hope to see modified gravity effects unsuppressed.

Tuesday (9:45-15:00)

- Mikhail Volkov (Tours U, CNRS)
Self-accelerated cosmologies and black holes
in the ghost-free bigravity
- Emir Gumrukcuoglu (Kavli IPMU)
Stability of cosmological solutions in massive
gravity
- Chunshan Lin (Kavli IPMU)
Cosmological perturbations in extended
massive gravity

Volkov “Self-accelerated cosmologies and black holes in the ghost-free bigravity”

- Solution in dRGT \rightarrow solution in bigravity
2 more excitations \rightarrow more solutions in bigravity
- FLRW with non-bidiagonal metrics
f & g separately FLRW
f-commoving observers \neq g-commoving observers
spherical symmetry & $T^t_r = T^t_t - T^r_r = 0$ & constraint eq
 \rightarrow PDE \rightarrow infinite # of solutions
massive gravity limit \rightarrow 2-parameter (k,q) family of special solutions
- FLRW with bidiagonal metrics
constraint eq \rightarrow 3 branches of solutions
homothetic one (c.f. Soda) & other two with $a/b = \text{const.}$

- Static black holes
 f & g share the same horizon and surface gravity
 $f = c^2 g \rightarrow S (\Lambda_g=0), SdS (\Lambda_g>0), SAdS (\Lambda_g<0)$
- SAdS can carry short hair while S and SdS cannot. For S and SdS , would-be hair inevitably accretes toward BH.
- The hair survives in $r_h = 0$ limit \rightarrow lump solution
- Regular star solution with matter source \rightarrow
 Vainshtein screening in dRGT and bigravity

Gumrukcuoglu “Stability of cosmological solutions in massive gravity”

- Homogeneous & isotropic FLRW cosmological solution exists in dRGT
- Linear perturbation analysis \rightarrow vanishing kinetic terms for 3 out of 5 d.o.f. \rightarrow nonlinear analysis is needed
- Almost FLRW Bianchi I + linear perturbation \rightarrow leading nonlinear perturbation on FLRW
- Negative kinetic term inevitable @ cubic order \rightarrow all FLRW solutions in dRGT are unstable...

- In analogy with tachyon/ghost condensation, why don't we consider large deviation from FLRW?
→ anisotropic FLRW attractor solution: FLRW $g_{\mu\nu}$ but $O(1)$ anisotropy in $f_{\mu\nu}$
- Stable dynamics near anisotropic FLRW possible

Lin “Cosmological perturbations in extended massive gravity”

- All FLRW solutions in dRGT are unstable...
- We need (i) to break FLRW symmetry or (ii) to extend the theory. Let's consider (ii).
- dRGT < quasi-dilaton, varying mass < bi-metric
- Self-accelerating solution in quasi-dilaton always have a ghost in scalar sector.
- Varying mass theory may have stable regime. However, a viable self-accelerating solution has not yet been found.

Tuesday (16:00-18:00)

- Tsutomu Kobayashi (Rikkyo U)
Vainshtein mechanism in Horndeski's general scalar-tensor theory (and in massive gravity)
- Tetsuya Shiromizu (Kyoto U)
Concern on "ground state" of spacetimes with/without black holes
- Rampei Kimura (Hiroshima U)
Vainshtein mechanism in quasi-dilaton massive gravity

Kobayashi “Vainshtein mechanism in Horndeski's general scalar-tensor theory (and in massive gravity)”

- Vainshtein screening due to nonlinearity:
 $r_c^2 \ddot{\pi}$ can be $O(1)$ even if $\pi \ll 1$
- Horndeski's theory (= covariant Galileon) \rightarrow 2nd order eom
- Spherically-symmetric, static perturbation around Minkowski \rightarrow 2nd order ode \rightarrow 1st integral gives algebraic eq's for 1st derivatives with source $A(r) \sim M(r)/r^3$ if we consider various coefficients as constants \rightarrow can be solved @ each r

- Can we connect almost GR inner solution to asymptotically flat outer solution?
- Condition for smooth matching:
 $P(x,A)=0$ has a single root in $(x_-,0)$ for any $A>0$
 \rightarrow no local extrema in $x<0$
- 2nd derivative may become large near transition region if the above condition is marginally fulfilled \rightarrow observable signature in e.g. cluster lensing?

Shiromizu “Concern on ground state of spacetimes with/without black holes”

- Various theorems in GR
Positive mass theorem
Strictly stationary vacuum \rightarrow Minkowski
BH no hair (static, stationary)
- Even in Einstein Gauss-Bonnet in $D > 4$,
Strictly stationary vacuum \rightarrow Minkowski
- In dynamical Chern-Simons gravity,
Strictly stationary vacuum \rightarrow Minkowski
BH no hair (static)
- How about massive gravity?

Kimura “Vainshtein mechanism in quasi-dilaton massive gravity”

- Quasi-dilaton: dRGT + scalar
- Decoupling limit \rightarrow galileon & bi-galileon terms
- Spherically-symmetric, static perturbation around Minkowski \rightarrow 2nd order ode \rightarrow 1st integral gives algebraic eq's for 1st derivatives with source $A(r) \sim M(r)/r^3 \rightarrow$ can be solved @ each r
- Can we connect almost GR inner solution to asymptotically flat outer solution?
- However, there is a ghost (c.f. Lin). Can we found a healthy background?

Wednesday (9:45-12:00)

- Kei-ichi Maeda (Waseda U)
Anisotropic universes in bigravity theory
and homothetic metrics
- Jiro Soda (Kyoto U)
Inflation in bimetric gravity

Maeda “Anisotropic universes in bigravity theory and homothetic metrics”

- Classical stability of FLRW
Validity of cosmic no hair conjecture
- f & g of the same Bianchi type (in class A)
 $C^c_{ab} = n^{cd} \varepsilon_{dab} + \alpha (\delta^1_a \delta^c_b - \delta^1_b \delta^c_a)$ with $\alpha = 0$
- g- & f-potentials depend only on difference between f-anisotropy and g-anisotropy
- $f = c^2 g$ solution is an attractor
- massive GW \rightarrow shear density $\sim 1/a^3 \rightarrow$ may remain until now \rightarrow observable signature?

Soda “Inflation in bimetric gravity”

- If $m < H \rightarrow$ shear may slow roll \rightarrow cosmic no-hair violated? Higuchi bound as hair censorship?
- Small anisotropy on top of homothetic dS massless mode & massive mode \rightarrow graviton mass in FP action is identified
- Higuchi bound always satisfied by homothetic dS
- Slow-roll pot \rightarrow almost homothetic dS attractor
- [Multi-metric gravity in metric formulation absence of BD ghost \rightarrow loops forbidden]

Let's start discussion now!