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Cosmic Ray Electrons and Positrons from Fermi Gamma-Ray Pulsars

"Focus on Indirect Dark Matter Searches" IPMU – University of Tokyo Kashiwa, December 11, 2009 1. (as per talk title) contribution to CR e+efrom Fermi gamma-ray pulsars

(work in collaboration with UC Santa Cruz undergrad Lev Gendelev and grad Michael Dormody)

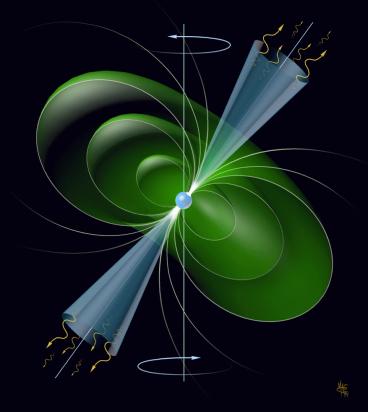
2. some thoughts about the "Fermi haze "

(work in collaboration with UC Santa Cruz grad Tim Linden)

• Rotation-powered Pulsars can seed e+e- direct pair production

(strong rotationally induced **electric fields** in the magnetoshpere accelerate and extract **e-** from stellar surface, which radiate **gamma rays**; gammas cascade produce **e+e- pairs**, escaping the magnetosphere from the polar cap regions)

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SNR & PWN shock acceleration

Propagation of charged species: diffusion equation

$$\frac{\partial f}{\partial t} = \frac{D(\gamma)}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial f}{\partial r} + \frac{\partial (P(\gamma) \ f)}{\partial \gamma} + Q.$$

$$f(r,t,\gamma)$$
 Distribution function

$$\gamma = E_{e^{\pm}}/m_e, \quad D(\gamma) \propto \gamma^{\delta}$$

$$P(\gamma) = p_0 + p_1\gamma + p_2\gamma^2$$

0: Coulomb; 1: Brems; 2: IC & Synch

(*) Atoyan, Aharonian, Volk, 1995

Approximate solution to the electron/positron distribution function^(*) (only IC and Synch losses – burst-like injection)

$$f(r,t,\gamma) = \frac{N_0 \gamma^{-\alpha}}{\pi^{3/2} r^3} (1 - p_2 t \gamma)^{\alpha - 2} \left(\frac{r}{r_{\text{dif}}} \right)^3 e^{-(r/r_{\text{dif}})^2}$$
$$\gamma < \gamma_{\text{cut}} \equiv \gamma_{\text{cut}}(t) = (p_2 t)^{-1}$$
$$p_2 = 5.2 \times 10^{-20} \frac{w_0}{1 \text{ eV/cm}^3} \text{ s}^{-1}$$
$$\begin{array}{c} \text{Cut-off in} \\ \text{Energy} \rightarrow \\ \text{Pulsar Lifetime} \end{array}$$
$$r_{\text{dif}}(\gamma,t) \simeq 2\sqrt{D(\gamma)t} \frac{1 - (1 - \gamma/\gamma_{\text{cut}})^{1-\delta}}{(1 - \delta)\gamma/\gamma_{\text{cut}}}.$$

Main feature of high-energy (10-1000 GeV) e⁺e⁻: they lose energy very efficiently

Energy losses $\sim E^2$, via synchrotron and inverse Compton

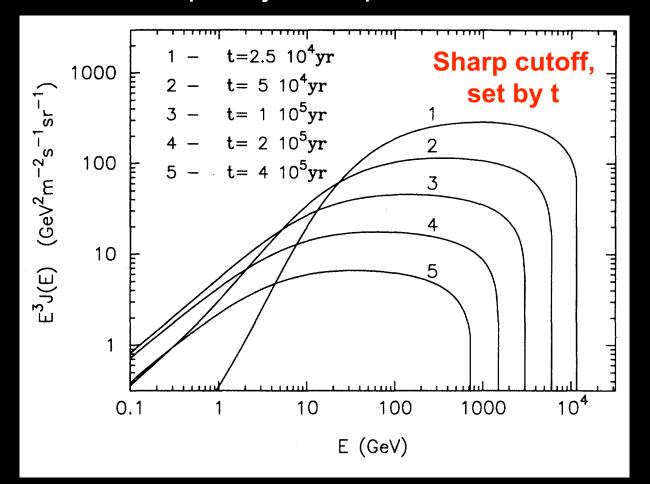
$$\frac{t_{\text{Lifetime}}}{yr} \approx 5 \times 10^5 \left(\frac{1 \text{ TeV}}{E}\right) \left[\left(\frac{B}{5 \ \mu \text{G}}\right)^2 + 1.6 \times \left(\frac{w}{1 \text{ eV/cm}^3}\right) \right]^{-1}$$

In conjunction with conventional diffusion models, this short radiative cooling time limits the sources of high energy electron/positron both in space and time

dist_{max}
$$\approx \sqrt{D_0 \times t} \approx 100 - 500 \text{ pc} \quad \left[D_0 \approx 10^{28} \text{ cm}^2 / \text{s} \right]$$

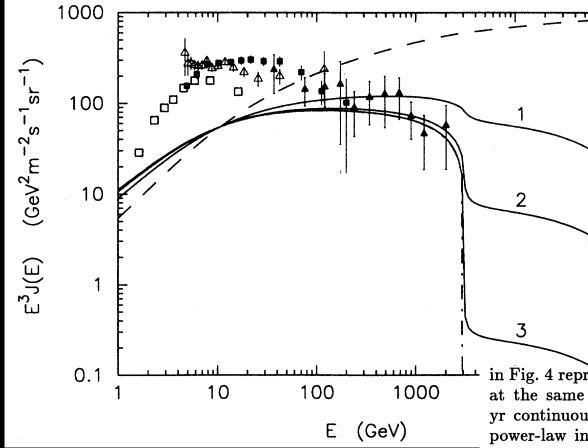
Astrophysical sources relevant for energetic e+e- production must be **young** (~10⁵ yr) and **nearby** (<kpc)

Example of a **burst**-like injection at different **times**, r = 100 pc, injection power-law: 2.2



Distance sets normalization, and affects spectrum, ($s=r/r_{dif}$, and more distant, more peaked) (*) Atoyan, Aharonian, Volk, 1995

The effects of a non-burst-like injection

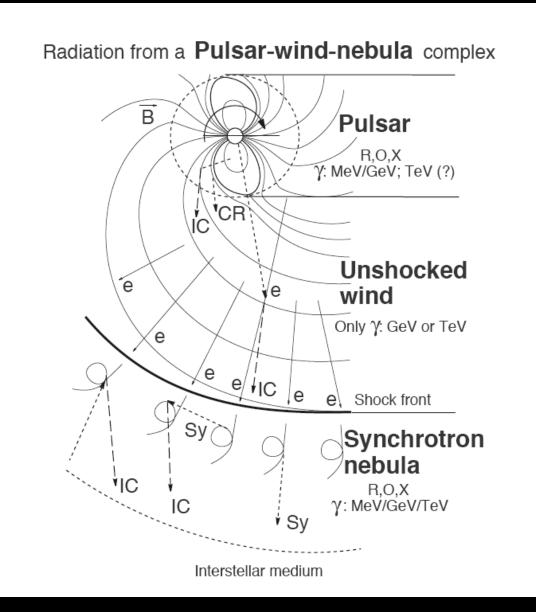


in Fig. 4 represent the fluxes of electrons from the source at the same distance r = 100 pc and of age $t = 10^5$ yr continuously injecting relativistic electrons with the power-law index $\alpha = 2.2$ into ISM, but with the total luminosity varying in time during $0 \le \tau \le t$ as

$$L_e(\tau) = \frac{L_0}{(1 + \tau/\tau_*)^k}$$
(24)

for three different values of the characteristic "decay" time τ_* : $\tau_*/t = 0.1$ (curve 1), $\tau_*/t = 0.01$ (curve 2), $\tau_*/t = 0.001$ (curve 3). This kind of time-dependent in-

(*) Atoyan, Aharonian, Volk, 1995



Electromagnetic processes and emission in a pulsar environment

Multi-Wavelength observations help:

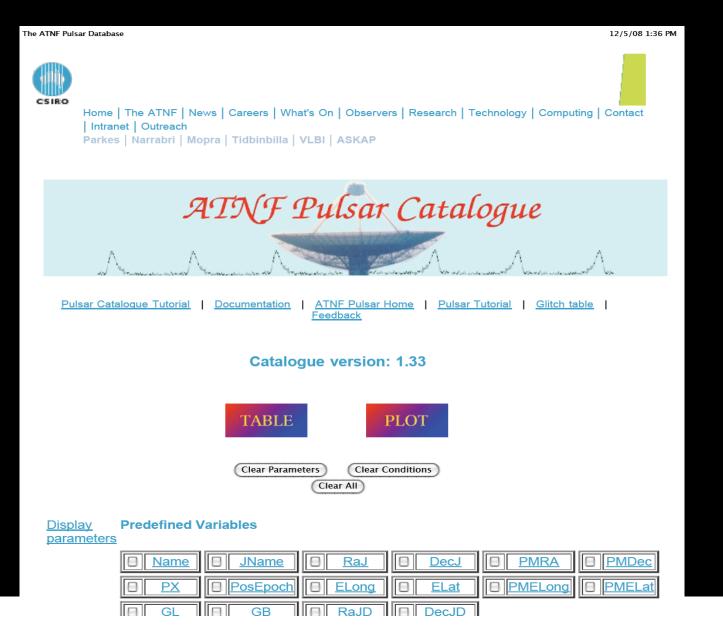
(i) set the scale of the total powerin relativistic e+e-

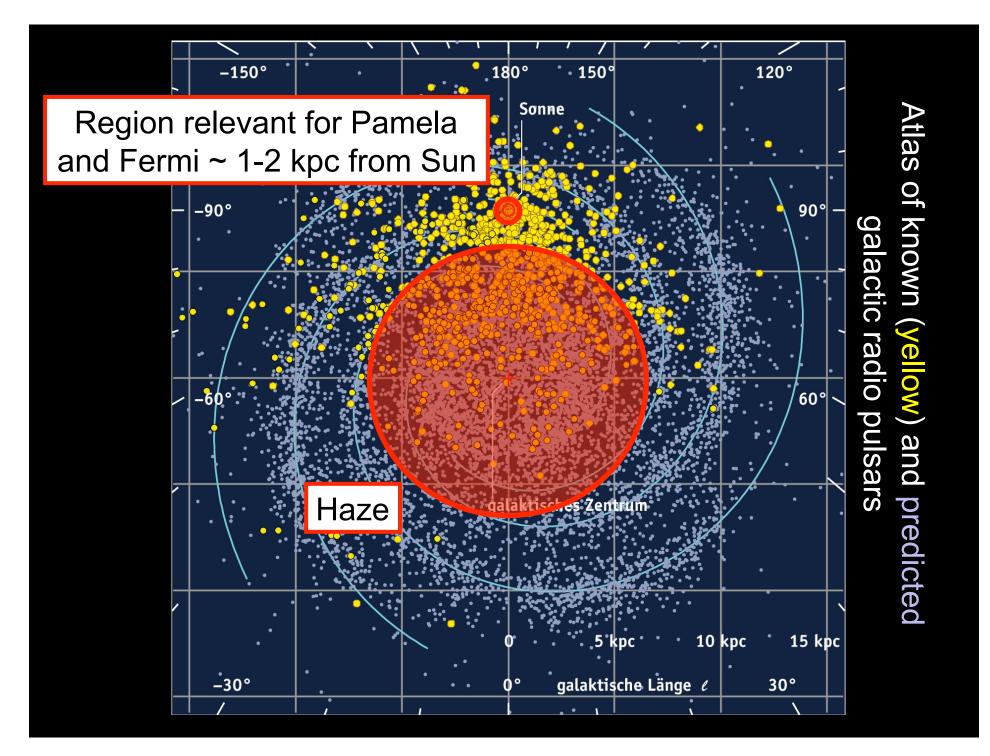
(ii) understand the e+e- spectrum

examples: • HESS J1825-137 • Vela X (HESS) • Geminga (Milagro)

(*) Aharonian and Bogovalov, 2002

An **asset** of the pulsar scenario: pulsars **exist**, detailed **catalogues**, very **accurate** data





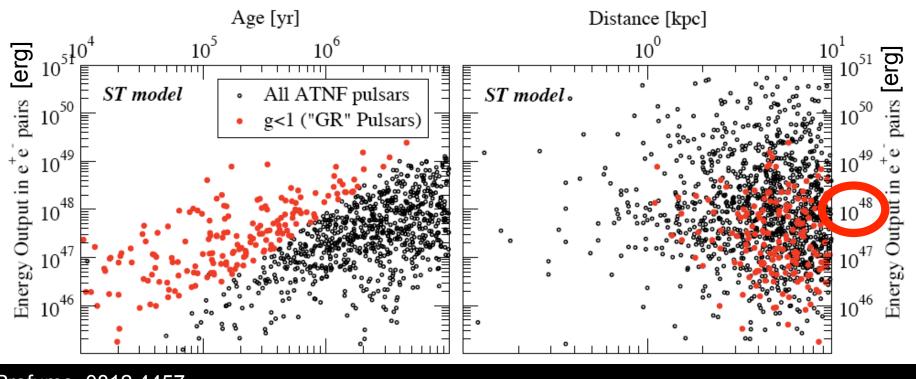
Not all pulsars are **gamma-ray** (and e+e-) pulsars

E.g., in the outer gap model, pair production sets in if

 $E_x E_\gamma > m_e^2$

This condition depends on magnetic field intensity and period

$$g = 5.5 \ P^{26/21} \ B_{12}^{-4/7} \ < \ 1.$$



Profumo, 0812.4457



Role of Fermi to assess the origin of high-energy CRE:

- 1. Accurate CRE **Spectral Information** (probably not conclusive by itself)
- 2. Local CRE source ? \rightarrow Compare the Inverse Compton and Bremss. emis. predicted from the measured CRE spectrum with diffuse gamma-ray data
- 3. Anisotropy: search for excess CRE from bright nearby pulsars (problem: pulsar proper motion !! also ongoing: East-West asymmetry search !!)
- Discovery and improved understanding of gamma-ray pulsars, guaranteed sources of e+e-



The Fermi First Pulsar Catalogue

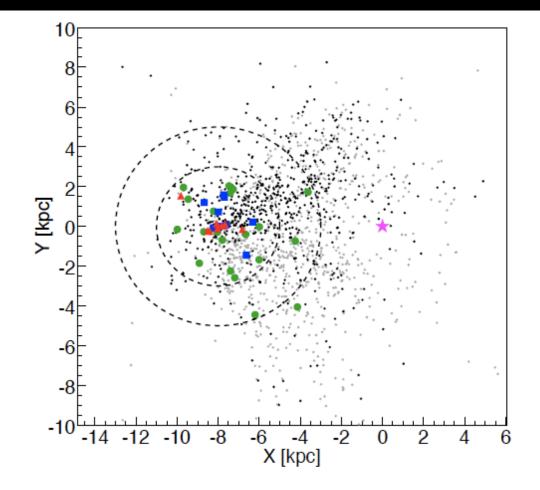
• Includes high-confidence (>) pulsed sources in the first 6-months of data

arXiv:0910.1608v1 [astro-ph.HE] 8 Oct 2009

- 16 gamma-ray selected (majority from blindsearch campaign)
- 24 radio-selected (via ephermerides, and of which 8 millisecond)
- 6 pre-Fermi GR pulsars

The First Fermi Large Area Telescope Catalog of Gamma-ray Pulsars

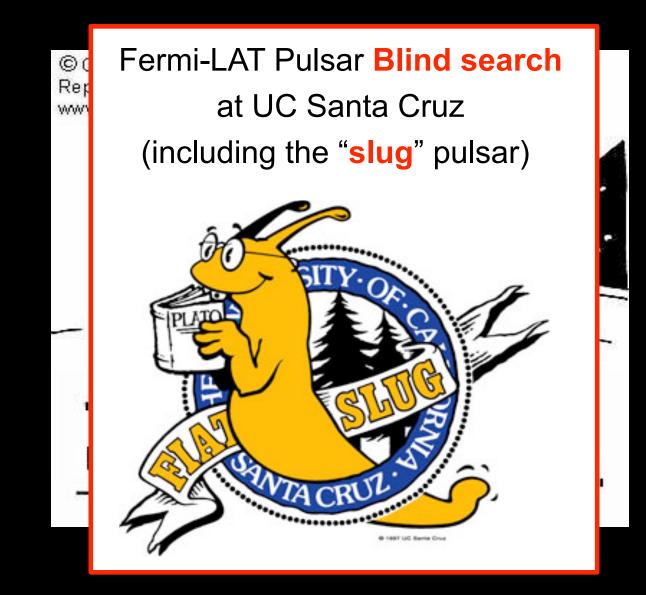
A. A. Abdo^{2,3}, M. Ackermann⁴, M. Ajello⁴, W. B. Atwood⁵, M. Axelsson^{6,7}, L. Baldini⁸ J. Ballet⁹, G. Barbiellini^{10,11}, M. G. Baring¹², D. Bastieri^{13,14}, B. M. Baughman¹⁵, K. Beehtol⁴, R. Bellazzini⁸, B. Berenji⁴, R. D. Blandford⁴, E. D. Bloom⁴, E. Bonamente^{16,17}, A. W. Borgland⁴, J. Bregeon⁸, A. Brez⁸, M. Brigida^{18,19}, P. Bruel²⁰, T. H. Burnett²¹, S. Buson¹⁴, G. A. Caliandro^{18,19,1}, R. A. Cameron⁴, F. Camilo²², P. A. Caraveo²³, J. M. Casandjian⁹, C. Cecchi^{16,17}, Ö. Çelik^{24,25,26}, E. Charles⁴, A. Chekhtman^{2,27}, C. C. Cheung²⁴, J. Chiang⁴, S. Ciprini^{16,17}, R. Claus⁴, I. Cognard²⁸, J. Cohen-Tanugi²⁹, L. R. Cominsky³⁰, J. Conrad^{31,7,32} R. Corbet^{24,26}, S. Cutini³³, P. R. den Hartog⁴, C. D. Dermer², A. de Angelis³⁴, A. de Luca^{23,35} F. de Palma^{18,19}, S. W. Digel⁴, M. Dormody⁵, E. do Couto e Silva⁴, P. S. Drell⁴, R. Dubois⁴, D. Dumora^{36,37}, C. Espinoza³⁸, C. Farnier²⁹, C. Favuzzi^{18,19}, S. J. Fegan²⁰, E. C. Ferrara^{24,1} W. B. Focke⁴, P. Fortin²⁰, M. Frailis³⁴, P. C. C. Freire³⁹, Y. Fukazawa⁴⁰, S. Funk⁴, P. Fuseo^{18,19} F. Gargano¹⁹, D. Gasparrini³³, N. Gehrels^{24,41}, S. Germani^{16,17}, G. Giavitto⁴², B. Giebels²⁰ N. Giglietto^{18,19}, P. Giommi²³, F. Giordano^{18,19}, T. Glanzman⁴, G. Godfrey⁴, E. V. Gotthelf²² I. A. Grenier⁹, M.-H. Grondin^{36,37}, J. E. Grove², L. Guillemot^{36,37}, S. Guiriec⁴³, C. Gwon², Y. Hanabata⁴⁰, A. K. Harding²⁴, M. Hayashida⁴, E. Hays²⁴, R. E. Hughes¹⁵, M. S. Jackson^{21,7,44} G. Jóhannesson⁴, A. S. Johnson⁴, R. P. Johnson⁵, T. J. Johnson^{24,41}, W. N. Johnson², S. Johnston⁴⁵, T. Kamae⁴, G. Kanbach⁴⁶, V. M. Kaspi⁴⁷, H. Katagiri⁴⁰, J. Kataoka^{48,49} N. Kawai^{48,50}, M. Kerr²¹, J. Knödlseder⁵¹, M. L. Koeian⁴, M. Kramer^{38,52}, M. Kuss⁸, J. Lande⁴, L. Latronico⁸, M. Lemoine-Goumard^{38,37}, M. Livingstone⁴⁷, F. Longo^{10,11}, F. Lopareo^{18,19} B. Lott^{26,27}, M. N. Lovellette², P. Lubrano^{16,17}, A. G. Lyne³⁸, G. M. Madejski⁴, A. Makeev^{2,27} R. N. Manchester⁴⁵, M. Marelli²³, M. N. Mazziotta¹⁹, W. McConville^{24,41}, J. E. McEnery²⁴ S. McGlynn^{44,7}, C. Meurer^{31,7}, P. F. Michelson⁴, T. Mineo⁸³, W. Mitthumsiri⁴, T. Mizuno⁴⁰, A. A. Moiseev^{25,41}, C. Monte^{18,19}, M. E. Monzani⁴, A. Morselli⁵⁴, I. V. Moskalenko⁴, S. Murgin⁴ T. Nakamori⁴⁸, P. L. Nolan⁴, J. P. Norris⁵⁵, A. Noutsos²⁸, E. Nuss²⁹, T. Ohsugi⁴⁰, N. Omodei⁸, E. Orlando⁴⁶, J. F. Ormes⁵⁵, M. Ozaki⁵⁶, D. Paneque⁴, J. H. Panetta⁴, D. Parent^{26,37,1}, V. Pelassa²⁹, M. Pepe^{16,17}, M. Pesce-Rollins⁸, F. Piron²⁹, T. A. Porter⁵, S. Raino^{18,19} R. Rando^{13,14}, S. M. Ransom⁵⁷, P. S. Ray², M. Razzano⁸, N. Rea^{58,59}, A. Reimer^{60,4}, O. D. M. Razzano⁸, N. Rea^{58,59}, A. Reimer^{60,4}, D. L. Linger, M. Razzano⁸, N. Rea^{58,59}, A. Reimer^{60,4}, D. L. Linger, M. Razzano⁸, N. Rea^{58,59}, A. Reimer^{60,4}, D. Linger, M. Razzano⁸, N. Rea^{58,59}, R. Linger, M. Razzano⁸, R. Razzano⁸, N. Rea^{58,59}, R. Linger, M. Razzano⁸, N. Rea^{58,59}, R. Linger, M. Razzano⁸, N. Rea^{50,59}, R. Linger, M. Razzano⁸, N. Rea^{58,59}, R. Linger, M. Razzano⁸, R. Linger, M. L



Galactic Center Gamma-Ray sel Millisecond P Other radio-loud

Fig. 3.— Galactic plane pulsar distribution (polar view). The star represents the Galactic center. The two circles centered at the Earth's position have radii of 3 kpc and 5 kpc. For pulsars with different possible distances, the nearer values from Table 5 are used. Note that the millisecond pulsars (MSPs), while having a significantly lower \dot{E} than the other pulsars (see Figure 8), are detectable due to their close proximity. The one exception (PSR J0218+4232) also exhibits a significantly higher \dot{E} than the other MSPs. Blue squares: gamma-ray-selected pulsars. Red triangles: millisecond gamma-ray pulsars. Green circles: all other radio loud gamma-ray pulsars. Black dots: Pulsars for which gamma-ray pulsation searches were conducted using rotational ephemerides. Gray dots: Known pulsars which were not searched for pulsations.

The Issue of naming new Pulsars



- Distance estimates from gamma-ray data
- 2. Simple + universal PSR emission model

The Contribution of Fermi Gamma-Ray Pulsars to the local Flux of Cosmic-Ray Electrons and Positrons

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Abstract. We analyze the contribution of gamma-ray pulsars from the first Fermi-
Large Area Telescope (LAT) catalogue to the local flux of electrons and positrons (e^+e^-) . We present new distance estimates for all Fermi gamma-ray pulsars, based

3. Estimate of **contribution** to local e+e- flux

Main results:

 gamma-ray pulsar contribution can be substantial from 10 of the brightest Fermi pulsars

10

11

- 6/10 are blind-search gamma-ray selected pulsars
- outline of regions where radio-pulsars might contribute without a gamma-ray signal

$$L_{\gamma} \approx w\dot{E} \approx C \times \left(\frac{\dot{E}}{10^{33} \text{ ergs}^{-1}}\right)^{1/2} \times 10^{33} \text{ ergs}^{-1}$$

w : Gamma-Ray conversion efficiency

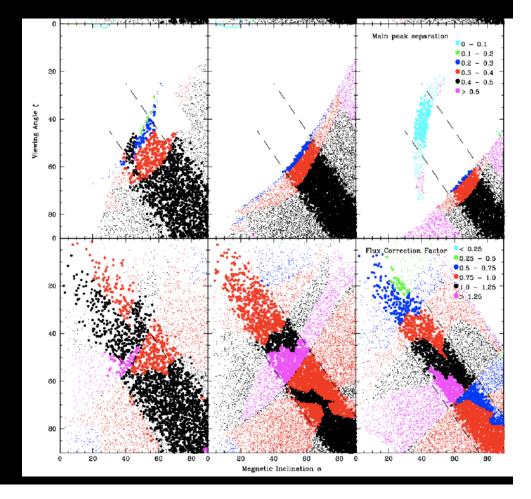
C : constant of O(1) – **fitted** with PSR subset with solid distance determinations

$$L_{\gamma} = 4\pi f_{\Omega} D^2 F_{\gamma},$$

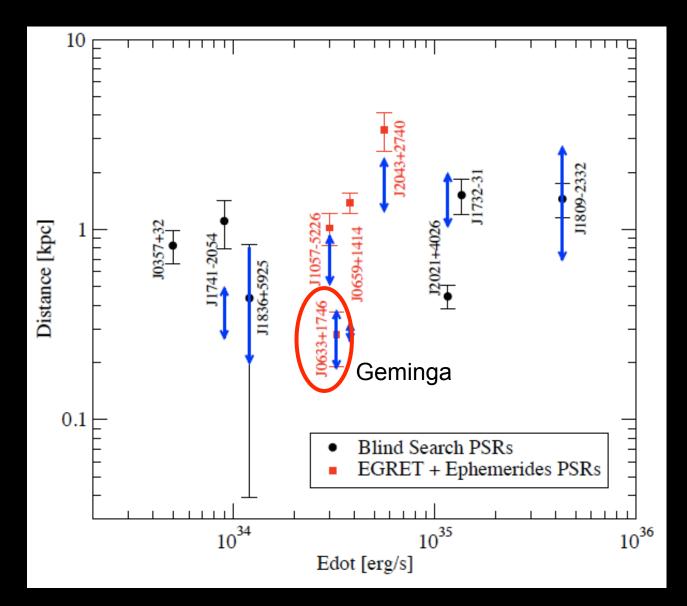
 f_{Ω} : flux correction factor – from Watters et al "Atlas"

 $f_{\Omega} = f_{\Omega}(\alpha, \zeta_E) = \frac{\iint F_{\gamma}(\alpha, \zeta, \phi) \sin(\zeta) \mathrm{d}\zeta \mathrm{d}\phi}{2 \int F_{\gamma}(\alpha, \zeta_E, \phi) \mathrm{d}\phi}.$

plus: standard error propagation



Distance Determinations: Results



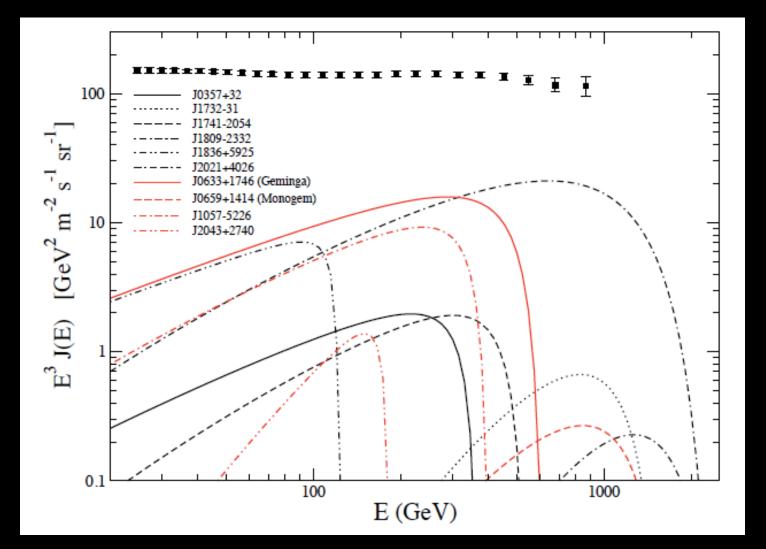
Contribution to the local e+e- Flux

$$Q(E, t, \vec{r}) = Q_0 \left(\frac{E}{1 \text{ GeV}}\right)^{-\Gamma} \exp[-E/E_{\text{cut}}]\delta(t - t_0)\delta(\vec{r})$$
$$E_{\text{cut}} = 1 \text{ TeV}, \qquad \Gamma = 1.7.$$
$$\int_{m_e}^{\infty} Q(E) dE = E_{\text{out}} = \eta \frac{\dot{E}t_{\text{ch}}^2}{\tau}, \quad \text{with } \tau \simeq 10^4 \text{ yr and } \eta = 0.4.$$

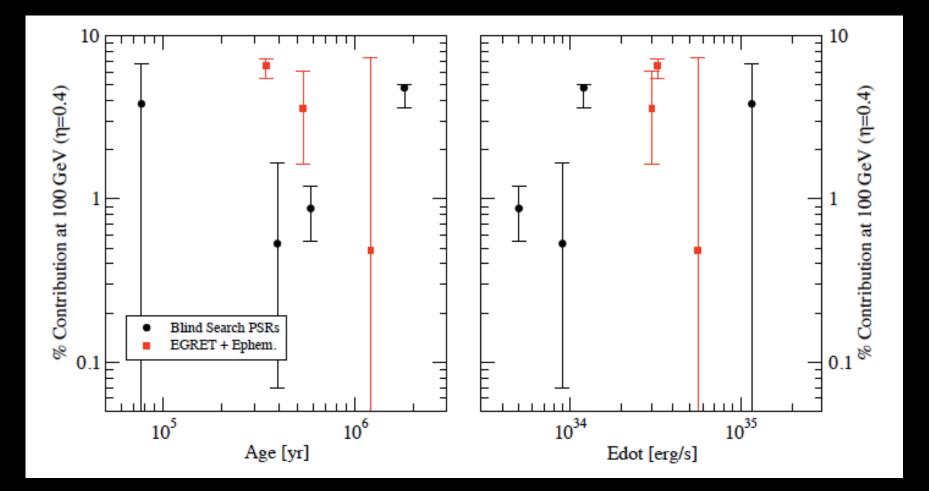
$$\frac{\partial N_e(E, t, \vec{r})}{\partial t} - D(E)\nabla^2 N_e - \frac{\partial}{\partial E}(b(E)N_e) = Q(E, t, \vec{r}),$$

$$D(E) = D_0(E/1 \text{ GeV})^{\delta} \qquad D_0 = 3.6 \times 10^{28} \text{ cm}^2/s \text{ and } \delta = 0.33$$

Contribution to the local e+e- Flux



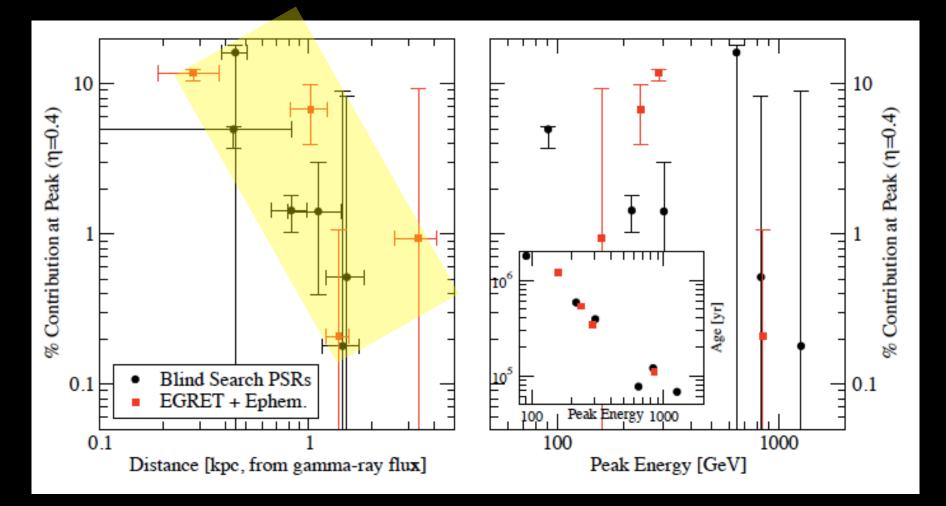
Contribution to the local e+e- Flux – a closer look (1)



Rather narrow age and spin-down luminosity ranges

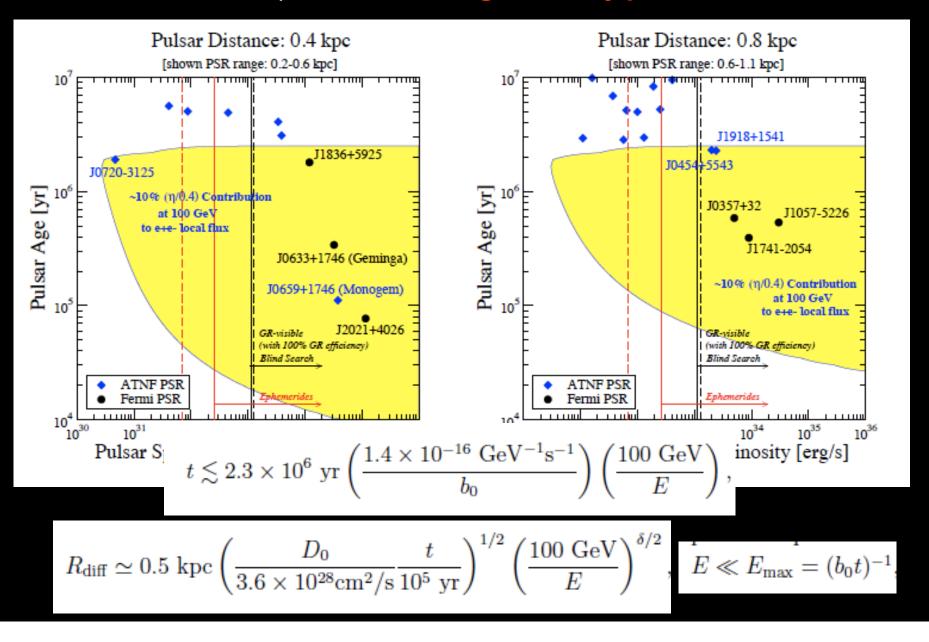
• flat contribution over those ranges

Contribution to the local e+e- Flux – a closer look (2)

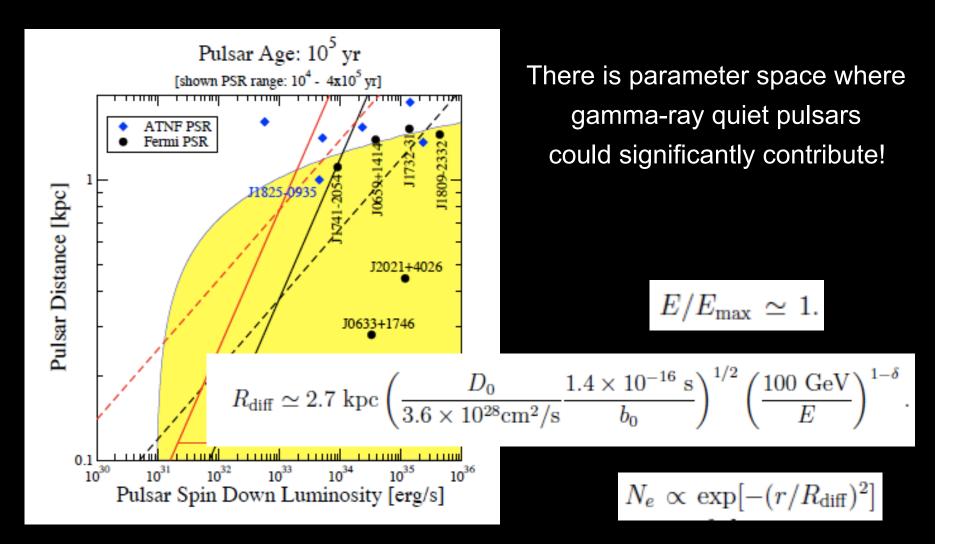


- Steeply declining contribution with distance
- Largest contributions between 100 and 1000 GeV

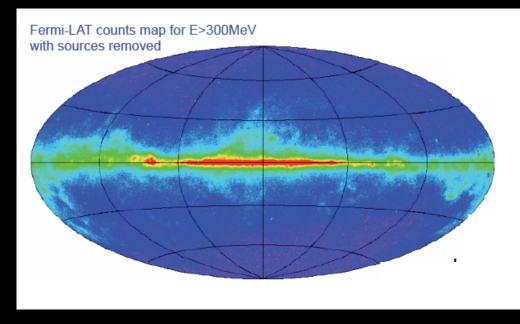
Do we expect known radio pulsar (~1,500) to be still more important than ~50 gamma-ray pulsars?



Do we expect known radio pulsar (~1,500) to be still more important than ~50 gamma-ray pulsars?



Some thoughts about the "Fermi haze"

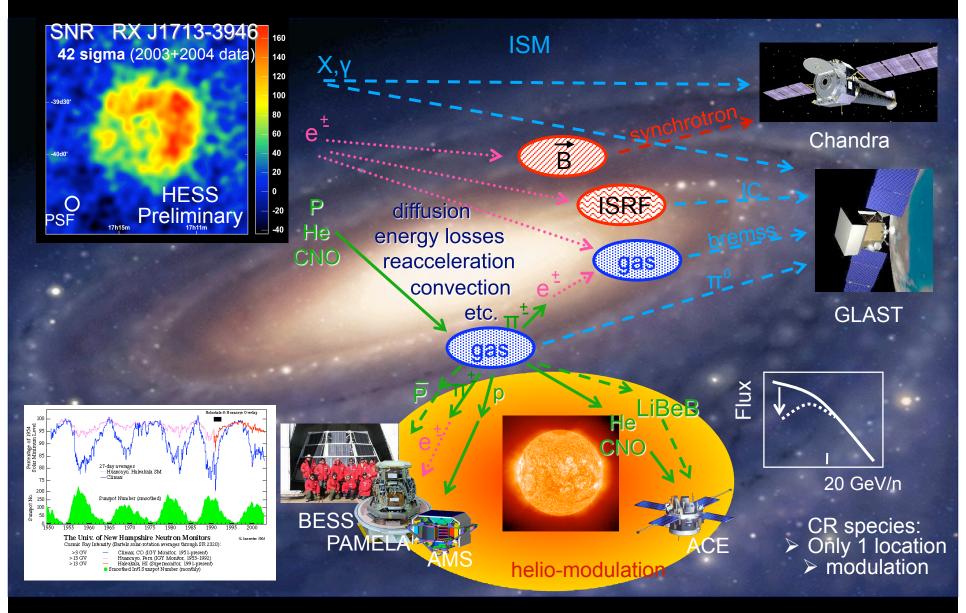


Three main contributions to the **diffuse** gamma-ray **emission**:

- Galactic emission
 - Neutral Pions inelastic hadronic cosmic-ray processes
 - e+e- Radiative Losses: Inverse Compton & Bremsstrahlung
- **Isotropic** emission (extra-galactic + cosmic ray mis-ID)

(*) J-M Casandjian, Fermi Symposium 2009

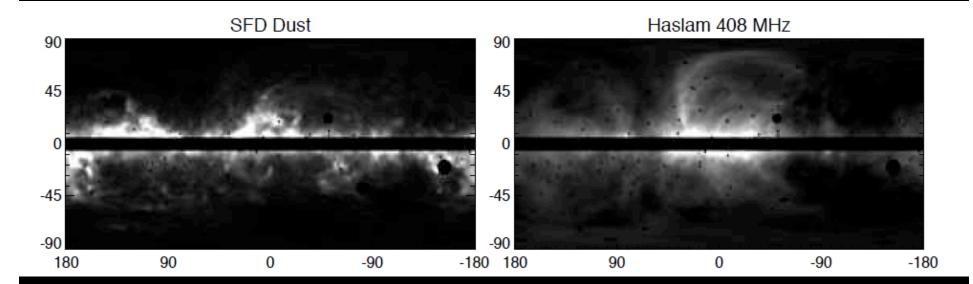
The standard full-glory approach



[slide from Igor Moskalenko]

Is there an **excess** gamma-ray diffuse **emission**? Relevant for: e+e- responsible for the haze, DM annihilation

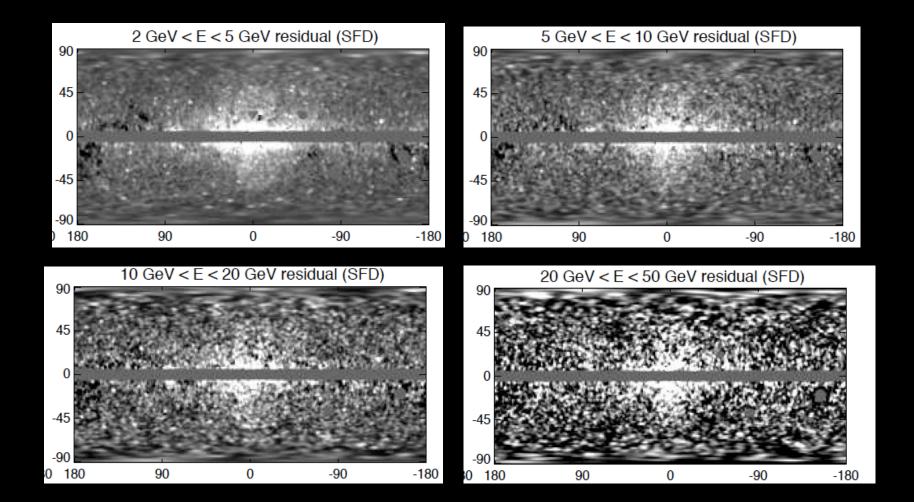
> Dobler et al: simple diagnostic: fit **spatial templates** to Fermi gamma-ray sky



Dust (far IR 100 μ emission) => ISM Gas => π^0 (no CR sources!) Synchrotron emission map

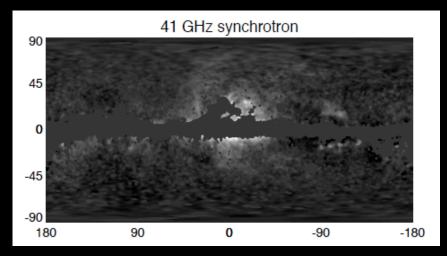
=> IC emission (same e+e-, but...)

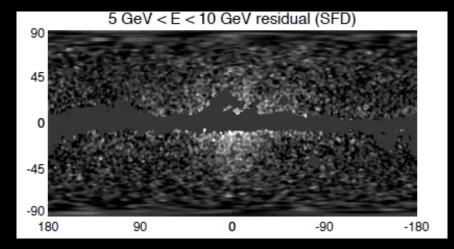
Large **residuals** in a region morphologically comparable to the microwave **WMAP** haze excess



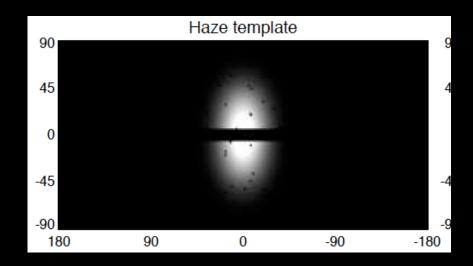
(*) Dobler et al, 0910.4583

The resulting "Fermi haze" has an intriguing morphology

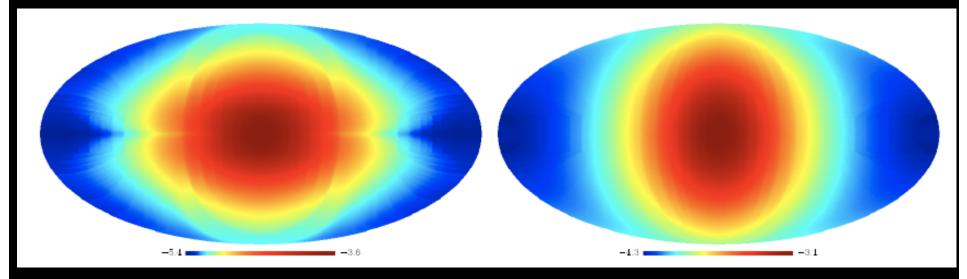




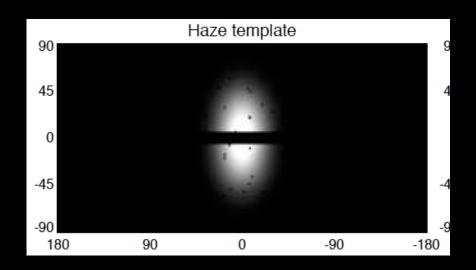
Try to fit with a **bi-variate gaussian** additional diffuse component (expected from e.g. e+e- producing WMAP haze, or DM secondaries)

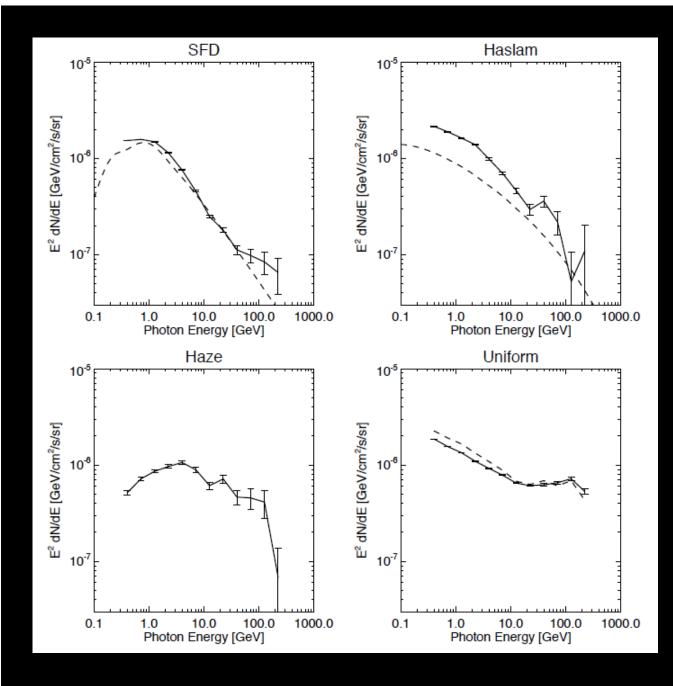


The resulting "Fermi haze" has an intriguing morphology



Try to fit with a **bi-variate gaussian** additional diffuse component (expected from e.g. e+e- producing WMAP haze, or DM secondaries)





Residuals are now at a "noise" level – satisfactory fit!

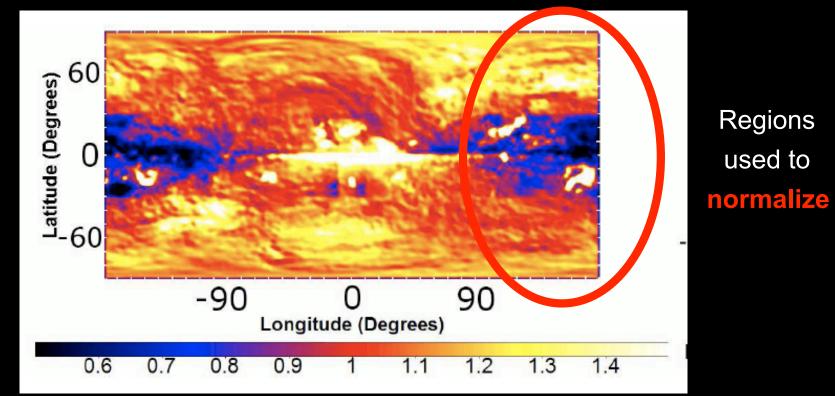
Proceed to use the relative normalization to calculate **spectra**

(*) Dobler et al, 0910.4583

Are spatial templates accurate enough to claim a Fermi haze?

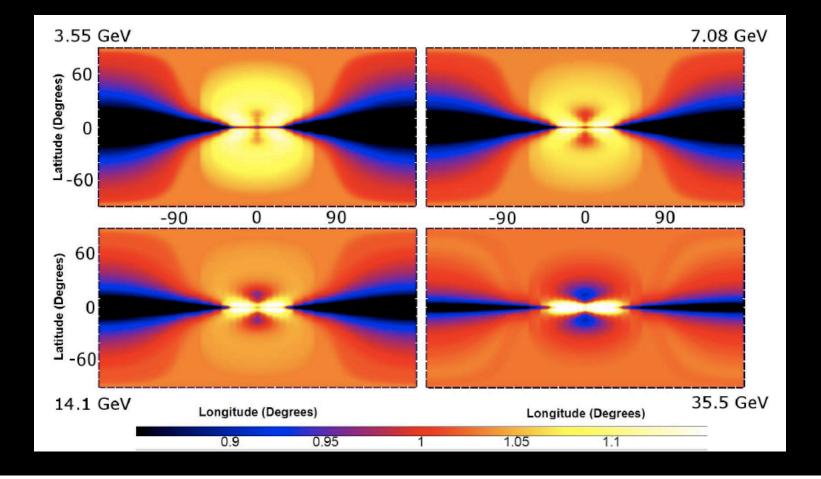
Test quantitatively by employing the predictions of the full-glory approach

First template - First assumption: gas (dust) traces π^{0} 's



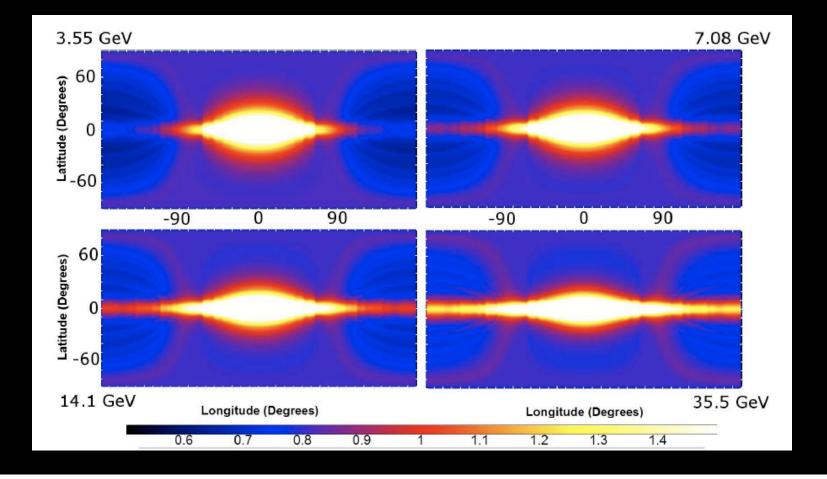
Map (norm. to 1) of ratio of π^0 's to **ISM** col. dens. Issue: **cosmic rays sources** live in the haze region! **Second Template - Second Assumption**: Synchrotron traces IC

Issue: while both are sourced by e+e-, the emissions follow the magnetic field and inter-stellar radiation field energy densities with potentially very different morphologies!



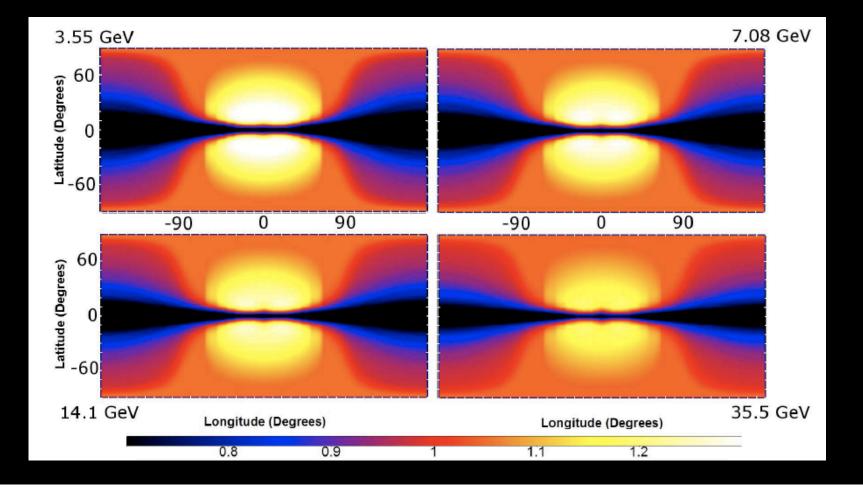
Second Template - Second Assumption: Synchrotron traces IC

Try a **constant** magnetic field model (retaining a detailed **ISRF** model) again, produce an artificial haze



Second Template - Second Assumption: Synchrotron traces IC

Assuming a "sharper" galactic magnetic field model (smaller z scale) one produces again an (energy-dependent) artificial haze



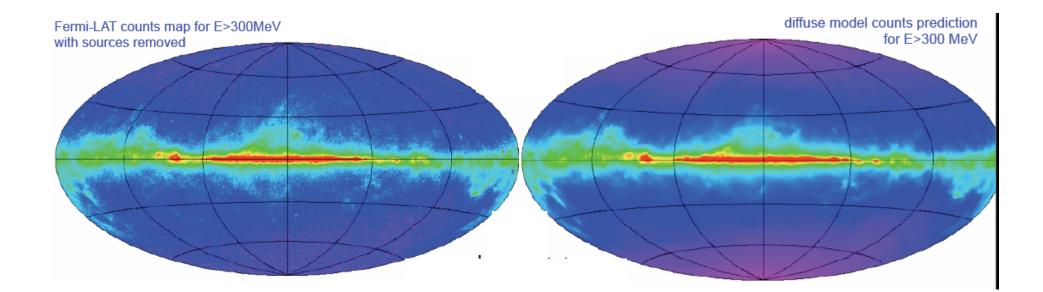
Systematic effects in the spatial templates artificially produce a haze and affect the low-energy determination of an excess

At large energies: very low statistics + very large cosmic ray contamination

A full-glory galactic cosmic ray simulation is necessary

but is it **sufficient**?!

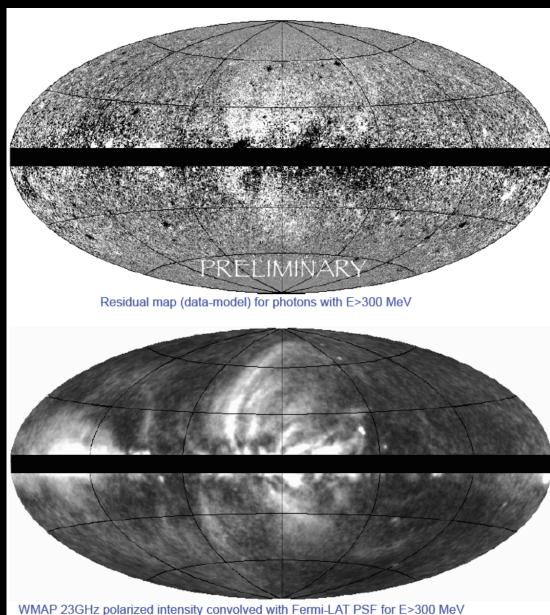
Fermi gamma-ray sky (source subtracted) and Galprop predictions (from Casandjian, poster @ Fermi Symposium, 2009)



(*) J-M Casandjian, Fermi Symposium 2009

The residuals trace the polarized synchrotron emission presumably from the nearby Loop I (North Polar Spur) supernova shell (~100 deg!!)

local cosmic ray structure maybe key to understand the diffuse emission beyond galactic cosmic ray simulations and models



(*) J-M Casandjian, Fermi Symposium 2009

Unfortunately charged cosmic rays don't travel straight lines and they are sometimes not produced by our favorite sources... Unfortunately charged cosmic rays don't travel straight lines and they are sometimes not produced by our favorite sources...

Yet, the field is in the midst of a unique and exciting spur of data

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A solid prediction: **boredom not in sight for theorists**!