

# Kaluza-Klein Dark matter

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Focus week on indirect Dark matter search

Dec 09, 2009

- SCP, Shu, Phys. Rev. D 79, 091702(*Rapid*) (2009)
- Chen, Nojiri, SCP, Shu, Takeuchi, JHEP 0909.078 (2009)
- Chen, Nojiri, SCP, Shu, Takeuchi, [arXiv 0908.4317]
- Csaki, Heinonen, Hubsiz, SCP, Shu, in preparation



# contents

- Dark matter problem in particle physics
- Kaluza-Klein Dark matter in UED & split-UED
- Fitting Pamela+Fermi excesses



# Dark matter problem

- WMAP data implies :  $DM \gg Baryon$
- Chandra-X observation of bullet cluster: DM is made of matter (not MOND)
- No candidate in the SM



# WIMP is a wonderful solution

- Consistent with the standard Big-Bang thermal history
- $O(100)$  GeV, weakly interacting stable massive particle can nicely fit WMAP result.

[Lee-Weinberg, PRL 1977]



# Stability of WIMP

- Stability of a particle can be guaranteed by a symmetry
- The simplest choice is Parity ( $Z_2$ )
- Examples: R-parity in MSSM, T-parity in Little Higgs, KK-parity in extra dimension, possibly many more



# Examples of WIMP candidates

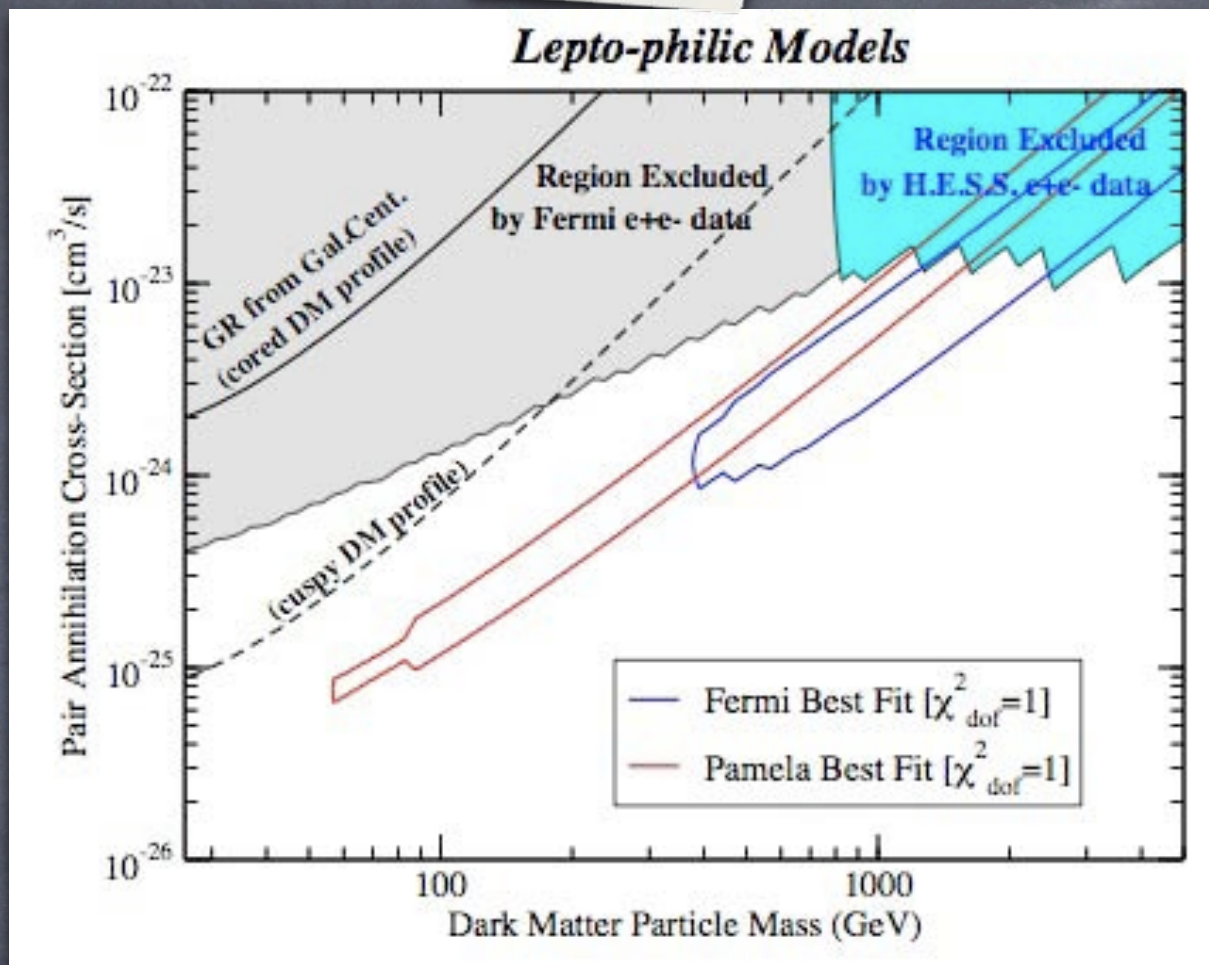
- LSP in MSSM with R-parity (neutralino,  $s=1/2$ , Majorana particle)
- LTP in Little Higgs with T-parity ( $A_H$ ,  $s=1$ ,  $m < 350$  GeV)
- LKP in extra dimension with KK-parity ( $B_1$ ,  $s=1$ , couples to hypercharge)
- possibly many more..



# Info from recent CR-data

- If DM is responsible for Pamela, ATIC, Fermi “excesses”, we can learn more about the details of DM properties.
- Pamela  $p^-/p$  and Fermi gamma-ray data seem to imply leptophilic DM. (neutralino disfavored, KK dm favored)
- Mass of DM can be determined by the “peak” position (ATIC=> 600–800 GeV, Fermi=>900 GeV or larger) (LTP disfavored, KK dm favored)





[Fermi-Lat: arXiv 0905.0636]

**DM annihilation** into charged leptons can explain the Fermi and Pamela if  $m > 400 \text{ GeV}$

\*Tesla Jeltema this workshop showed that non-detection of clusters (end EGB) by Fermi-LAT "excludes" high mass models with  $m > 1\text{--}2 \text{ TeV}$  fitting the Pamela positron excess.

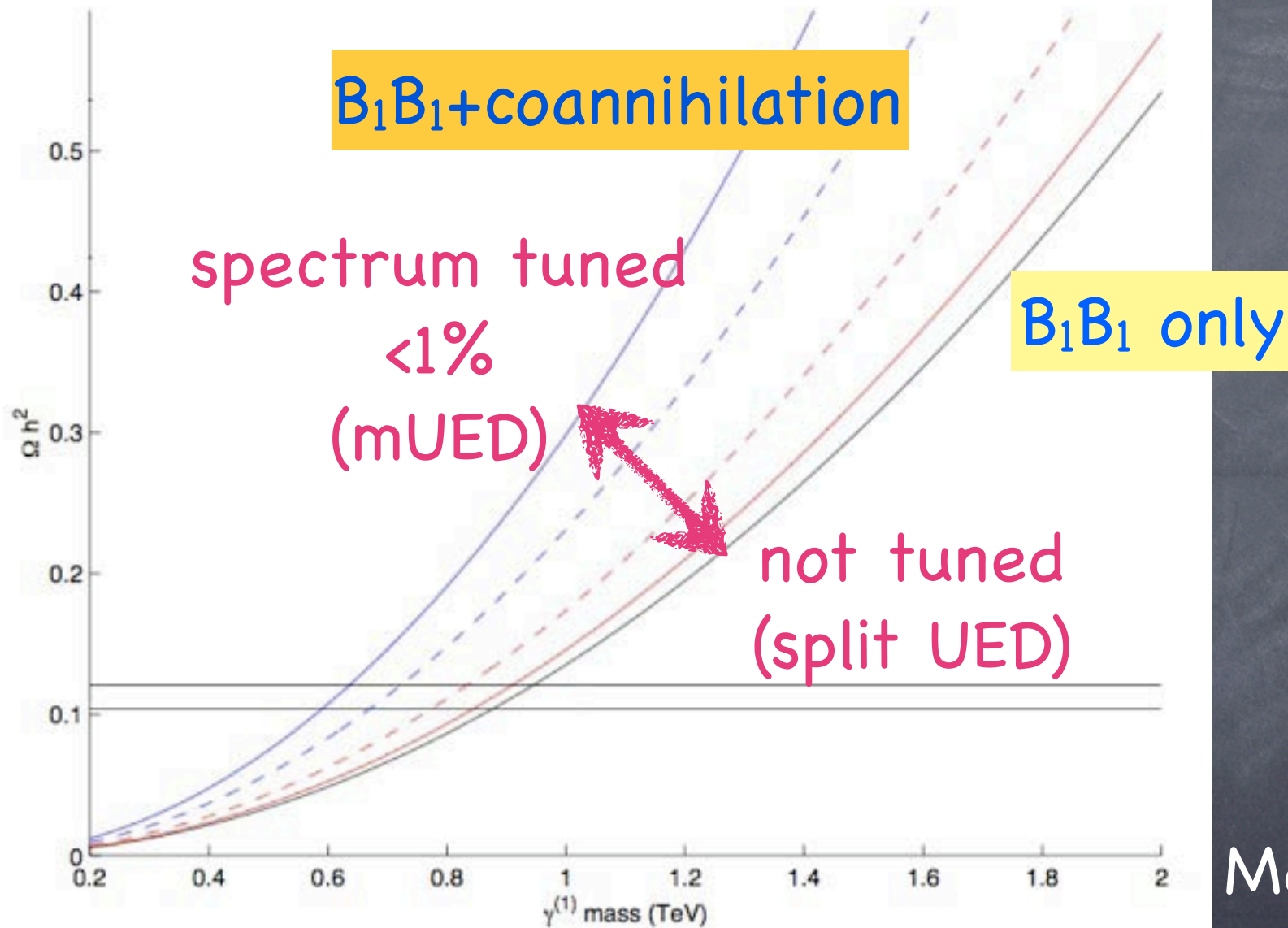


Relic  
density

big



small



Mass

**The Abundance of Kaluza-Klein dark matter with coannihilation.**

Fiona Burnell, (Princeton U.) , Graham D. Kribs, (Oregon U.) . Sep 2005. 38pp.

Published in **Phys.Rev.D73:015001,2006.**

e-Print: **hep-ph/0509118**

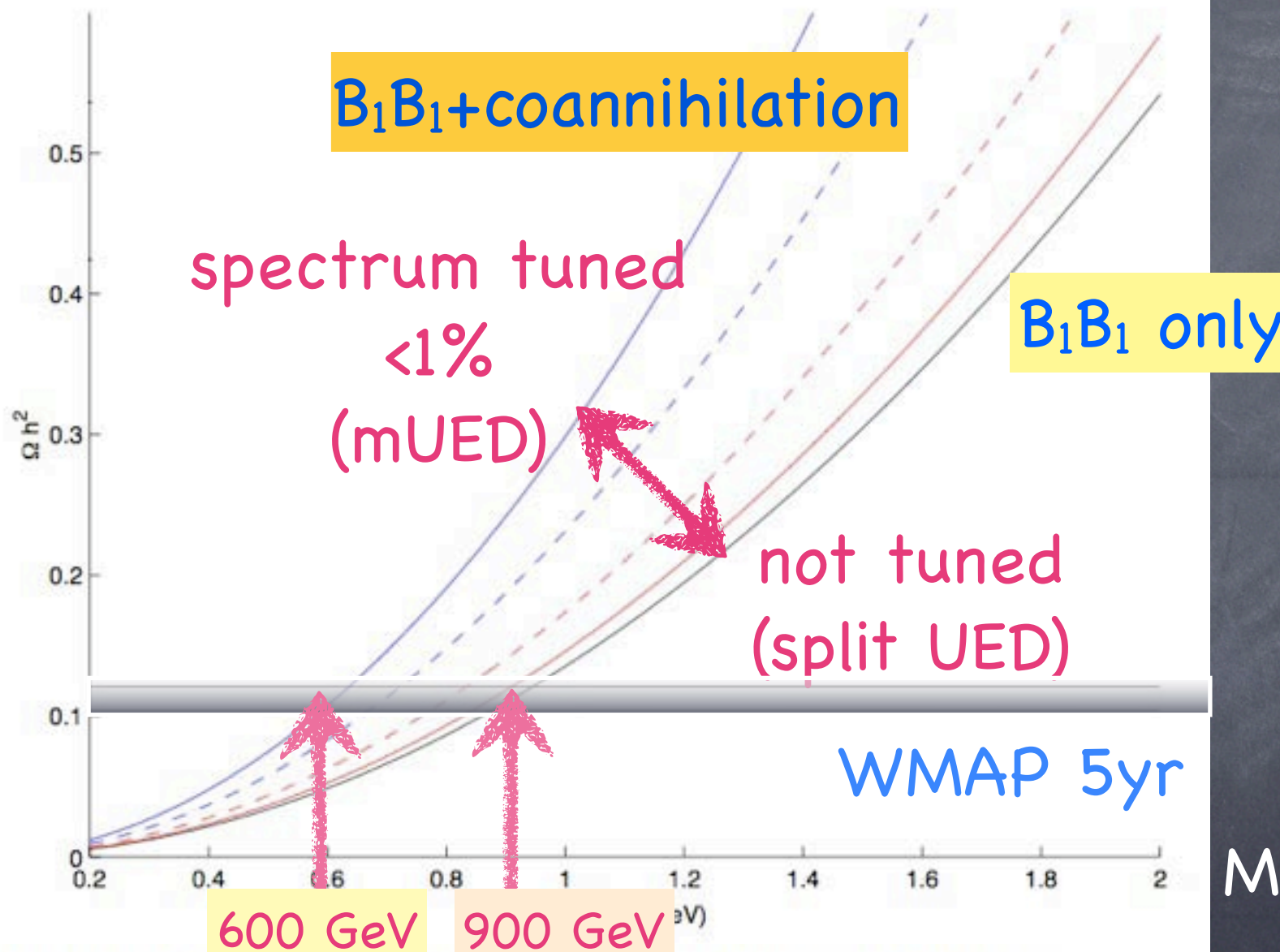


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e-Print: [hep-ph/0509118](http://arxiv.org/abs/hep-ph/0509118)



# Kaluza-Klein DM in UED & split-UED



# Split-UED

- A model with the minimal extension of spacetime whose low energy effective theory is nothing but the SM.  $D=5$ ,  $M^4 \times S^1/Z_2$ ,  $G=SU(3) \times SU(2) \times U(1)$
- **All matters in the bulk.** Matter=Q (3,2,1/6),  $U^c$  (3,1,-2/3),  $D^c$  (3,1,1/3), L(1,2,-1/2),  $E^c$ (1,1,1)
- Bulk Dirac Mass term introduced. ( $m \rightarrow 0$  limit corresponds conventional “UED”)



# Theory space

5D orbifold gauge theory  
 $G \supseteq SU(3) \times SU(2) \times U(1)$   
 $\Psi \ni (Q, U^c, D^c, L, E^c)$

split-UED  
arbitrary  $m_5$

mUED  
 $m_5=0$



# Bulk Action

the most generic, lowest order 5D-Lorentz & gauge invariant action

$$\mathcal{L}_5^{UED} = \sum_{ij} \frac{i}{2} (D_M \bar{\Psi}_i \Gamma^M \Psi_j - \bar{\Psi}_i \Gamma^M D_M \Psi_j) \\ + \mathcal{L}_H + \mathcal{L}_{g,W,B}$$

$$\Psi = (Q, U^c, D^c, L, E^c)$$

$$\Gamma^M = (\gamma^\mu, i\gamma_5)$$

$$D = \partial - ig_3 \lambda G - ig_2 \frac{T}{2} \cdot W - ig_1 \frac{Y}{2} B$$

without loss  
of generality  $m_{ij} = \mu_i \delta_{ij}$     5 mu's for each gen.



# Bulk Action

the most generic, lowest order 5D-Lorentz & gauge invariant action

$$\text{split-}\mathcal{L}_5^{UED} = \sum_{ij} \frac{i}{2} (D_M \bar{\Psi}_i \Gamma^M \Psi_j - \bar{\Psi}_i \Gamma^M D_M \Psi_j) \\ - m_{ij}(y) \bar{\Psi}_i \Psi_j + \mathcal{L}_H + \mathcal{L}_{g,W,B}$$

$$\Psi = (Q, U^c, D^c, L, E^c)$$

$$\Gamma^M = (\gamma^\mu, i\gamma_5)$$

$$D = \partial - ig_3 \lambda G - ig_2 \frac{T}{2} \cdot W - ig_1 \frac{Y}{2} B$$

without loss  
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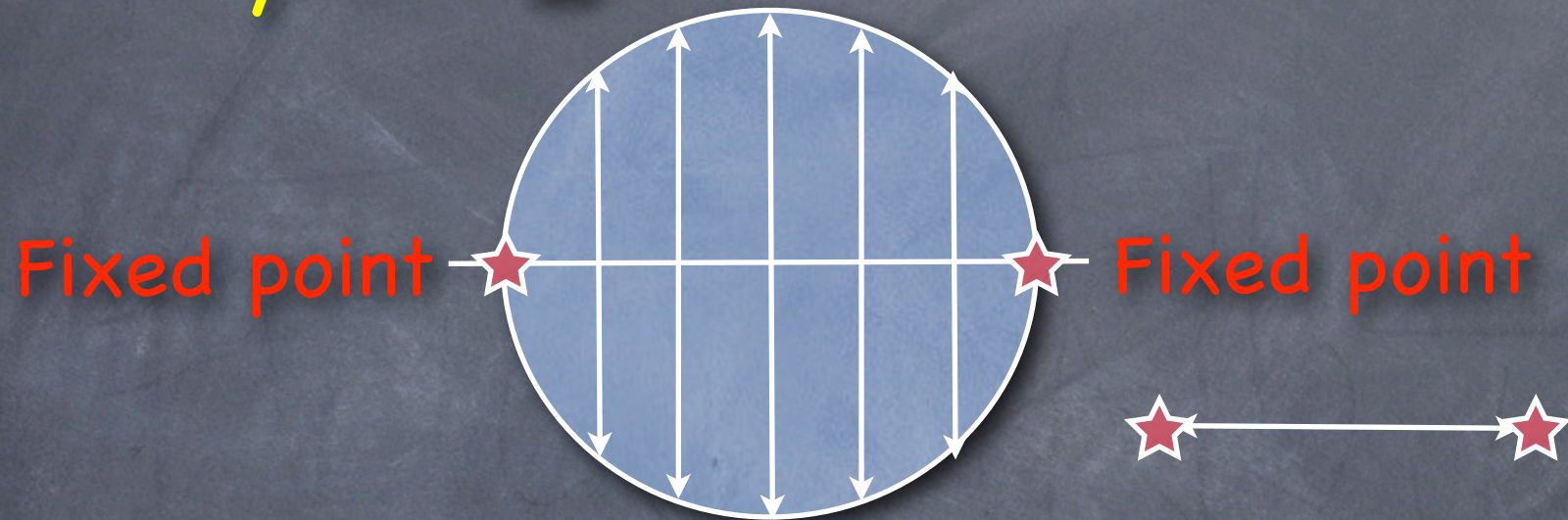


# Property #1: Orbifold

- The theory (5D) is **vectorlike**, but the SM (4D) is a **chiral theory**. The best thing we hope is that **a chiral theory appears as a low energy effective theory** below compactification scale.
- **Orbifold compactification** is a good example. Unnecessary chiral states are projected out by orbifold BC.
- The simplest orbifold:  $S^1/Z_2$



$$S^1 / Z_2$$



- $Z_2$ : Identify the opposite points, there are two fixed points
- Equivalent to an interval with two boundaries
- Orbifold condition determines BCs for fields



On orbifold  $S^1/Z_2$ , a field has a “ $Z=\pm 1$ ” parity

$$\Psi_Z(x, y_f) \rightarrow \Psi_Z(x, y_f - \epsilon) = Z\gamma_5 \Psi_Z(x, y_f + \epsilon)$$

$$\Psi_{L/R} = \sum_n \psi_{L/R}^n(x) f_{L/R}^n(y)$$

$$\Psi_Z(y_f - \epsilon) = Z\gamma_5 \sum_n \psi_L^n f_L^n(y_f + \epsilon) + \psi_R^n f_R^n(y_f + \epsilon)$$

$$= Z \sum_n -\psi_L^n f_L^n(y_f + \epsilon) + \psi_R^n f_R^n(y_f + \epsilon)$$

$f_L$  and  $f_R$  have  
the opposite  
parities.

$$f_L^n(y_f - \epsilon) = -Z f_L^n(y_f + \epsilon)$$

$$f_R^n(y_f - \epsilon) = +Z f_R^n(y_f + \epsilon)$$

If one of them is odd (satisfying Dirichlet BC)  
, the other one is even (satisfying Neumann BC)



Neumann BCs are imposed on:

$$H, G_{\mu}^A, W_{\mu}^a, B_{\mu}, Q_L, L_L, U_R, D_R, E_R$$

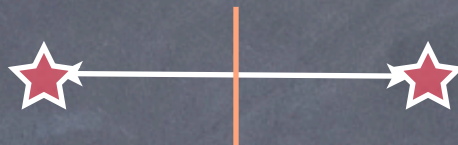
Dirichlet BCs are imposed on:

$$G_5^A, W_5^a, B_5, Q_R, L_R, U_L, D_L, E_L$$

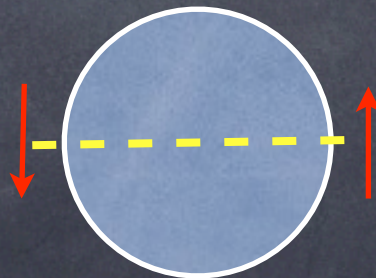
\*The resultant zero mode spectrum is the same as the SM.



# Property #2: KK-parity

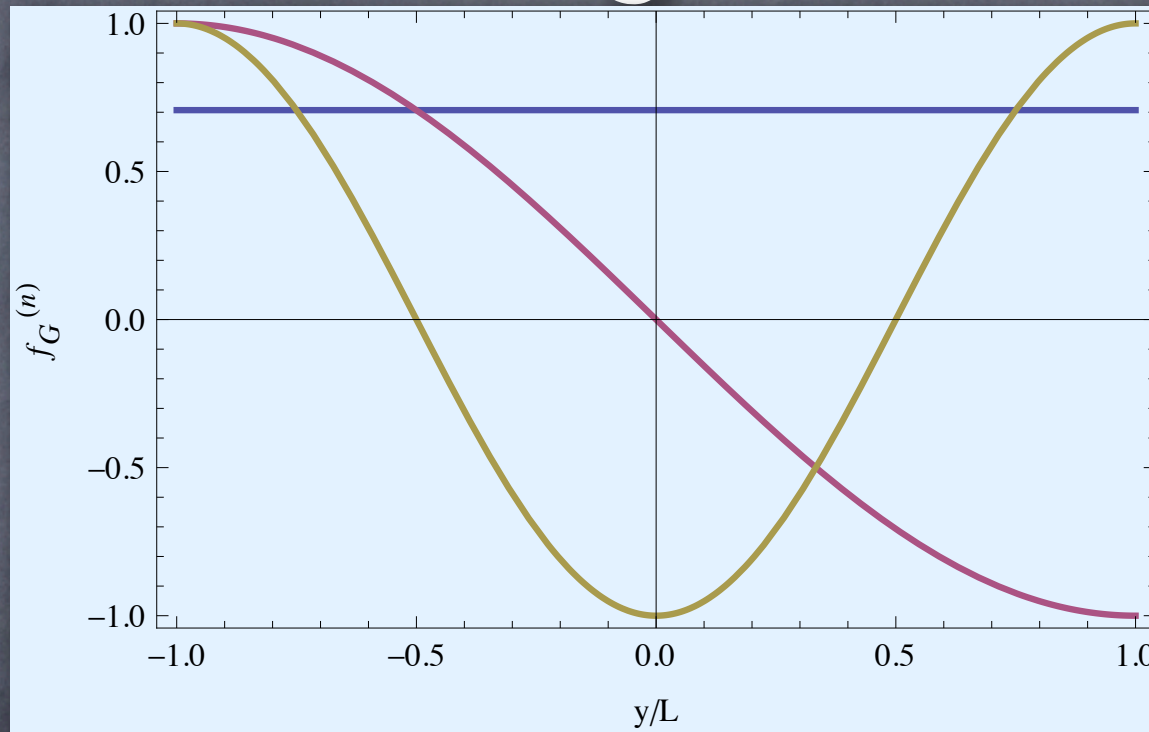


- The reflection **symmetry about the mid point** of extra dimension.
- **Remnant symmetry of KK-number conservation** (=momentum conservation along 5th direction coming from translational invariance which is broken by **fixed points**)





# e.g. Wave function of KK gluons



1. Gluon respects Neumann BCs.  $\Rightarrow$  have zero mode
2.  $n=\text{even}$  Kk modes are symmetric (parity even)  
 $n=\text{odd}$  KK modes are antisymmetric (parity odd)



# Property #3:

## Bulk mass

No "chirality" in odd dim.

- The minimal spinor in 5D is Dirac-like spinor

$$\Psi = \Psi_L(1/2, 0) + \Psi_R(0, 1/2)$$

- Dirac mass is generically allowed. No symmetry forbids it. This mass is nothing to do with the electroweak symmetry breaking. Zero mode remains massless regardless this mass:

$$m \bar{\Psi} \Psi$$



- In split-UED, we introduce **an odd mass** so that **KK-parity is respected**.

**Inversion** about the middle point

$$\begin{array}{ccc} \star & \xrightarrow{y} & \star \\ -L/2 & & L/2 \end{array} \quad S^1/Z_2$$

$$\Psi(x, y) \rightarrow \Psi(x, -y) = \pm \gamma_5 \Psi(x, y)$$

$$\bar{\Psi} \Psi \rightarrow -\bar{\Psi} \Psi$$

Thus, we need an intrinsically “odd” mass term for keeping KK-parity.

$$M_5(y) \bar{\Psi} \Psi$$

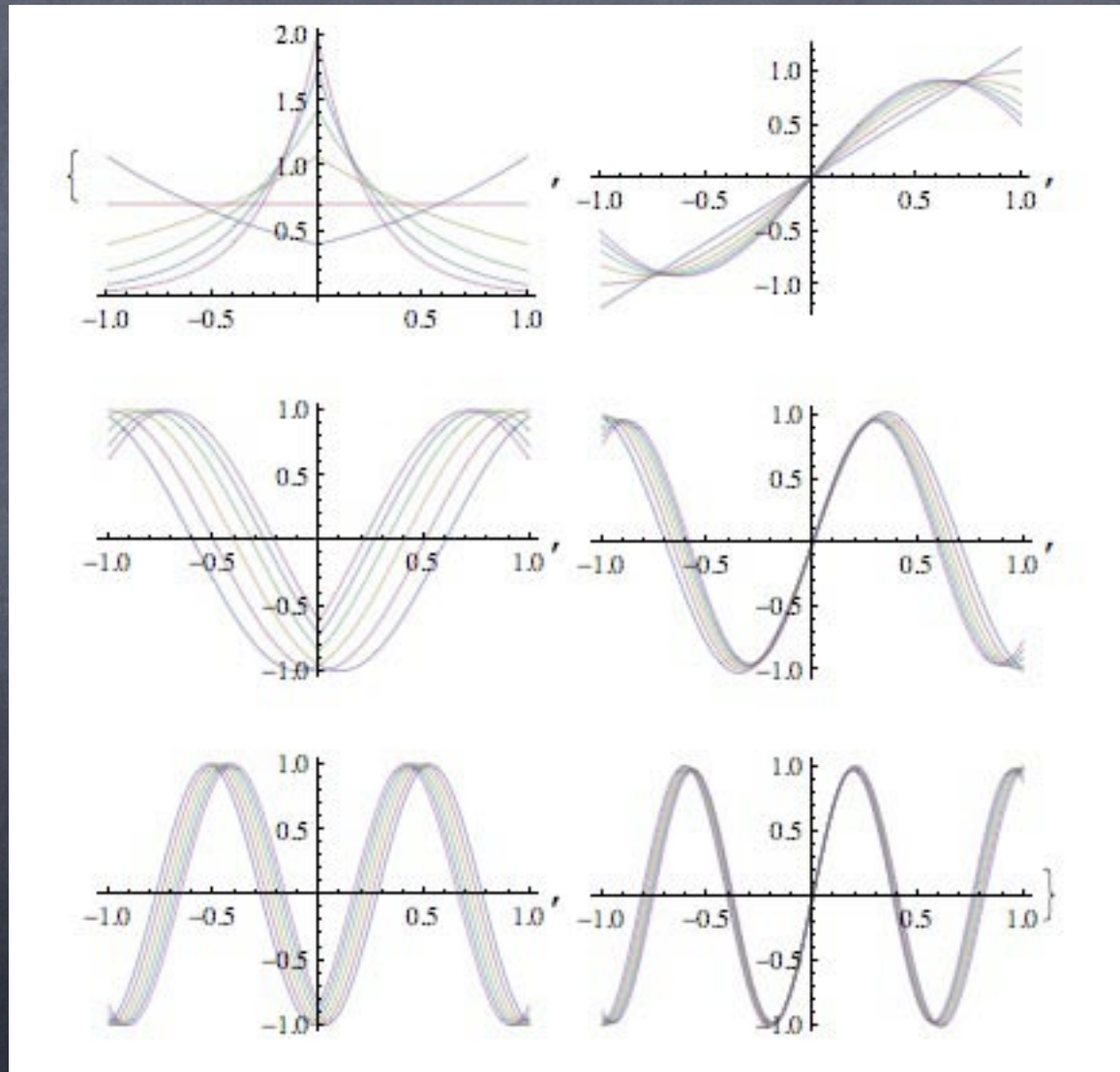
$$M_5(y) \rightarrow M_5(-y) = -M_5(y)$$

The simplest  
choice:

$$m(y) = m \theta(y) \approx m \tanh(\mu y)$$



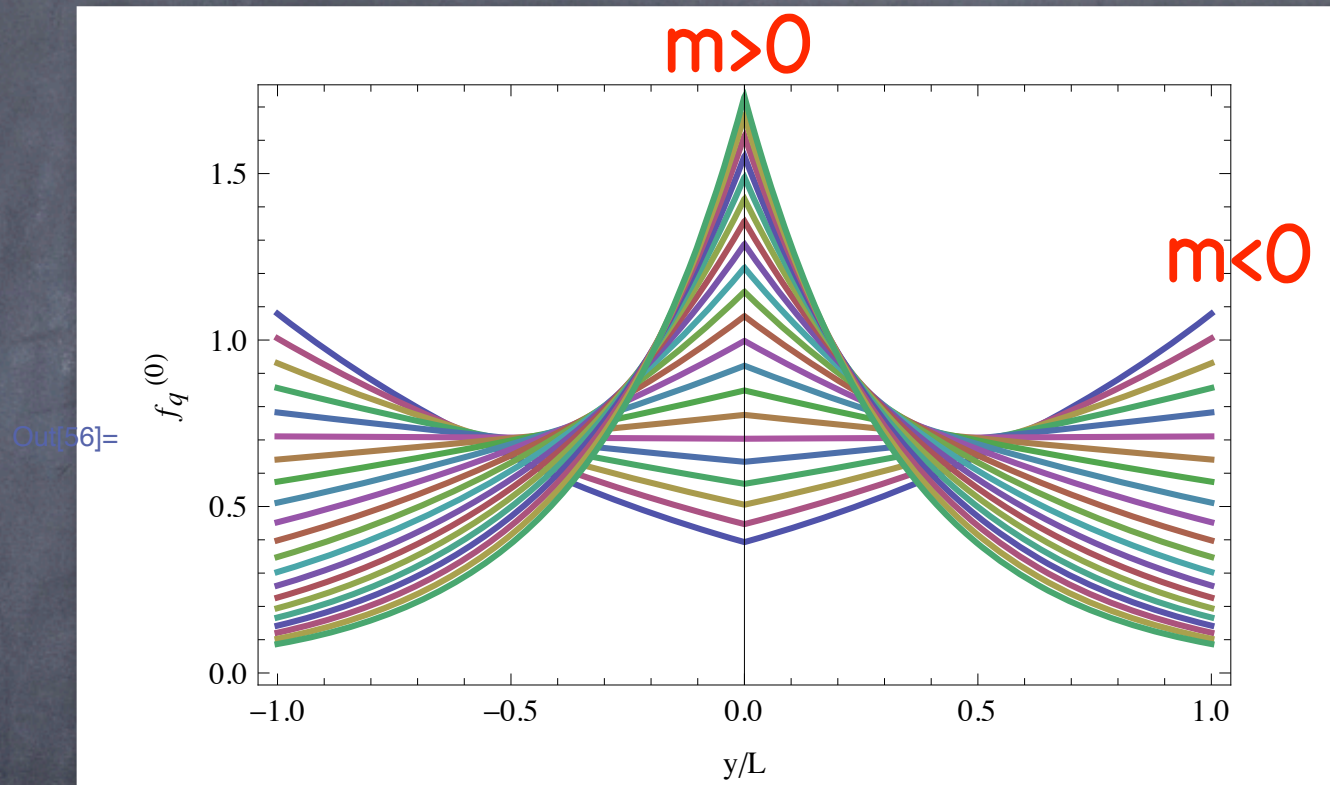
# Wave functions for KK modes





Zero mode  
:all massless

$$m(y) = m\theta(y)$$



$$(\partial_y + m(y))f_R^0 = 0$$

$$f_R^0 \sim e^{-\int dy m(y)} = e^{-m|y|}$$



# KK-gauge coupling

$$\begin{aligned}\mathcal{L}_{\text{int}} &\ni \int dy g_5 \bar{\psi} T^a \gamma^\mu \psi G_\mu^a \\ &\ni \sum_{nlm} g^{eff, nlm} \bar{\psi}^n T^a \gamma_\mu \psi^l G_\mu^{ma} \\ g_n^{eff, nlm} &= g_5 \int dy f_{\bar{\psi}}^n f_\psi^l f_G^m\end{aligned}$$

The integral of the **wave function overlap** gives the effective coupling.



SM-SM-SM vertex is allowed  
(i.e. even-even-even)

even

even

$$g \propto \int_{-L}^L dy \psi_{\text{even}} \psi_{\text{even}} \psi_{\text{even}} \neq 0$$

even

- 👁 KK-even particle has a **symmetric** wave fn.
- 👁 SM particles are zero modes (i.e. KK-even)



SM-SM-DM vertex is **NOT** allowed

odd

even

$$g \propto \int_{-L}^L dy \psi_{\text{odd}} \psi_{\text{even}} \psi_{\text{even}} = 0$$

Chen, Nojiri, Park, Shu, Takeuchi (2009)

Chen, Nojiri, Park, Shu (2009)

even

- LKP is stable  $\Rightarrow$  Dark Matter
- Never produced singly

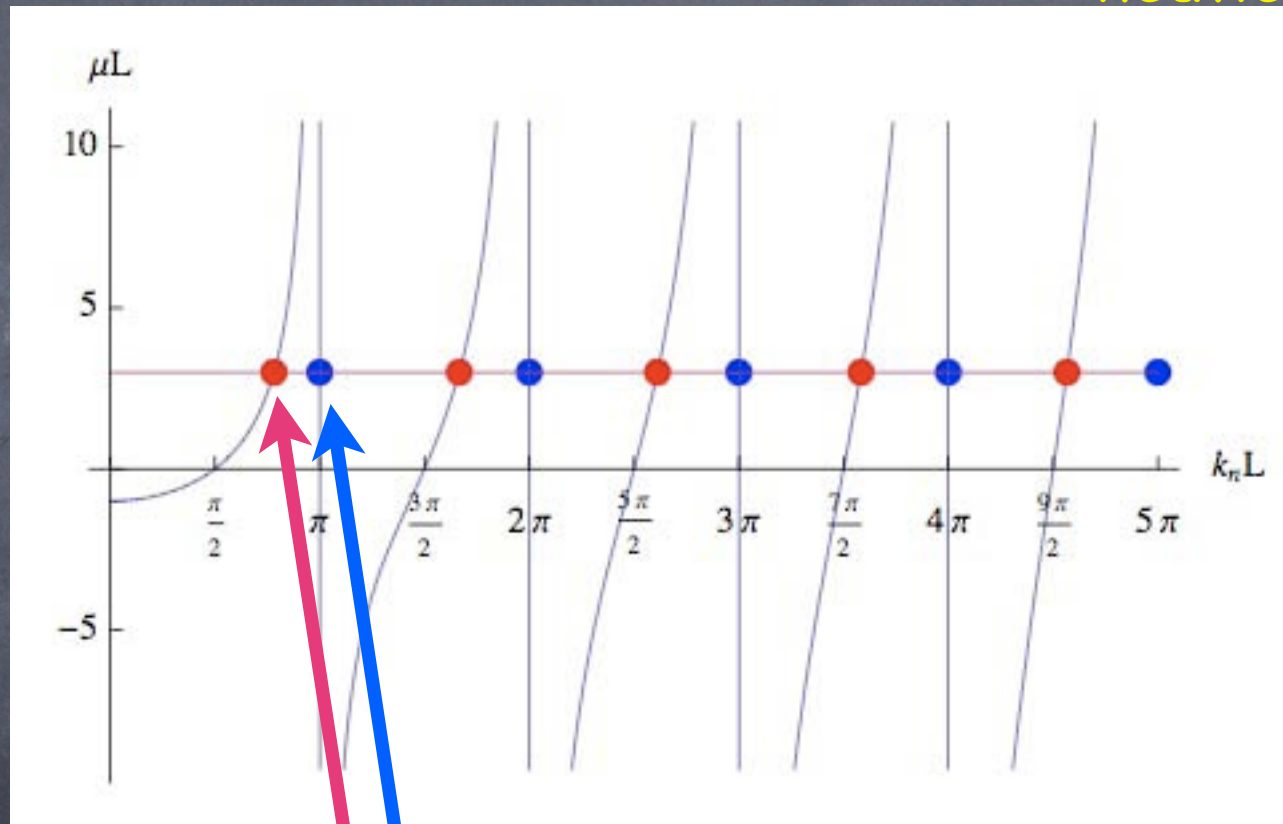
“KK parity mimics R-parity”



# Kaluza-Klein spectrum (Tree)

$$m_n^2 = \mu^2 + k_n^2$$

Fermions get  
heavier with  $m_5$



[Csaki, Hainonen, Hubsiz, Park, Shu (in preparation)]

$$\mu = -k_n \cot k_n L \text{ (odd)}$$

$$k_n = n\pi/L \text{ (even)}$$



# KK photon (=B<sub>1</sub>) is LKP

$$\delta(m_{B(n)}^2) = \frac{g'^2}{16\pi^2 R^2} \left( \frac{-39}{2} \frac{\zeta(3)}{\pi^2} - \frac{n^2}{3} \ln \Lambda R \right)$$

$$\delta(m_{W(n)}^2) = \frac{g^2}{16\pi^2 R^2} \left( \frac{-5}{2} \frac{\zeta(3)}{\pi^2} + 15n^2 \ln \Lambda R \right)$$

$$\delta(m_{g(n)}^2) = \frac{g_3^2}{16\pi^2 R^2} \left( \frac{-3}{2} \frac{\zeta(3)}{\pi^2} + 23n^2 \ln \Lambda R \right)$$

$$\delta(m_{Q(n)}^2) = \frac{n}{16\pi^2 R} \left( 6g_3^2 + \frac{27}{8}g^2 + \frac{1}{8}g'^2 \right) \ln \Lambda R$$

$$\delta(m_{u(n)}^2) = \frac{n}{16\pi^2 R} (6g_3^2 + 2g'^2) \ln \Lambda R$$

$$\delta(m_{d(n)}^2) = \frac{n}{16\pi^2 R} \left( 6g_3^2 + \frac{1}{2}g'^2 \right) \ln \Lambda R$$

$$\delta(m_{L(n)}^2) = \frac{n}{16\pi^2 R} \left( \frac{27}{8}g^2 + \frac{9}{8}g'^2 \right) \ln \Lambda R$$

$$\delta(m_{e(n)}^2) = \frac{n}{16\pi^2 R} \frac{9}{2} g'^2 \ln \Lambda R.$$

1-loop RG-running

1. KK photon -0.2%

2. KK gluon +30%

3. KK quarks +14%

4. KK leptons +1%

**Bosonic supersymmetry? Getting fooled at the CERN LHC.**

Hsin-Chia Cheng, (Chicago U., EFI), Konstantin T. Matchev, (Florida U. & CERN), Martin Schmaltz, (Boston U.)

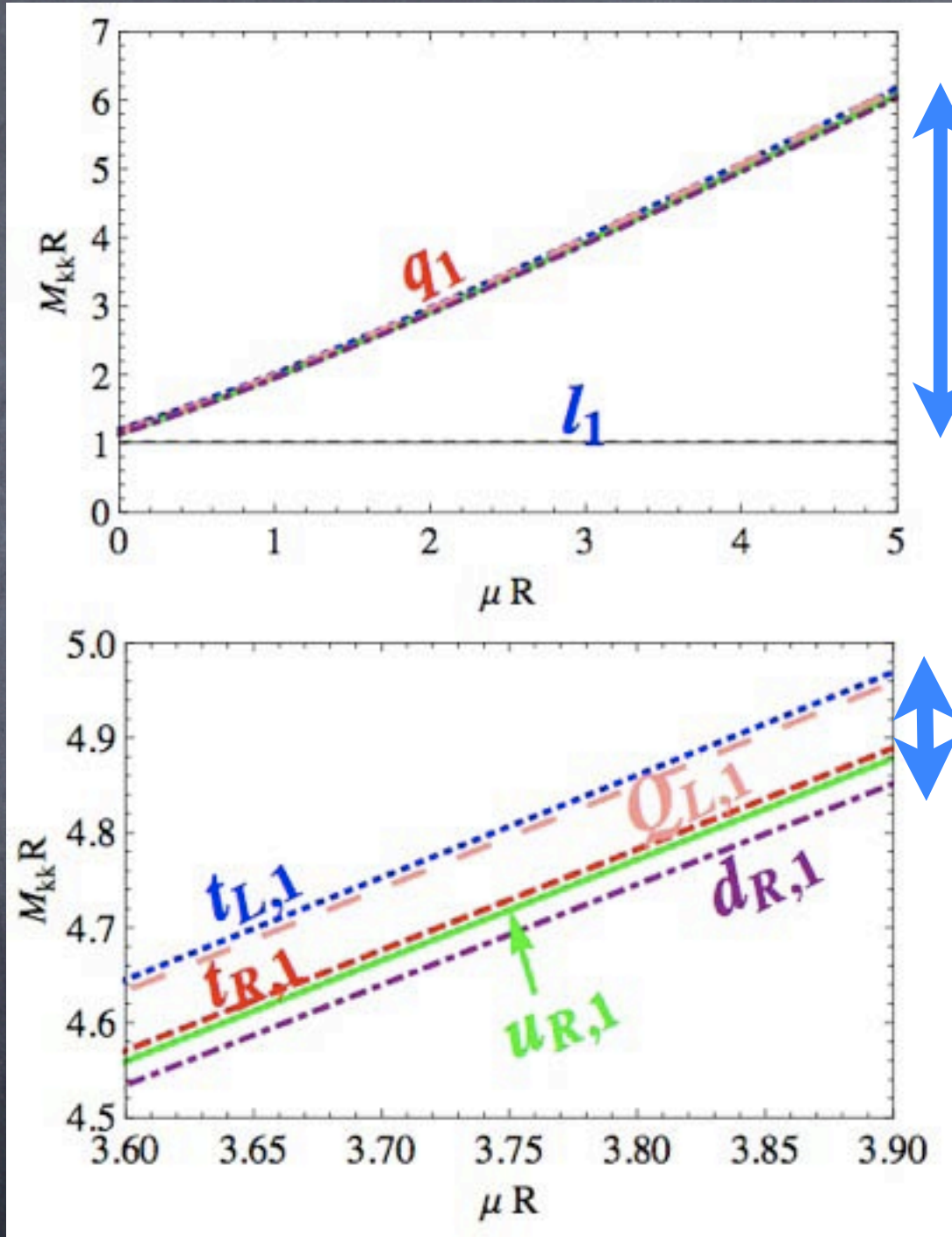
CERN-TH-2002-108, BUHEP-02-21, May 2002. 6pp.

Published in **Phys.Rev.D66:056006,2002.**

e-Print: [hep-ph/0205314](http://hep-ph/0205314)



bulk mass  
for  
quarks  
(1-loop  
RG-  
running)



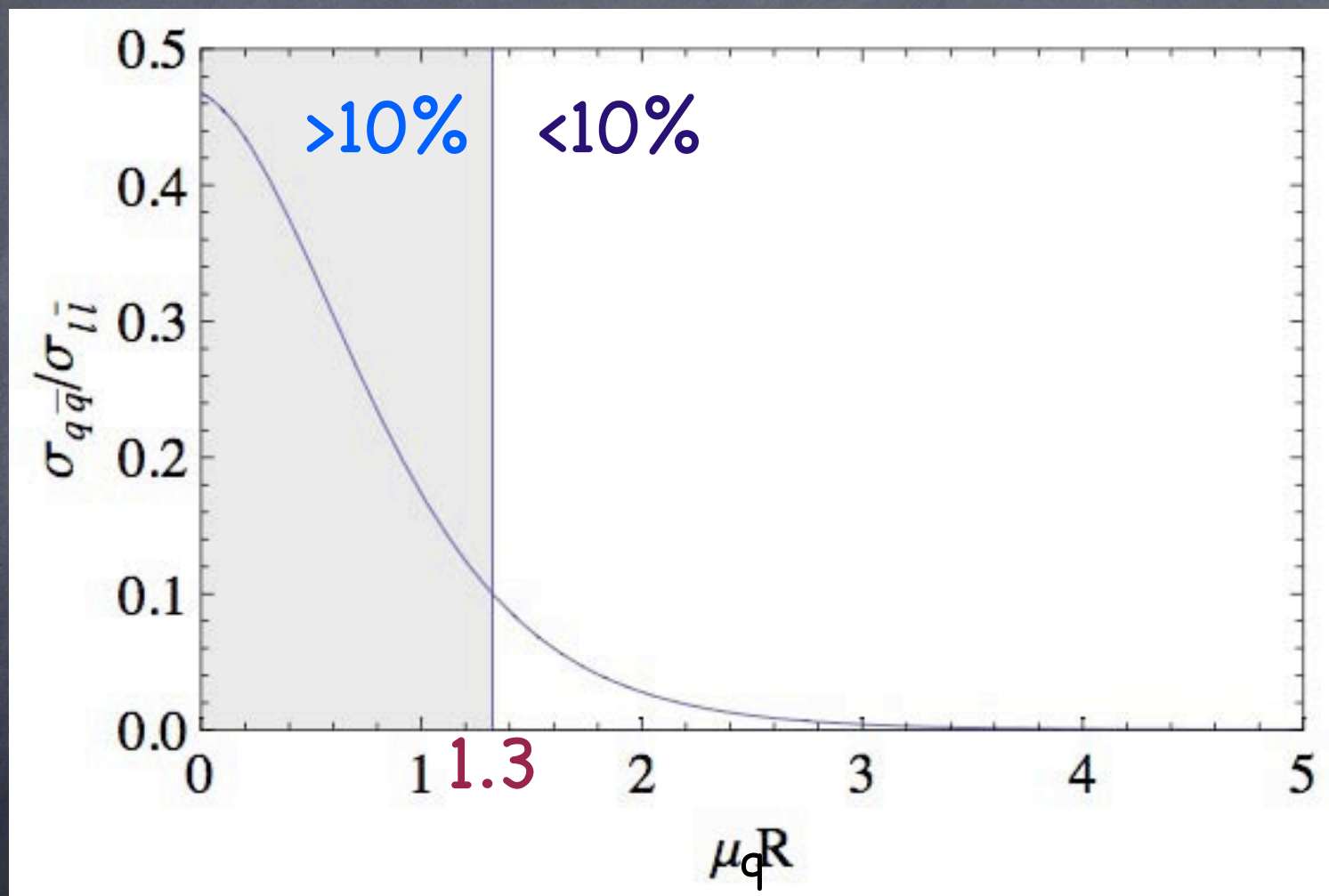
mainly  
from  $m_5$

# 1-loop

[Chen, Nojiri, SCP, Shu, Takeuchi JHEP0909.078(2009)]



# Branching Fraction $B_1 B_1$



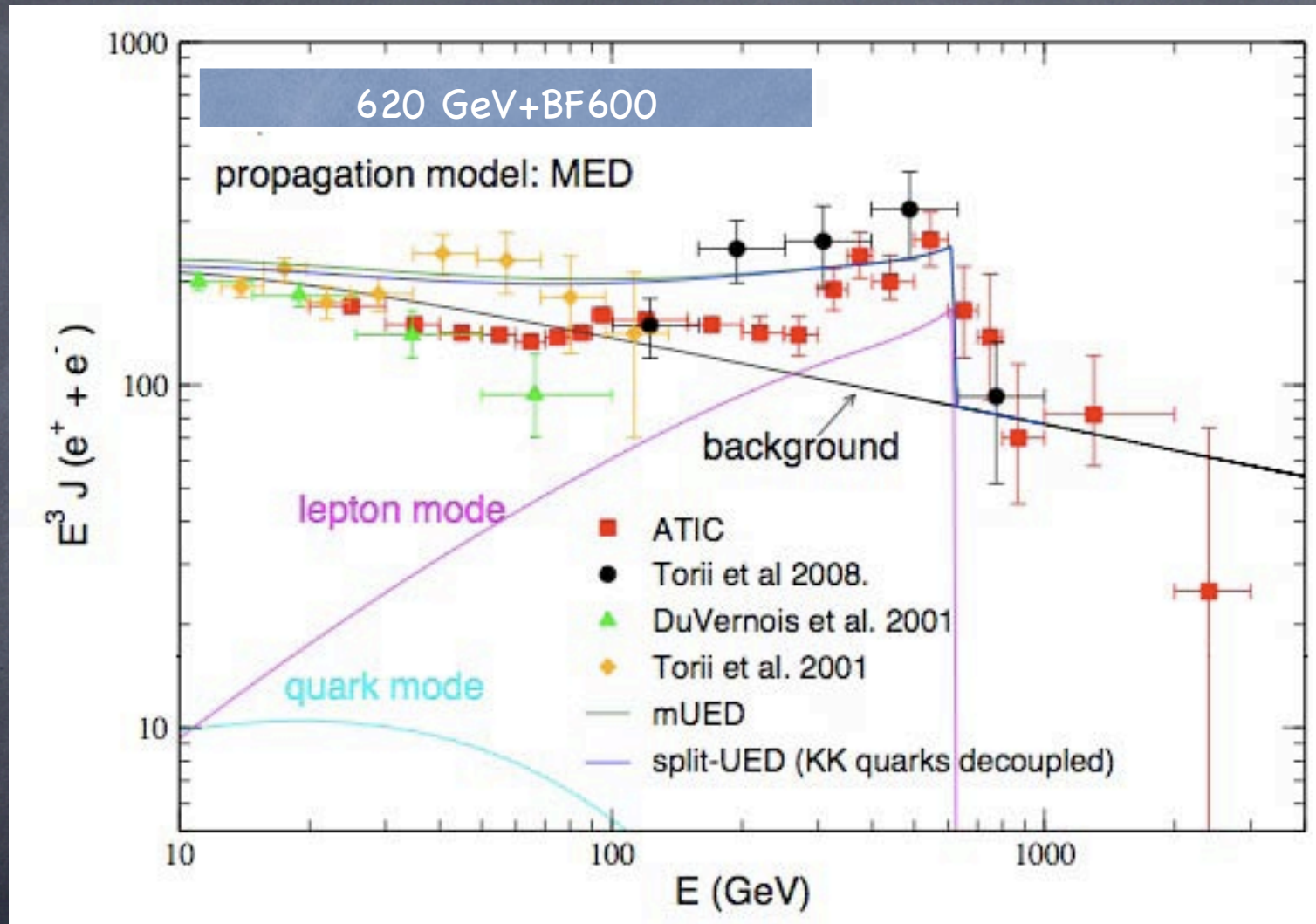
[Chen, Nojiri, SCP, Shu, Takeuchi JHEP0909.078(2009)]



# Fitting Pamela+ATIC with $m=620$ GeV



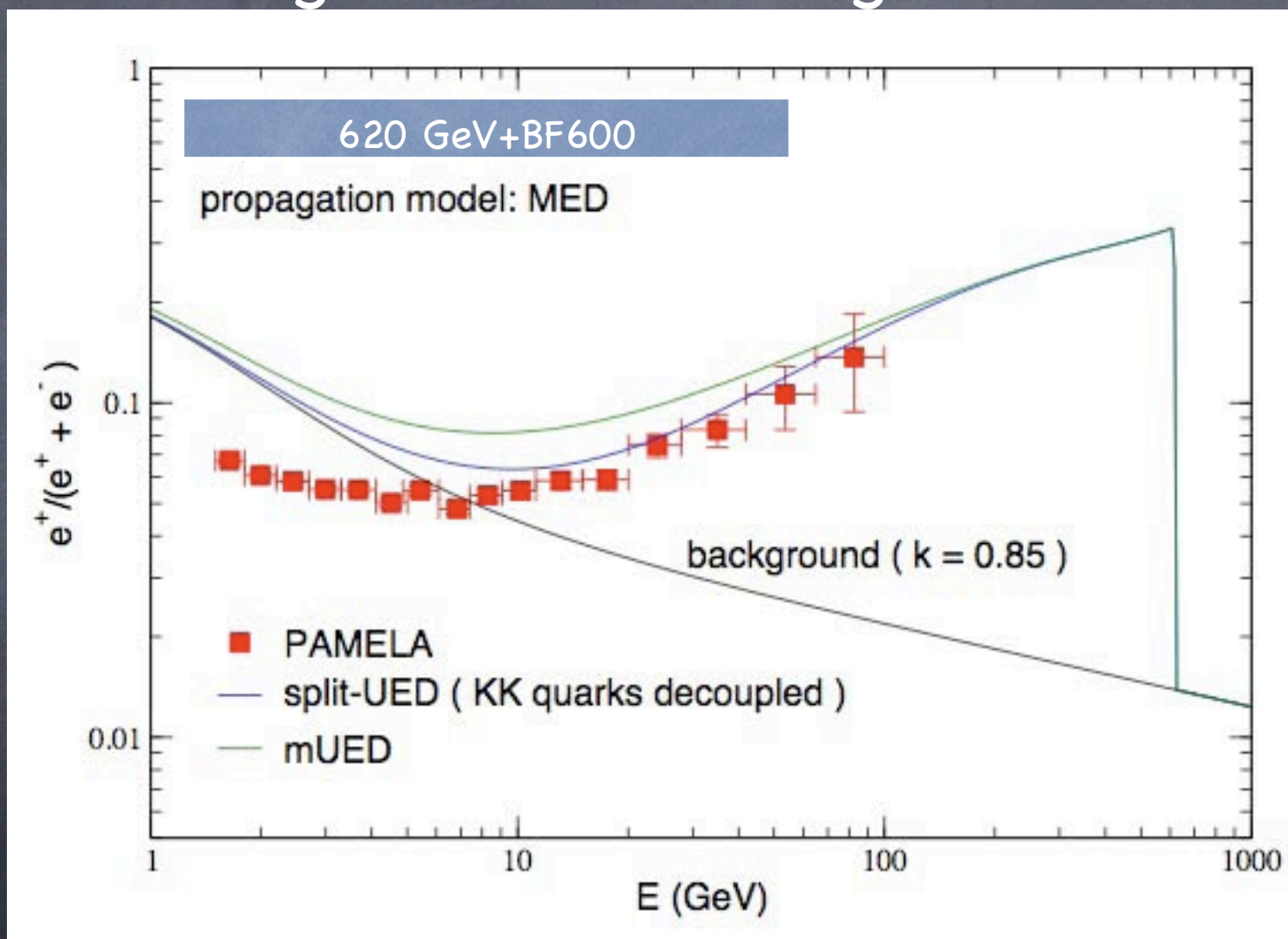
# Fitting ATIC ( $e^-e^+$ ) with 620 GeV+BF600



[Chen, Nojiri, SCP, Shu, Takeuchi JHEP0909.078(2009)]  
Here the position of peak determines  
the mass of DM: 620 GeV



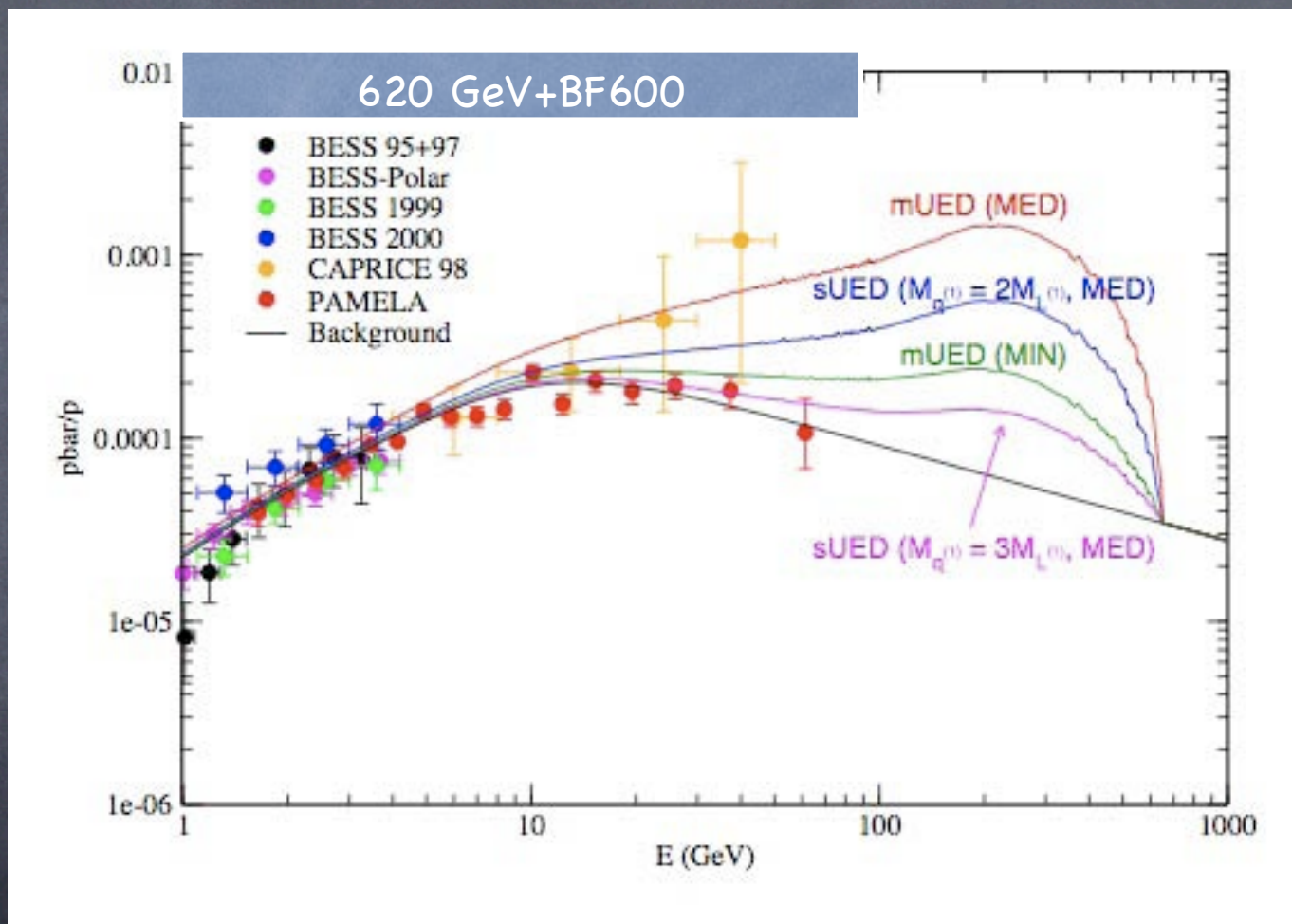
With the same mass & BF  
fitting Pamela is also good!



[Chen, Nojiri, SCP, Shu, Takeuchi JHEP0909.078(2009)]



# No Anti-proton excess



[Chen, Nojiri, SCP, Shu, Takeuchi JHEP0909.078(2009)]

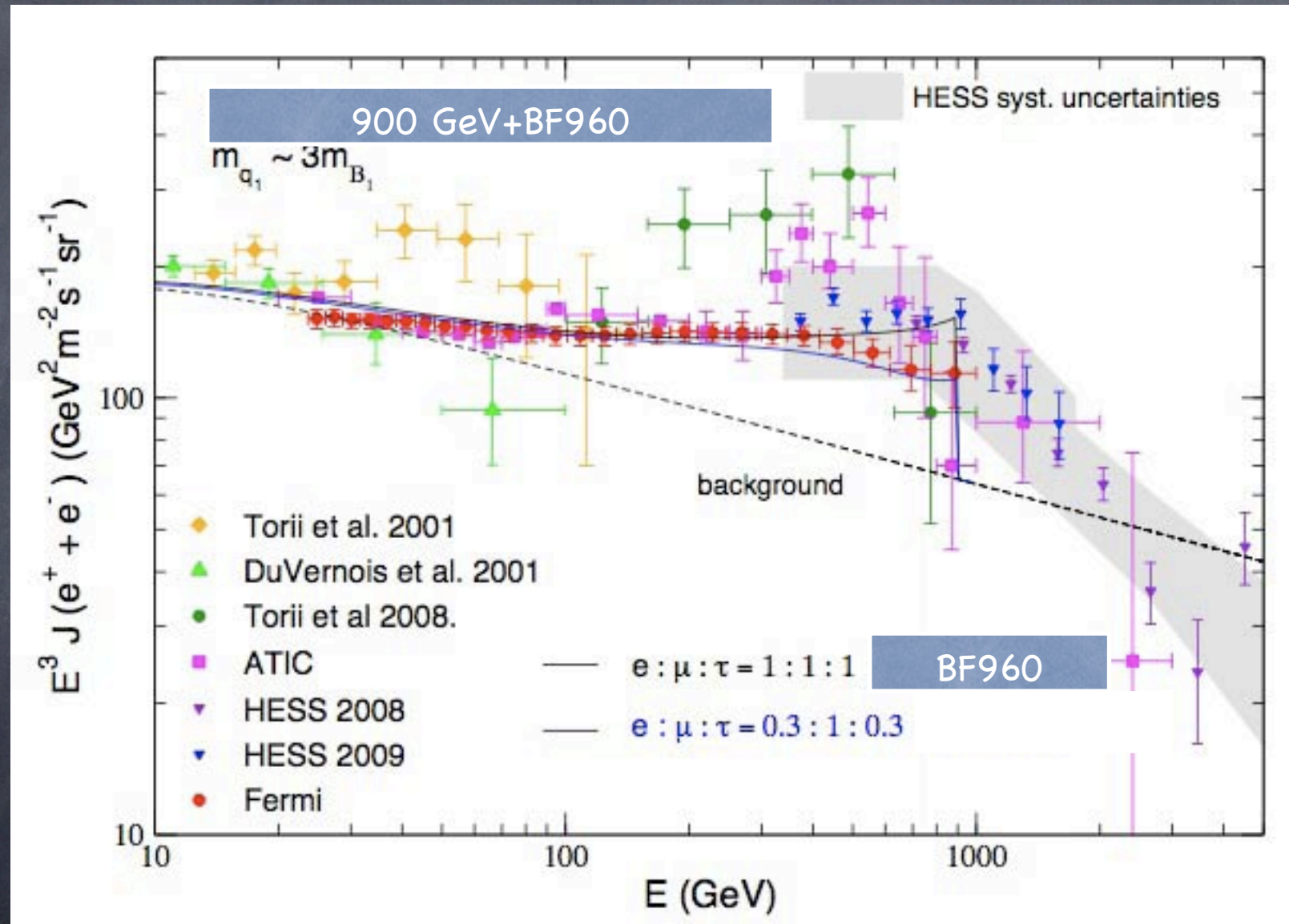


# Fitting Pamela+Fermi/LAT with $m=900$ GeV



Fermi-LAT  $e^+e^-$

KK DM provides good fit!  
with 900 GeV

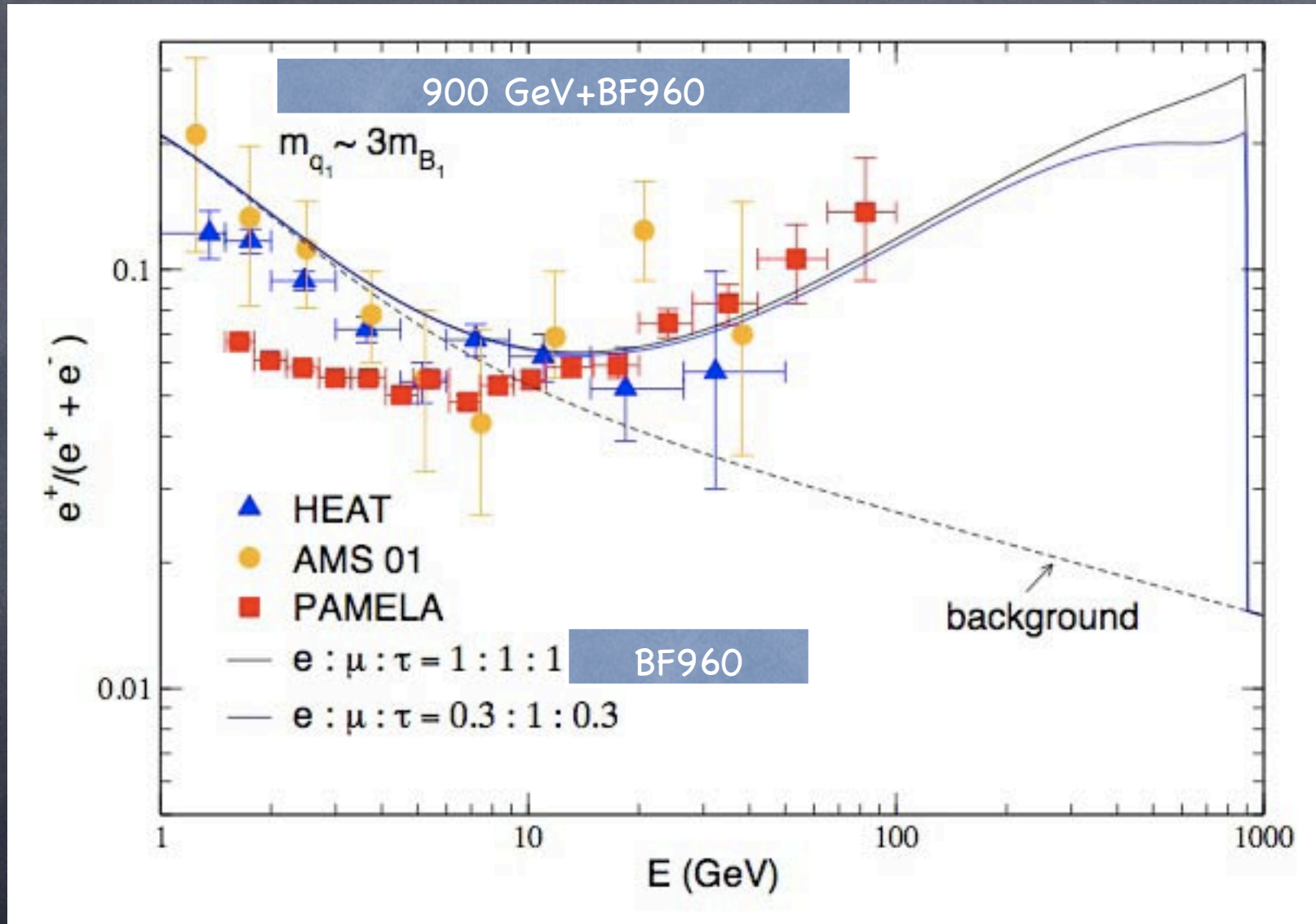


[Chen, Nojiri, SCP, Shu, arXiv:0908.4317]



# Pamela e+ fraction

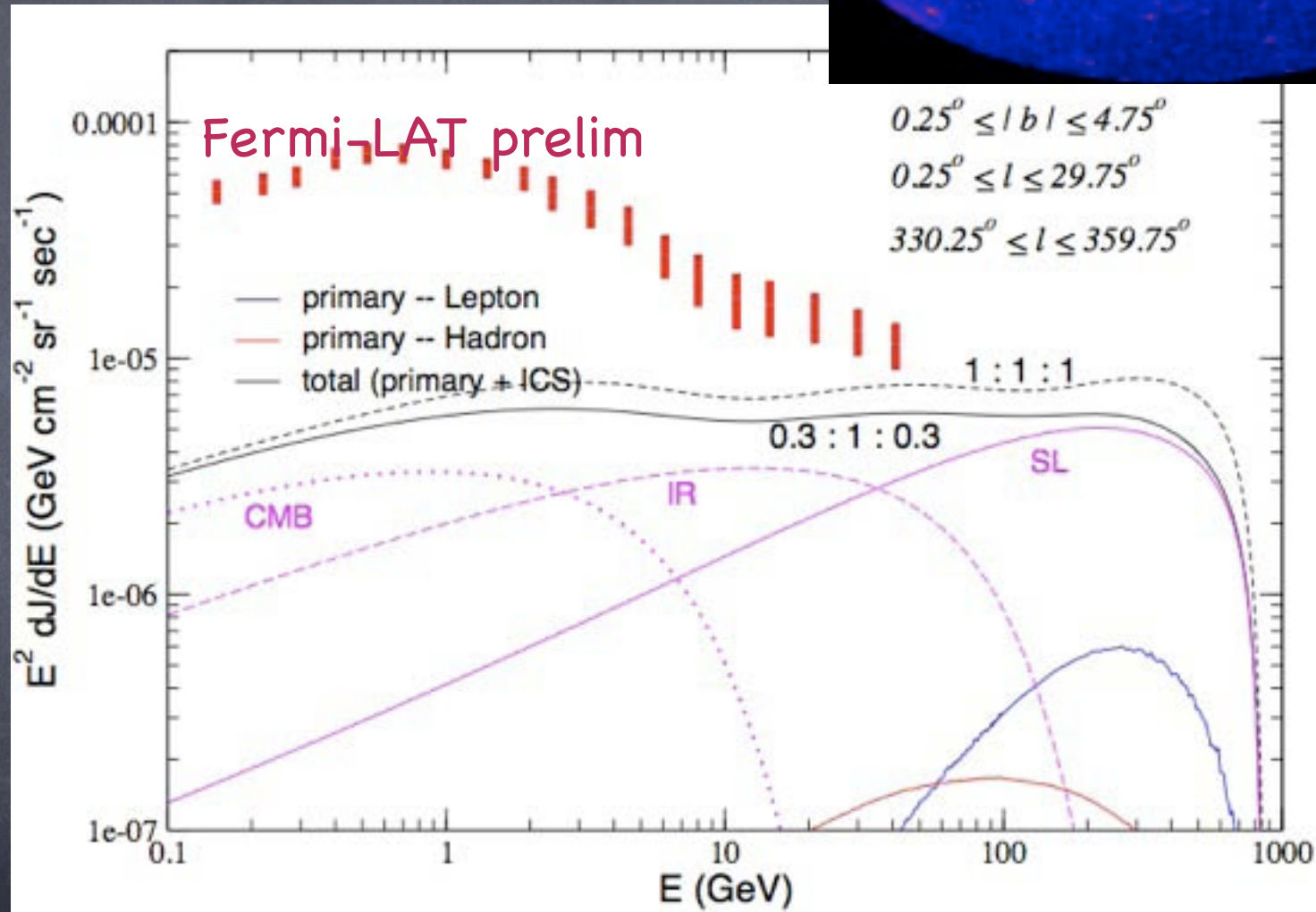
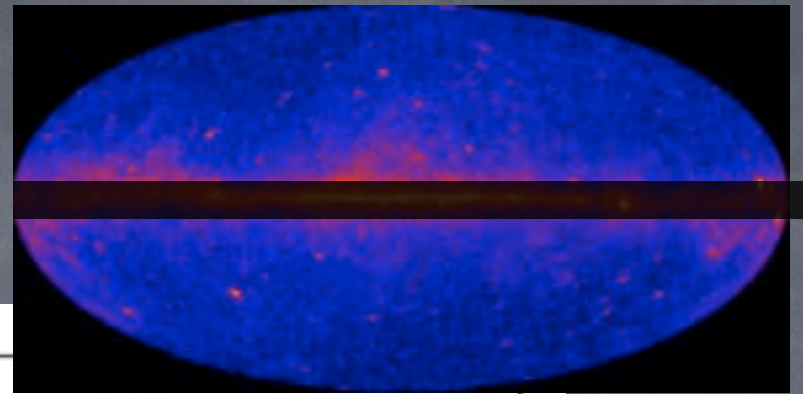
[Chen, Nojiri, SCP, Shu, arXiv:0908.4317]



[Chen, Nojiri, SCP, Shu, arXiv:0908.4317]



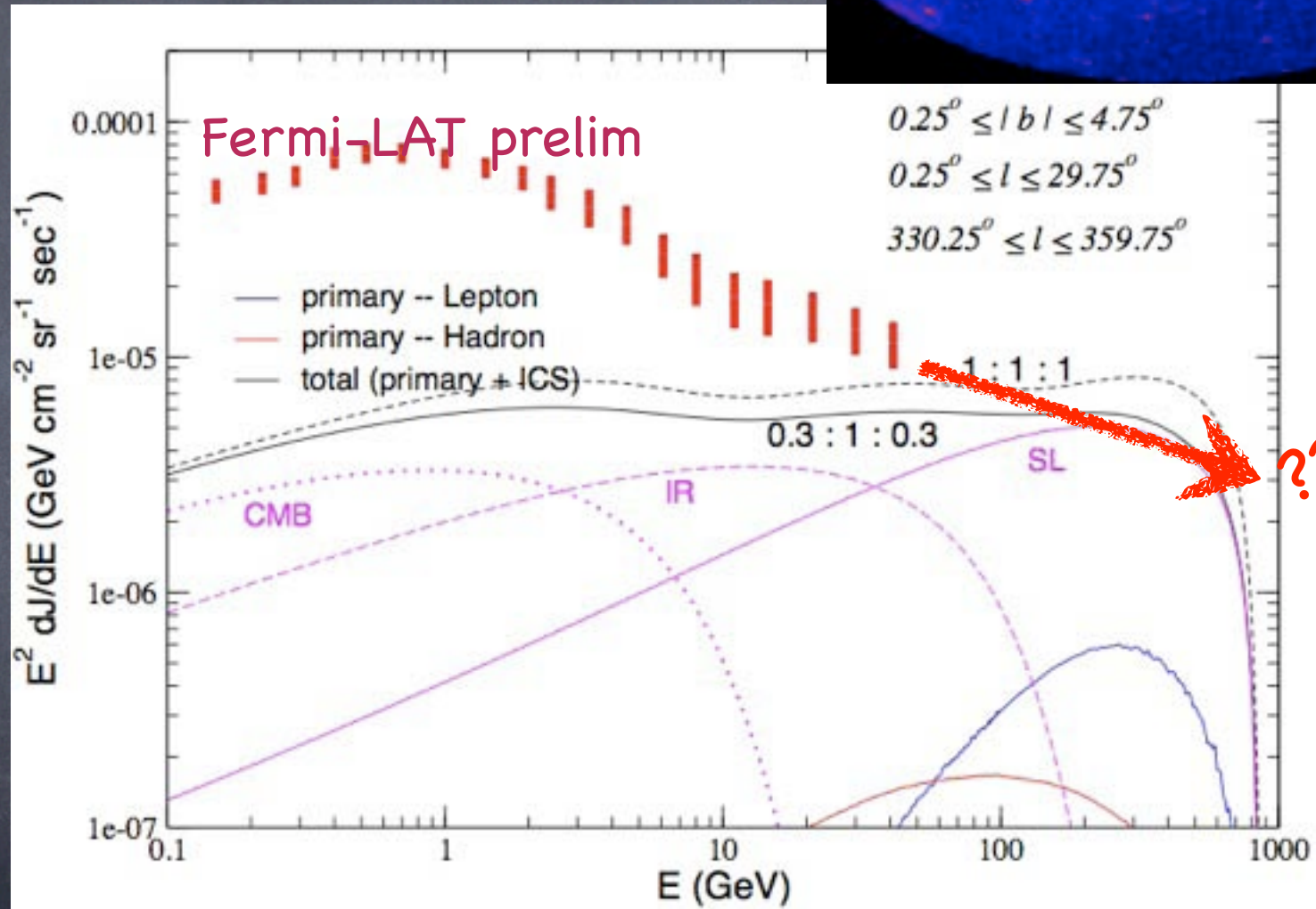
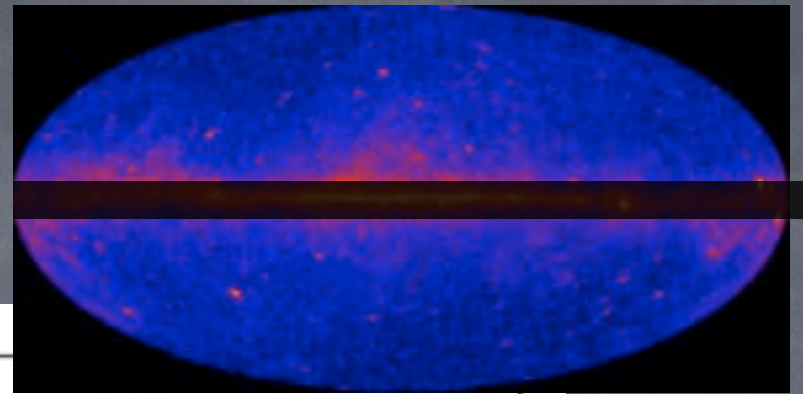
# Fermi-LAT $\gamma$ from central region $0.25^\circ \leq |b| \leq 4.75^\circ$



[Chen, Nojiri, SCP, Shu, arXiv:0908.4317]



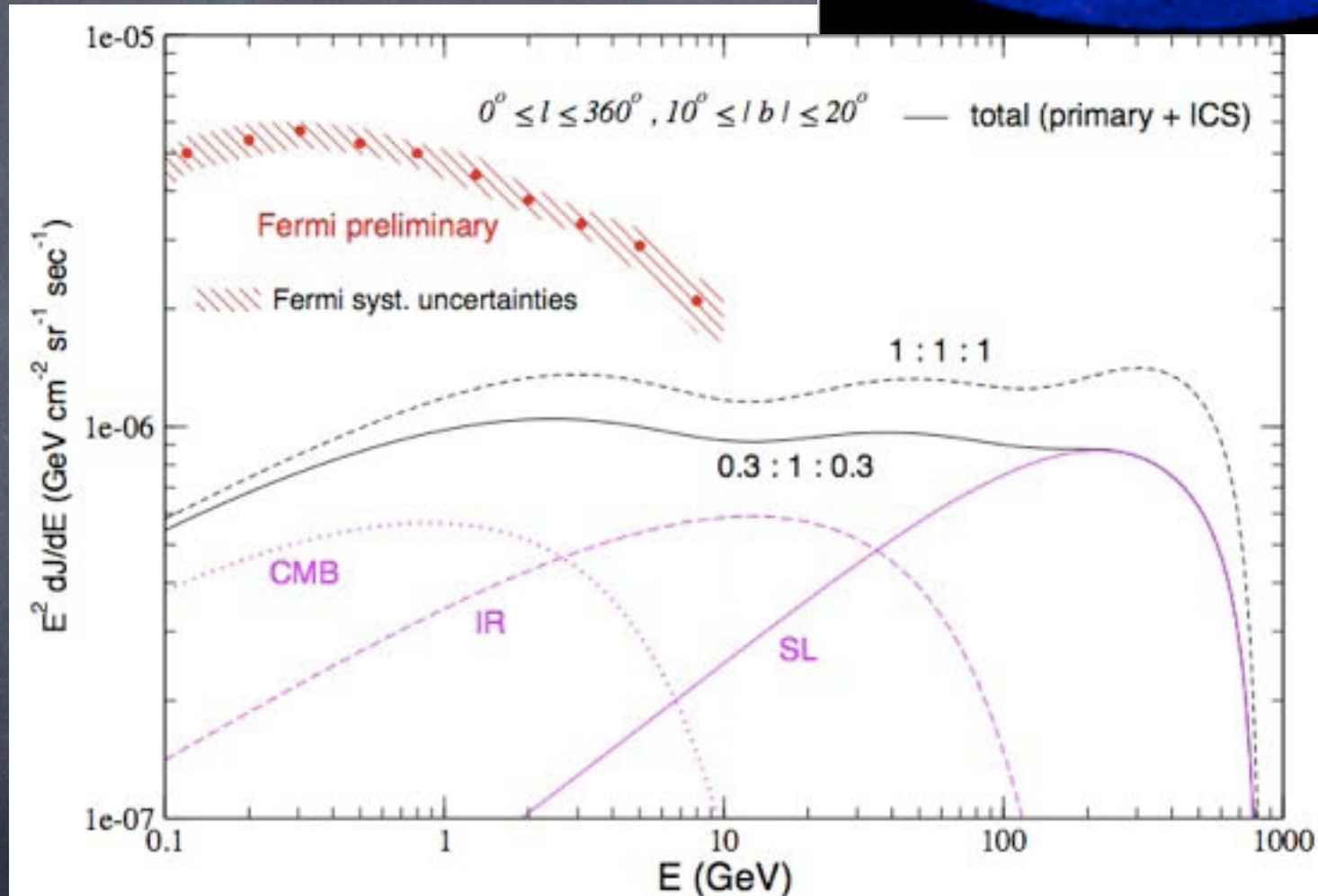
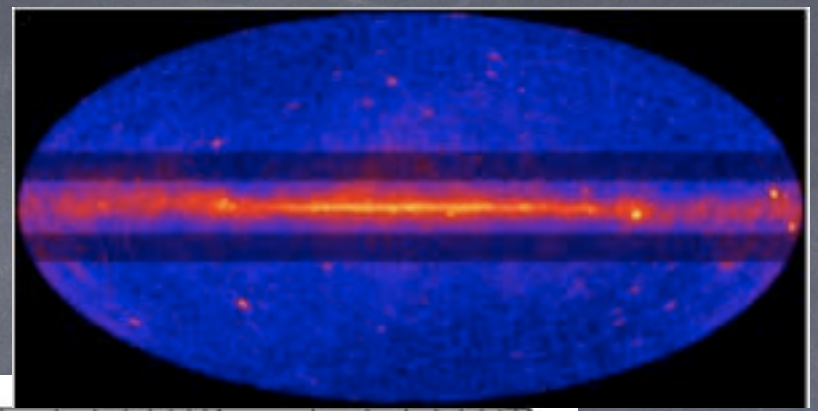
# Fermi-LAT $\gamma$ from central region $0.25^\circ \leq |b| \leq 4.75^\circ$



[Chen, Nojiri, SCP, Shu, arXiv:0908.4317]



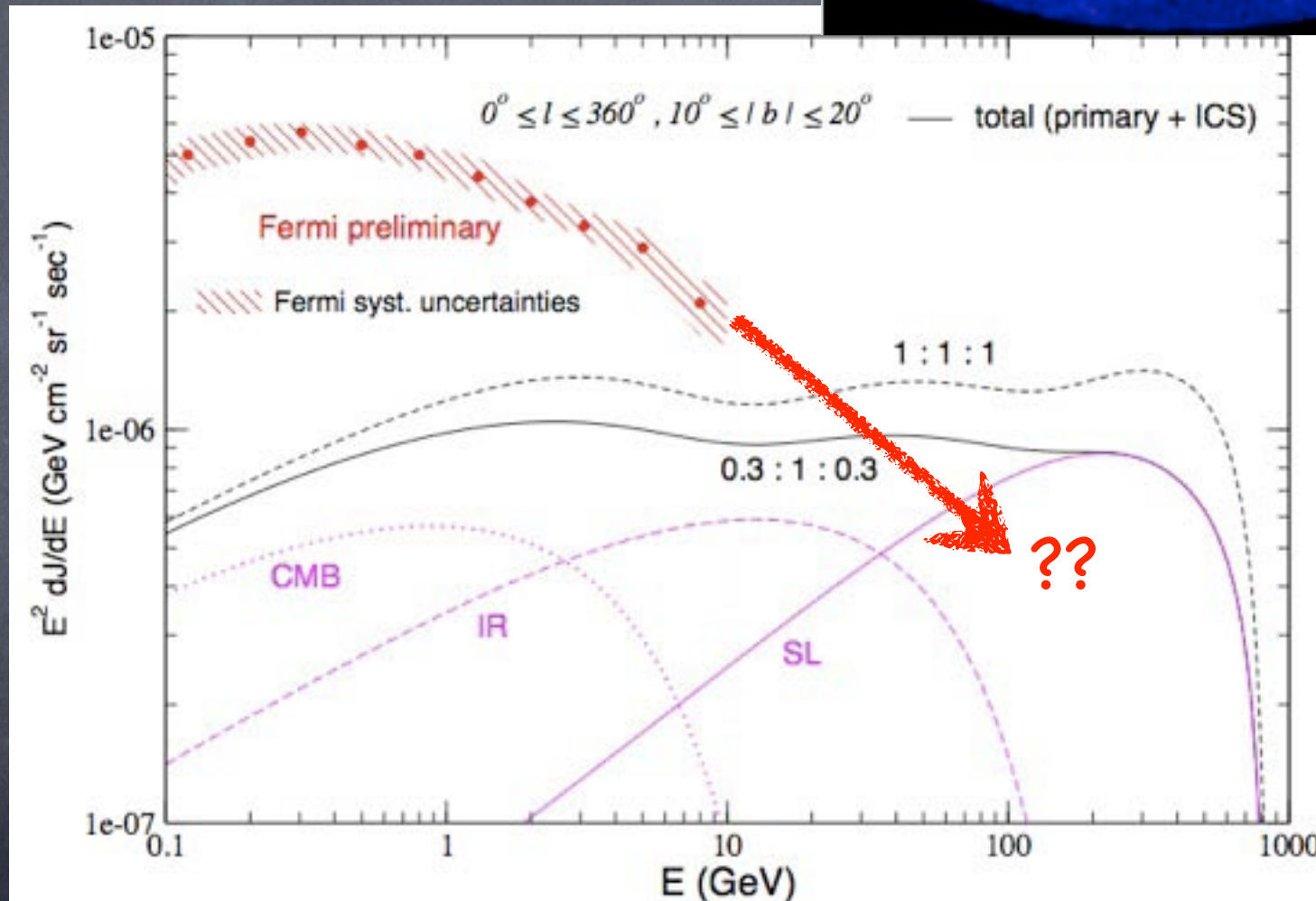
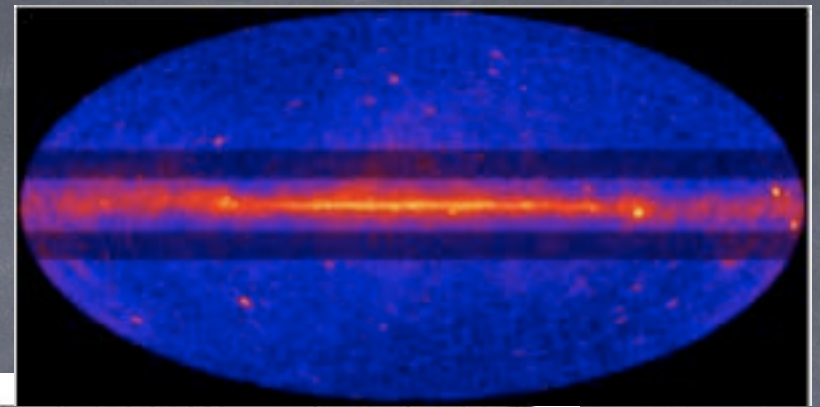
# Fermi-LAT $\gamma$ from intermediate region $10^\circ \leq |b| \leq 20^\circ$



[Chen, Nojiri, SCP, Shu, arXiv:0908.4317]



# Fermi-LAT $\gamma$ from intermediate region $10^\circ \leq |b| \leq 20^\circ$



[Chen, Nojiri, SCP, Shu, arXiv:0908.4317]



# Conclusion/Discussion

- In extra dimension models with KK-parity LKP can be a natural WIMP.  $B_1$  is the LKP in UED models.
- Leptophilic  $B_1$  in split-UED can fit PAMELA, FERMI (or ATIC) with  $1/R=900$  (or 600), which provides the right relic abundance of DM.
- Gamma-ray excess in  $E_\gamma > O(100)\text{GeV}$  can be a signal of KK dark matter. The LHC will also prove or disprove this idea.